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INVESTIGATION OF THE EXTERNAL FLOW ANALYSIS
FOR DENSITY MEASUREMENTS AT HIGH ALTITUDE

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NASA RESEARCH GRANT NSG 1630



Department of
Mechanical and
Aerospace Engineering

INVESTIGATION OF THE EXTERNAL FLOW ANALYSIS FOR
DENSITY MEASUREMENTS AT HIGH ALTITUDE

FINAL REPORT

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I. INTRODUCTION

Accurate experimental determination or verification of aerodynamic force coefficients (C_D and/or C_L) requires accurate simultaneous measurements of the forces (or accelerations) and the dynamic pressure ($q = 1/2 \rho U^2$). Comparison with theoretical predictions requires independent knowledge of the density (ρ) and velocity (U) to establish the proper values of the non-dimensional parameters such as Reynolds number (Re) and Mach number (M). These parameters in principle require independent measurement of temperature and measurement or inference of viscosity. Under hypersonic conditions during the early phases of re-entry the Mach number becomes a secondary parameter while the relevant Reynolds number is based on viscosity within the gas layer near the vehicle and is only weakly dependent on free stream temperature. During the earliest part of re-entry, independent knowledge of density is necessary to establish the degree of rarefaction generally measured by the Knudsen number $Kn = \lambda/L$ where λ is a relevant mean free path and L the characteristic physical dimension (either vehicle size for the overall flowfield or entrance dimensions for the local behavior at the instrument).

The Shuttle Upper Atmosphere Mass Spectrometer (SUMS) Experiment⁽¹⁾ is designed to provide independent measurement of $q = 1/2 \rho U^2$ within the high altitude range. When combined with information of vehicle velocity, it will provide independent determination of upper atmosphere density and

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coupled to accelerometer data will give the aerodynamic force coefficients within a regime difficult to simulate on the ground. The experiment is primarily intended to provide information between about 80 Km and 140 Km where rarefaction effects on the force coefficients are most important for a vehicle of the size of the Space Shuttle. It is also a regime where information on the atmosphere is relatively sparse as it lies below the altitude traversed by satellites and above that regularly assessed by ground launched meteorological vehicles. The interpretation of the measurements, however, requires an adequate understanding of the flowfield around the Space Shuttle within the vicinity of the SUMS experiment in order to provide the proper data reduction procedure and an assessment of the accuracy of the results.

At sufficiently low altitudes (below about 80 Km for the Space Shuttle), conventional pitot probe measurements can provide the dynamic pressure with straightforward data reduction and relatively minor corrections. At sufficiently high altitude (above about 150 Km), free molecular theory can be used to infer free stream conditions from surface measurements. The forces, however, are small and of little interest while the measurements require instruments of high sensitivity and are therefore difficult. In the intermediate range of altitudes where SUMS is designed to provide data, the typical molecular mean free path is of the same order as the characteristic vehicle dimensions. Figure 1 shows the

variation with altitude of the free stream Knudsen number ($Kn_\infty = \lambda_\infty/D$) based on the free stream mean free path λ_∞ and the Shuttle diameter D at the location of the SUMS orifice just ahead of the wheel well. Note that $Kn_\infty = .01$ at about 87 Km and $Kn_\infty = 10.0$ at about 136 Km. Intermolecular collisions can, therefore, neither be neglected (free molecular theory) nor represented by the resultant transport properties (continuum theory) over the major portion of the SUMS measurement regime. The gas properties at the entrance to the instrument are, therefore, dependent on a flowfield that can only be determined on the basis of a "molecular" theory.

In addition to the above "external flow problem", needed to establish properties near the surface of the vehicle, the entrance region of the instrument is typically either smaller or comparable to a local mean free path. In such circumstances, the connection between the gas properties at some distance into the internal plumbing and those at the vehicle surface, can be very sensitive to the velocity and angular distribution of the incoming molecules. This requires both a high degree of detail from the "external flow" and a local analysis that must assess the molecular behavior at the instrument entrance. We shall refer to this as the "entrance problem". Figure 1 also shows an approximate band of Knudsen numbers for the entrance region of the SUMS experiment. Kn_s is based on the mean free path λ_s at the vehicle surface and the orifice diameter ($d_o = .235$ cm) with surface properties fitted between free molecular results at high altitude and continuum Newtonian values at low altitude. Kn_c is a similar Knudsen number based on

conditions behind the "entrance" tube with λ_c estimated on free molecular results using Hughes and deLeeuw theory⁽²⁾ at high altitude, with continuum constant pressure results applied at low altitude. Note that conditions within the entrance tube range from clearly free molecular behavior above about 110 Km to transitional behavior near 80 Km with fully continuum results only approached at the lowest altitude of interest.

The subsequent connection between the properties immediately behind the entrance tube and the measurements at the mass spectrometer shall be referred to as the "internal problem". The analytical procedures for calculating pressure profiles through the internal plumbing are well established and will be further verified by instrument calibration⁽¹⁾.

During the preliminary phase of NASA Grant NSG 1630 (July 1979 to November 1979), the feasibility of examining the "external flow problem" for the Space Shuttle nose region within the relevant altitude range was established. A previously developed Direct Simulation Monte Carlo Computer Code^(3,4,5) was found to be suitable as the starting point for this geometry and altitude range. Preliminary results were obtained at 87, 95, 105, 115 Km altitudes.

During the subsequent grant periods (November 1979 to September 1982) improvements in the modelling of the geometry and the molecular interactions have been incorporated in the external flow computer code. A number of runs at altitudes of 87, 95, 105, and 115 Km have been made to obtain a range of the relevant parameters and to provide input information

at the SUMS entrance location. The "entrance problem" has been examined both by using published information (2,6,7) , and a previously developed Monte Carlo code for internal geometries (8) . Because of the combination of entrance geometry (very long tube) and the range of local Knudsen number over the altitudes considered, a totally new "entrance" computer code had to be developed. This code provides the connection between the flux information at the orifice entrance obtained from the external code and the local gas properties behind the entrance tube where the gas is in equilibrium with the "cold" walls of the internal plumbing. This new code has only been exercised to a limited extent, but preliminary results relating the pressure within the tube behind the tile to the free stream dynamic pressure have been obtained. This information coupled with an appropriate calibration of the mass spectrometer provides the basis for a viable data reduction procedure of the SUMS experiment.

Section IIA contains a brief description of the operation of the EXTERNAL computer code (the detailed code is attached in Appendix A). Section IIB describes the issues associated with geometric modelling of the shuttle nose region and the modelling of intermolecular collisions including rotational energy exchange and a preliminary analysis of the vibrational excitation and dissociation effects. Section IIC discusses the selection of the trial runs and presents the major results.

Section IIIA contains a brief description of both the

original version and the modified present code (INTERNAL) for the entrance problem (Appendix C contains the code listing). Section IIIB contains a discussion justifying the selection of geometric, collisional and surface modelling parameters used for the trial runs. Section IIIC presents the preliminary results and discusses the major effects.

Section IV presents the conclusions that can be drawn from the present study, provides a preliminary estimate of the data reduction procedure and suggests future work.

II. EXTERNAL FLOWFIELD

The physical properties of the gas monitored by the SUMS instrument are not those of the free stream but are altered both by the intermolecular interactions in the external flowfield and by the combination of intermolecular and surface interactions within the entrance orifice and tubing leading to the instrument. External flowfield effects can be summarized in terms of the relation between local "stagnation" pressure and the free stream dynamic pressure ($q = 1/2\rho U^2$) at sufficiently low altitudes. Within the "transition" regime of interest (80-140 Km) the gas entering the orifice is neither in equilibrium with the surface nor simply related to the free stream. The only currently available technique for describing the flowfield within this regime and providing sufficiently detailed data on the physical state of the gas at the surface is the Direct Simulation Monte Carlo Code.

A. DIRECT SIMULATION MONTE CARLO COMPUTER CODE

"Monte Carlo" is the technique of using a simulated situation and random numbers to generate solutions from which information for the real case is then deduced statistically. The Monte Carlo approach ranges from being a strictly mathematical technique for evaluating the complicated multi-dimensional Boltzmann collision integral to a complete simulation of a number of molecules, with randomness

introduced only in the initial conditions. A modification of this latter approach is the one used in the present development. It consists of simultaneously following a large number of particles which yields, to some degree, a "direct simulation" of the processes taking place. Because there are finite limits on computer storage space, a modification to the direct simulation technique was developed by G.A. Bird (Ref. 5) wherein the real gas is simulated by several thousand "sample" particles populated into cells of the sample space considered. For collision calculations, all the particles in one cell are used as a representation of the local distribution function from which collision pairs are chosen at random, but in proportion to their collision probability based on the real gas. This preceding discussion applies to a general program incorporating the direct simulation procedure. A specific computer program for the generalized three-dimensional program for axisymmetric bodies in a hypersonic multi-fluid flow is described below.

The program (EXTERNAL) conducts numerical experiments with a model multi-component gas. The real gas is simulated by several thousand molecules which may be thought of as a representative sample of the many billions of molecules in the corresponding real gas. The positions and velocity components of the simulated molecules are stored in the computer and typical collisions are computed among them as a time parameter is advanced. Since the flow is at an angle of attack to the body, three position coordinates, three velocity components and appropriate internal energy levels

must be stored for each simulated molecule.

The computation of collisions starts at zero time with the molecules moving along the flow axis at the required freestream Mach number. The body is inserted into this flow at the zero time and the desired steady flow is obtained as the large-time solution of the resulting unsteady flow.

The free-stream flow vector lies in the x-y plane. The simulated region is bounded by the x-y plane as a plane of symmetry, an outer cylindrical boundary (the axis of the cylinder is the x axis), and two planes parallel to the y-z plane. These boundaries must be set sufficiently far from the body to eliminate interference. The simulated region is divided into a number of cells which are sufficiently small for the expected change in flow properties across the cells to be small.

The first step is to generate the initial, or zero time, configuration of molecules. The molecules are distributed over the simulated region and the velocity components assigned are appropriate to a gas in Maxwellian equilibrium and moving at the required Mach number. The body is then inserted into the flow and the molecules are allowed to move and collide among themselves. The move and collide processes are uncoupled by computing a number of collisions appropriate to a time interval Δt_m equivalent to a small fraction of the mean time between collisions, and then moving the molecules through distances appropriate to Δt_m and their instantaneous velocities. The distortion produced in the molecular paths by this approximation is small as long as Δt_m is small

compared with the mean time between collisions, and smaller than the typical transit time of a molecule through a cell.

Since the change in flow properties over the width of one cell is assumed small, the molecules in a cell at any instant may be regarded as a sample of the molecules at the location of the cell. The relative location of the various molecules within the cell can then be disregarded. A pair of molecules is chosen at random from those within the cell under consideration and is retained or rejected in such a way that the probability of retention is proportional to the relative collision probability for the appropriate interaction law. When a pair is retained, a typical collision is computed between the two molecules and the new velocity components and internal energies are stored in place of the old ones.

In general, the relative number ratio of the species of molecules in the multi-component gas will differ from unity, requiring the computation of different types of collisions. There is, therefore, one time counter for each type of collision in each cell. For each collision, the correct time counter is advanced for the cell by an amount appropriate to the collision parameters. The probability of collision, and therefore the time advancement per collision, is made proportional to the number of molecules in the cell, and the relative velocity and cross-sections of the colliding molecules. Collisions are computed in each cell until all the time counters have advanced through at least a time Δt_m . When this procedure has been carried out for all cells, the

overall experiment time is advanced by Δt_m and the molecules are moved through appropriate distances.

The set of molecules in each cell changes as the molecules are moved and appropriate conditions must be applied at the boundaries of the region being simulated. Every boundary is treated as a source of molecules with velocity components representative of molecules moving in thermal equilibrium at the appropriate fraction of the free-stream Mach number. (The fractional Mach number is determined by the cosine of the angle between the local boundary normal and the flow direction.) Any molecule which moves outward across any boundary is regarded as being lost and is removed from the sample. The plane of symmetry (the x-y plane) is regarded as a specularly reflecting surface. Interactions with the body are also computed. The body consists of a number of conic sections rotated about the axis of symmetry. Each section must be separately specified according to the coefficients of the defining equation, a procedure to be described later in the report. For the purpose of computing the momentum and energy transfers to the surface, each region of the body can be subdivided into smaller sections. Within these smaller sections, the following three parameters must be specified: wall temperature/gas temperature, energy accommodation coefficient for each species, and tangential (momentum) accommodation coefficient for each species.

After the flow has settled down to a steady state, the number flux, momentum and energy transfers to the surface are

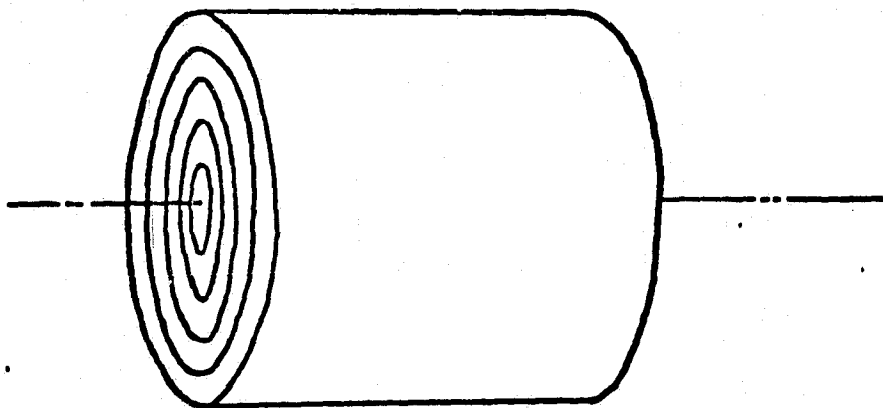
accumulated and used to compute the aerodynamic data. The time required to establish steady flow is usually assumed to be close to the time required for the free-stream flow to traverse several body lengths. The overall number flux, drag, and heat transfer coefficients are determined, along with their distribution along the surface.

In addition, it is possible to generate data on the body surface which can be used as input to the companion program, INTERNAL, described separately. (INTERNAL) computes the flow field regime inside an axially symmetric cavity, such as might be used for a spacecraft-borne sensor. The input data needed for this computation includes the molecular distribution functions present at the orifice to the cavity, the orifice being on the surface of the spacecraft.) This data consists of velocity and internal energy samples in three coordinate directions for all species of the mixture.

Flow field properties are also computed. Instantaneous values are sampled at appropriate time intervals and these are time-averaged for greater accuracy. Number densities, velocities and temperatures are printed for each cell.

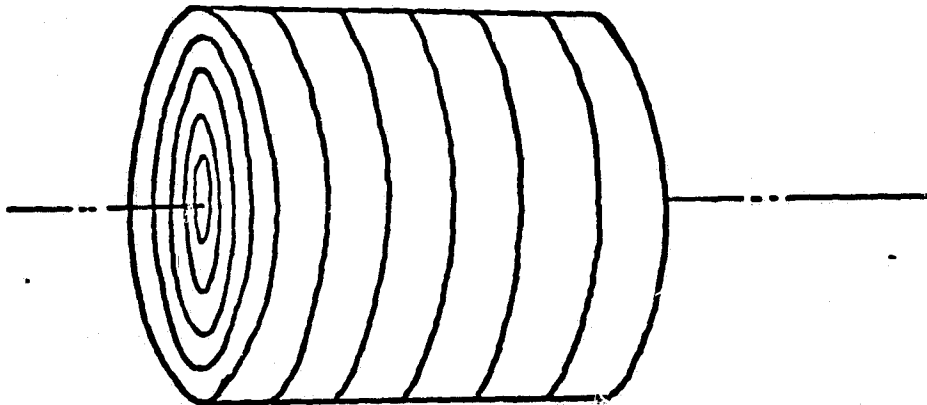
The numerical experiment takes place in a cylindrical block of space whose axis is coincident with the axis of revolution of the conic surfaces comprising the test body. This space is subdivided into cells in which the flow field properties of the experiment can be monitored. Cylindrical surfaces concentric with the axis partition the space into nested cylindrical volumes, as shown:

ORIGINAL PAGE 18
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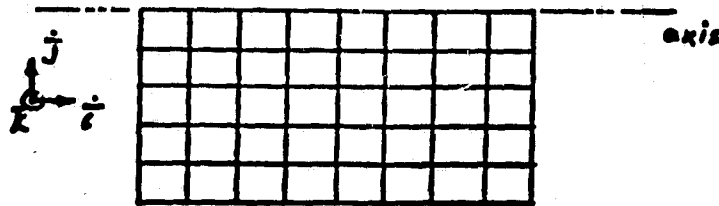
Planes parallel to the end-faces of the cylinders divide the cylinders into a stack of nested rings.

ORIGINAL PAGE 17
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Finally, planes perpendicular to the preceding planes and passing through the axis, called radial planes, divide up the rings in the azimuthal direction, producing cells which are wedge-shaped pieces of rings. The geometry is more easily visualized if one considers a trace of the cell configuration in a radial plane. A typical planar trace is shown:

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The axis shown is the axis of symmetry of the sample space (and of the test body). This axis is considered to be the x axis. When this direction must be specified in vector algebra computations, a unit vector \bar{i} is assumed in the x direction.

Assume that the planar trace shown above is bounded by the x -axis as described and by the $-y$ axis. This plane is at 0° azimuth angle and is called the zero plane. It is the plane normally depicted when describing the sample space.

A unit vector \bar{j} points in the $+y$ direction. The z direction points out of the paper, and in this direction is the unit vector \bar{k} , which is given by $\bar{i} \times \bar{j}$.

The flow velocity is in the direction $\bar{i} \cos \alpha + \bar{j} \sin \alpha$, where α is the defined angle of attack between 0° and 90° . Since the flow is thus in the xy plane, there is symmetry in the z direction. That is, any condition in effect at $+z$ is also in effect at $-z$. Thus azimuth angles need be specified only from 0° to 180° where 0° is in the $-y$ direction, 90° is in the $+z$ direction, and 180° is in the $+y$ direction.

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Now the fact is that the gas density in the vicinity of the stagnation point of the body can become many tens of times higher than the density far from the body. It is thus desirable to use small cells in this region while the cells are larger in the regions of relatively low density. This partitioning of cell sizes is accomplished in two ways.

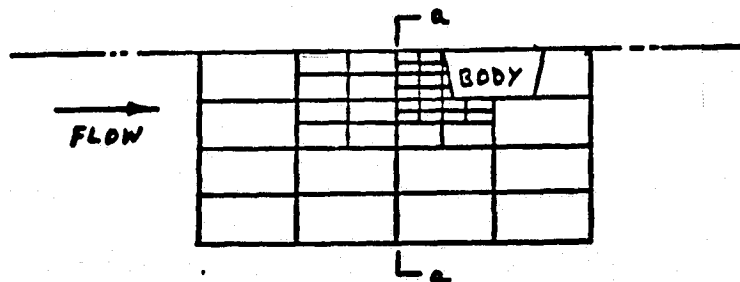
First, the rings can be divided azimuthally into two different sizes. This is done by specifying an angle, called THETAZ, and the number of wedge divisions both below and above this theta plane, called NWEDGE₁ and NWEDGE₂. ("Below the theta plane" means azimuth angles between 0° and THETAZ, and "above the theta plane" means azimuth angles between THETAZ and 180°.)

Second, the cells below the theta plane can be subdivided in the axial and radial directions down to second and then third level cells. In the 0° radial plane representation of the sample space, this would appear as large rectangles being sub-divided into small rectangles.

In this way, the sample space geometry can be tailored to the configuration of the test body angle of attack to the flow. The following examples are presented to clarify the above statement. Assume for all cases that the test body in question is a short cylinder. As explained in the section TEST BODY, the cylinder cannot have flat end faces, so the ends are cones with apex angles of about 175°.

- a) For 0° angle of attack, the flow impinges directly on the left face of the cylinder.

It is thus desirable to have, if possible, constant azimuthal angles since there is no angles since there is no azimuthal assymetry. One possible configuration is therefore: $\text{THETAZ} = 180^\circ$ $\text{NWEDGE}_1 = 6$, $\text{NWEDGE}_2 = 0$ (producing 30° wedges), and the axial and radial directions can be subdivided any convenient way, producing a zero radial plane that looks like:



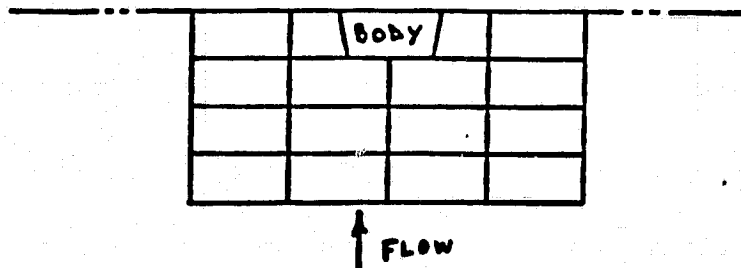
In this type of geometry, all radial planes are the same as the zero plane.

b) For 90° angle of attack, the flow impinges on

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the curved cylindrical surface of the cylinder at the bottom. It is thus necessary to have small azimuthal wedges on the lower portion of the cylinder at and near the stagnation point, while larger wedges will suffice on the upper portions of the body (in the wake of the flow).

For instance, an acceptable set of parameters is: $\text{THETAZ} = 60^\circ$, $\text{NWEDGE}_1 = 3$, $\text{NWEDGE}_2 = 3$ (producing 20° and 40° wedges), and again the axial and radial directions can be subdivided in any convenient way. Any radial plane up to 60° looks like the radial plane in example (a), with the body and smaller cells centered about a-a, while any radial plane between 60° and 180° looks like:



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- c) For an angle of attack between 0° and 90° , the configuration looks generally like that of example (b)). In this case, however, the theta plane generally should be at an azimuth angle which is relatively low (near 60°) for a high angle of attack ($45^\circ - 90^\circ$) and relatively high (near 120°) for a low angle of attack.

Because the test body is located on the axis of the cylindrical sample space, for each particle that interacts directly with the body, many more do not. In the interest of minimizing the program running time necessary to permit a statistically sufficient number of particles to strike the body, the computation makes use of zonal weighting factors. That is to say, each particle in the sample-space in reality represents one or more particles, the actual number depending upon the weighting factor of the zone in which the particle currently exists.

Up to five cylindrical boundaries are selected across which the zonal weighting factors change. These boundaries

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are specified in terms of the number of first level cells between the axis and the boundary. The change in the zonal weighting factor across each boundary can be given by:

$$LF_n = \frac{LD_{n+1} + LD_n}{n-1} \quad n=1,2,\dots,5$$

$$\prod_{j=0}^{n-1} LF_j$$

where the LD values are the number of first level cells between the axis and the cylindrical boundaries; and $LF_0=1$, $LD_6=LD_5=LD_K=NH$ where $K \geq K_{last}$.

This equation is the result of having the zonal weighting factors defined in such a way that they are equal to the ratio of sample space volume in the zone above the weighting-factor boundary to the sample space volume in the zone below the weighting factor boundaries. The importance of the region near the axis can be emphasized by choosing LF values larger than those given above.

The sample space is populated with a distribution of gas molecules. Each molecule is assigned a velocity, a rotational

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energy and a position, such that the sample space is uniformly (albeit in a random manner) filled with a gas in thermal equilibrium and flowing at the required Mach number. Each molecule is assigned a number corresponding to its zonal weighting factor. The molecule thus represents in actuality a number of molecules (including itself) equal to the zonal weighting factor. While the molecule moves in such a way as to stay within the given zone, its weighting factor does not change. If it crosses a weighting factor boundary while moving in toward the axis, a number of molecules is added to the distribution. The number of molecules added is equal to:

$$\frac{\text{old weighting factor}}{\text{new weighting factor}} - 1$$

For instance, assume that above a boundary, the zonal weighting factor is 6, and below the boundary it is 3. Hence $6/3 - 1$, or 1, molecule must be added to the distribution when the molecule crosses the boundary. This is clearer if one considers that above the boundary, the molecule represents a total of 6 molecules. When the molecule drops below the boundary, it can only represent 3 molecules, so another molecule must be added to the distribution to represent the other 3 molecules.

The added molecule(s) is (are) assigned the same position and velocity components as the original molecule. While this does not approximate true kinetic theory at first, in practice the positions and velocities are soon randomized

by collision processes.

On the other hand, if a molecule crosses a weighting factor boundary while moving away from the axis of the sample space, there is a probability that it must be dropped from the molecular distribution. The probability is given by

$$1 - \frac{\text{old weighting factor}}{\text{new weighting factor}}$$

The random number generator is used to generate a histogram of disappearing molecules to match the actual probability of disappearance. The whole idea behind using weighting factors is to increase test body - flow field interaction in a given running time. Thus, when body surface quantities are accumulated (like flux, energy, momentum, heat), they are accumulated in terms of the weighted number of molecules striking the body. This is particularly important if the body exists in two or more weighting zones, so that surface quantities in different zones can be correctly compared.

The program is set up to handle a test body consisting of a sequence of connected conic sections rotated about a common axis of symmetry. Some typical surfaces which can be considered without modification are sections of spheres, cones, cylinders, ellipsoids, hyperboloids, and paraboloids. A disc perpendicular to the symmetry axis cannot be handled without modification, since it represents a multi-valued function in r . Cones with very large apex angles ($\sim 180^\circ$) are

used in place of discs. The procedure used to specify body surfaces is to generate a form of the standard equation of the surface in question, substitute into this equation the actual values of the constants, and reduce the standard form to the following type of equation:

$$Ax^2 + By^2 + Cz + D = 0$$

where A, B, C, and D are numerical values

These coefficients are used as input data to the program. The program requires additional parameters for each conic surface. Each surface can be axially divided into several segments for the purpose of accumulating the body surface parameters like flux, heat transfer, etc. The x-coordinate on the right side of the segment is required as data. (The segments are divided azimuthally by the same radial planes that partition the sample-space.) In addition, the temperature, energy accommodation coefficients and tangential accommodation coefficients for each species are required for each segment. The temperature of a surface is normalized by the free-stream temperature.

The tangential accommodation coefficient specifies the fractional part of incoming tangential momentum that is lost to the surface. For the reflection model used in this program, this also represents the fractional part of the molecules colliding with the body surface whose collisions are diffuse. The remainder collide specularly.

The energy accommodation coefficient specifies the ratio

of the net molecular energy flux absorbed by the body to that energy which would be absorbed if all re-emission were appropriate to equilibrium at the surface temperature of the body. This coefficient together with the tangential accommodation coefficient and the surface temperature determines the effective temperature of the diffuse component of the re-emitted molecules in the reflection model used in this program.

It is also possible to collect velocity samples of the colliding molecules on a restricted number of surface segments to generate distributions for the internal flow program. For this purpose, molecules that collide with the body surface are considered in two classes. The class called "UNCOLLIDED" consists of "free-stream" molecules. That is, these molecules have not previously collided with the body surface or with any other molecules other than "free-stream" molecules. The other class, called "COLLIDED", consists of molecules which have previously collided with either the body surface or with other "COLLIDED" molecules.

Appendix A contains the full listing of the computer code in Fortran IV. The main program consists of dimensioning statements coupled to a fairly detailed description of the input cards (using the NAMELIST input). The main operating program is called RUN which in turn calls the appropriate subroutines. Figure 2 shows a schematic of the flow chart for the operation of the Program EXTERNAL.

B. Modelling of the SUMS Problem

1. Geometric Modelling

The computer code described above requires an axisymmetric geometry of the body, although the flow vector can be at an arbitrary direction. The Space Shuttle nose geometry in the vicinity of the SUMS orifice has to be modelled by an equivalent body of revolution. A paraboloid of revolution around an axis inclined at about 8° with respect to the shuttle axis models the lower surface cross sections of the body both in the symmetry plane and in the transverse direction reasonably well. Figure 3 shows a sketch of the actual and modelled Space Shuttle nose geometry. This model will be called the parabola model, and the flow direction of 32° with respect to the axis will be used to represent the 40° angle of attack of the Space Shuttle.

Both for the purpose of benchmarking the test runs against continuum results and to achieve better resolution an alternative axisymmetric flow model was examined. This model consists of a hyperbola rotated about the flow velocity vector passing through the stagnation point. Figure 4 shows a typical cell geometry for this model. This is also the model used by Professor Clark Lewis for his continuum calculations. As will be seen from the results in the next section the latter (axisymmetric) model is questionable as a representation of the flow in the stagnation region of the

Space Shuttle, at least at higher altitudes.

2. Modelling of Cross Sections

In order to avoid uncertainties associated with the choice of hard sphere cross-section to best model the "typical" collision, an energy dependent inverse power law cross-section collision code has been incorporated in the program. The power law exponent and reference cross-section is chosen to provide the best fit to the viscosity temperature dependence over the range between the wall and stagnation conditions which represents the energy range of the important collisions. Figure 5 shows that viscosity can be matched to within a few percent over the relevant range with a single choice of exponent and reference conditions

Using the exponent of $N=.552$ and a reference temperature of 1000°K four axisymmetric cases of a hyperbola at 0° have been run to simulate 87KM, 95KM, 105KM and 115KM altitudes. Since in these cases rotational energy was not included these represent a fictitious monatomic gas of $\gamma=5/3$. Figure 6 shows "smoothed" temperature profiles normalized to the free stream temperature. Note that at the two lower altitudes a relatively "flat" Rankine Hugoniot (R.H.) region exists, while at 105KM R.H. conditions are barely reached and at 115KM the temperature peaks at about 70% of R.H. temperature. The shock layer in all cases is, however, very much thicker than the inviscid result based on nose radius R_N . Note, however, that the hyperbola is very blunt and it is not clear that the nose radius is the appropriate dimension or that this models properly the flow about the Space Shuttle Nose.

3. Rotational Energy Exchange

On the basis of previous experience in modelling rotational energy exchange in Reference 4 the model of Larsen and Borgnakke (9) was chosen as appropriate for the blunt geometry under investigation. The External Flow Program was therefore re-coded to include an arbitrary number of internal energy modes but with a single relaxation time related to the parameter ϕ which ranges from $\phi=0$ representing no exchange to $\phi=1$ simulating maximum available exchange at each collision. An indication of the effect at 115KM is shown in Figure 7 and 8 for the most rapid energy exchange ($\phi=1$). Note that, as expected, the peak temperature is lower for the $\gamma=1.4$ case because some of the energy goes into rotation. Also note that significant non-equilibrium exists between translation and rotation even at this maximum energy exchange rate corresponding to $\phi=1$. The effect on surface properties shown in Figure 8 is virtually non-existent for the shear, but noticeable on the heat flux and pressure.

4. Preliminary Attempts at Comparison with Continuum Results

Some preliminary comparisons of both the monatomic ($\gamma = 1.667$) and diatomic ($\gamma = 1.4$) runs for the axisymmetric hyperbola were made with continuum results provided by Prof. C.H. Lewis based on continuum theory^(10,11). Among the surface properties only the pressure distribution agreed reasonably well. Since this is the property least sensitive to detailed flow field behavior the agreement is not a very sensitive test. A typical comparison is shown on Figure 8. A comparison of heat flux is not shown on this figure as it is not the same scale at this high altitude. The continuum result of Prof. Lewis gives the stagnation point heat transfer coefficient $C_H = Q/1/2\rho U^3 = 2.8$ which is physically unrealistic. At 105KM the continuum heat transfer coefficient C_H is below one but still appreciably above the Monte Carlo results. At 95KM the Monte Carlo results give heat transfer that is almost twice as large as Prof. Lewis's result for $\gamma = 1.4$.

The most significant discrepancy between the continuum and Monte Carlo results arises in the shock layer thickness. While the definition is somewhat arbitrary, a comparison of the subsonic region in the vicinity of the stagnation point does give a good indication of the extent of influence of the "downstream" portions of the body. Figure 9 shows the Monte Carlo results for the $M=1$ line at 115KM together with the shock line from C.H. Lewis. This discrepancy led to a whole re-examination of the modelling of the nose geometry.

5. Re-examination of Geometric Modelling

Since the subsonic region appears to extend to the "shoulder" wherever the hyperbola is terminated as shown in Figure 9, the applicability of the hyperbola model becomes suspect, at least for the Monte Carlo calculation where upstream influence cannot be eliminated. In order to further assess the effect of the geometric model on the results a comparison of the hyperbola at 0° to the parabola at 32° are presented in Figures 10 and 11. It can be seen from Figure 10 that the subsonic layer in the vicinity of the SUMS orifice is significantly different in the two cases. Figure 11 shows that while the pressure is not dramatically affected by the model the heat flux and shear are significantly altered. On the basis of these results it was concluded that while the runs for the hyperbola may indicate trends they are neither representative of the SUMS region on the Space Shuttle nor good candidates for benchmarking with the

continuum results. All subsequent runs were therefore made using the geometric model of a parabola at 32° angle of attack.

6. Inelastic vibrational Excitation and Dissociations

Because of the high energies of collision between free-stream and reflected molecules, inelastic collisions (certainly vibration and dissociation) are in principle possible. At the highest altitude the number of such collisions is expected to be insignificant because the cross-sections are low enough so that a typical molecule will reach the body with a negligible probability of a previous vibrationally exciting or dissociating collision. At the lower altitudes 87 to 95Km the number of collisions suffered by a typical molecule before reaching the body surface is measured in the tens or hundreds and therefore inelastic cross-sections that are only a few hundredths of the elastic and rotational ones may produce significant effects. The detail needed to properly model these collisions in the Monte Carlo Programs far exceeds the available experimental information, which primarily gives overall rate constants. An investigation of the best combination of analytical and empirical information was initiated early within the grant period. A theoretical attempt to couple low energy vibrational excitation experimental information to the highly non-equilibrium high energy collisions through theoretical work was initiated. Appendix B contains a Master's thesis

presenting the formalism and giving the initial results of a theoretical formalism. Based on those results, coupled to some limited experimental data, a method for determining the probabilities of specific outcomes in individual collisions in the Monte Carlo Programs can be developed. The increase in computing times, however, may make the feasibility of using such a code, for anything but benchmarking, prohibitive.

C. MAJOR RESULTS FOR EXTERNAL FLOW

The primary effort during the grant period up to February 1981 was spent in developing the code for external flow and establishing the appropriate geometry to model the region of the Space Shuttle in the vicinity of the SUMS experiment. The major results to that date are presented in the renewal proposal for the period February 1, 1981 to July 30, 1982 which also served as a progress report on the previous grant period. (12) We will only summarize those results here and update them on the basis of additional external flow computations.

A representation of the shuttle nose geometry as a paraboloid of revolution around an axis 8° from the actual shuttle axis (Figure 3) was found to adequately model the windward side of the shuttle in the vicinity of the SUMS orifice. Computations for both this model at an effective angle of attack of 32° and an alternative axisymmetric flow about hyperboloid model centered on an axis through the nominal stagnation point, demonstrated that the paraboloid model at angle of attack is necessary to adequately model conditions near the SUMS orifice. The typical body and computation cell geometry is shown in Figure 12.

A reasonable indication of the flowfield can be obtained by examining density contours or Mach number contours about the body. Figure 13 shows the sonic lines and the $M = 5$ lines indicating approximately the outer extent of the "shock" layer at 95 Km and 115 Km, within the plane formed by

the velocity vector and the axis of the modelled paraboloid. Note the greater "shock" layer thickness at the higher altitude. Also note the fact that even at 95 Km, the shock layer is a large fraction of the local body dimension such as nose radius. The "shock" thickness is a major portion of the entire "shock layer" casting doubt on any calculation incorporating a thin "shock" assumption. Figure 14 shows some density contours on the "windward" side of the body at 95 Km altitude as well as the sonic line. Note the rather constant density rise towards the body with no discernible separation between "shock" and "boundary layer." Also shown are estimates of the stagnation streamline and another streamline along which the velocity only goes slightly subsonic. Note the rather gradual turning along the latter streamline and the rather diffuse nature of the shock layer even at this altitude where the nominal free stream Knudsen number is 0.04. In order to give a better picture of the three-dimensional aspect of the flowfield, a sketch of one-half of the paraboloid and some contour plots of $M = 1$. and $M = 5$. lines are shown on Figure 15.

Information in Figures 13 through 15 gives some indication of the nature of the flowfield in the vicinity of the SUMS orifice. The ultimate objective, however, is to establish properties at the vehicle surface. Figure 16 shows the variation of two surface fluxes (normal pressure and heat flux) at 95 Km along the four cross-planes shown in Figure 15. The normal pressure p_s normalized by $1/2\rho U^2$ ($C_p = p_s / 1/2\rho U^2$) is shown on the right side of the figure while a heat

transfer coefficient ($C_H = Q/1/2\rho U^3$) is shown on the left. Figure 17 shows the variation versus angle in the cross-plane nearest the one containing the SUMS orifice, and compares the results to some theoretical and semi-empirical predictions. The pressure coefficient, as expected, lies approximately between the free molecular (C_p (F.M.)) value and the modified Newtonian (C_p (NEWT)) value, up to about 90° . (Note that at 90° around the axis the local angle (β) between the velocity vector and the surface normal is approximately 64° at this cross-plane.) The heat transfer coefficient lies substantially below the free molecular value. It is also compared to a semi-empirical extrapolation of experimental results of stagnation heat transfer on hemispheres presented in reference 13. Direct comparison is clearly dubious due to the substantially different body geometry and the implicit assumption of totally local behavior contained in the $\cos\beta$ variation with local angle of the surface to the velocity vector. Some additional degree of uncertainty is contained in the choice of "body size" that is used to evaluate the correlation parameter K_R^2 .

The ultimate objective of the external flow calculations is to provide information on the properties of the gas entering the SUMS orifice. It is the potential non-equilibrium nature of the entering distribution of molecules that is responsible for the difference between the surface "pressure" and the gas pressure within the internal plumbing around the mass spectrometer. The non-equilibrium aspect is most commonly represented by a temperature jump and a

velocity slip within the continuum formalism. Under the highly rarefied conditions at the upper end of the altitude range of interest, even that description may be inadequate because of the highly non-Boltzmann distribution of the molecules arriving at the surface. Figure 18 shows the distribution function of the flux of molecules arriving at the SUMS orifice at 95 Km and 115 Km obtained from the external program. The flux distribution is plotted both versus molecular velocities normal to the local surface and tangential within the plane formed by the free stream velocity vector and the local normal. Note that at 95 Km one could fit the distribution by a Maxwellian with some temperature different than the body temperature and a slip velocity which is comparable to the free stream mean molecular speed. At 115 Km the distribution is clearly composed of two components, the free stream molecules arriving unperturbed at the surface, and the collided molecules having a broader distribution possibly representable by another Maxwellian. Clearly this potential bimodal character of the incoming molecules must be recognized in the evaluation of the "entrance" problem at the SUMS entrance.

As a companion to the initial version of the external code (EXT) an internal code (INT) was developed (8). The objective of this code was to determine local properties and surface fluxes in a cavity connected by a tube to the exterior surface. The input molecular distribution is obtained from the surface flux information provided by the external code. While the code is in principle general to allow a wide variety of geometric configurations it is optimized for a short tube-large cavity geometry. Early attempts to apply this code directly to the SUMS inlet geometry resulted in very long running times with little assurance that a steady state solution had been achieved. A total recoding of the entrance problem was therefore implemented resulting in a code (INTERNAL) that allows greater flexibility in handling the long tube geometry of SUMS as well as incorporates all of the changes in intermolecular collisions, numbers of species and rotational energy exchange developed for the code EXTERNAL.

A. Entrance Computer Code INTERNAL

The purpose of the program described in this section is to determine the fluid field inside a cavity which consists of a connected sequence of conic sections rotated about an axis of symmetry. The cavity is considered in two parts: a main cavity on whose interior surface the sensor will be

located, and an inlet tube whose orifice is presumed to be at the exterior surface of a spacecraft. A detailed code listing is given in Appendix C.

The input data for this program includes a molecular distribution function which is obtained from an external-flow run. Some of the other data refers to input parameters of the external-flow run. Thus it is seen that the pair of programs can be run as a set, computing the conditions inside a cavity which exist for a given set of conditions in the undisturbed free-stream flow. The programs have not yet been directly coupled, although they are written with this intent.

The program was constructed by turning an external-flow program inside out. In doing this, the basic molecule/body collision mechanism is preserved with only minor changes, while the molecule/molecule collision mechanism are not changed in any way.

A general description of the way in which the internal-flow program conducts the experiment is unnecessary since the description in the preceding section generally applies here. Any important differences will be described as they occur.

The numerical experiment takes place in a cylindrical block of space quite similar to that used in external flow. However, the inlet region must be cylindrical in shape while the cavity region can be defined by conic sections.

There is no subdivision of 1st level cells into smaller sized cells. Also, the number of azimuthal wedges is specified for a full 360° (since there is no plane of symmetry determined by an input flow direction) and is the

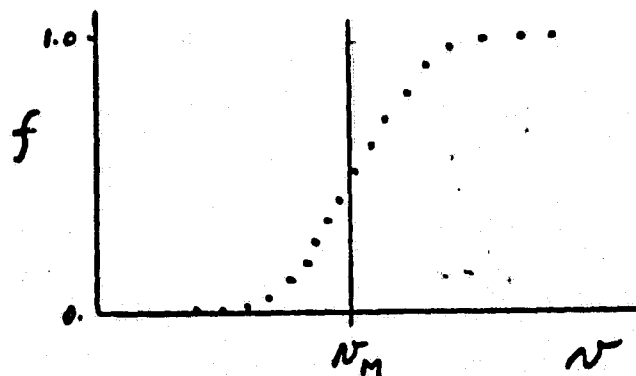
same for both regions. No weighting factors are used as the desired surface fluxes on the walls generally occur near the outer edges of the sample space.

The orientation of the sample space is determined by the geometry of the particular body surface segment in external flow which is used as the orifice area of internal flow. The body surface normal of external flow becomes the x-axis (axis of symmetry) of internal flow; the direction given by the cross product of the x-axis of external flow and the body surface normal becomes the z-axis of internal flow; and the direction given by the cross product of the above cross product and the body surface normal becomes the y-axis.

The zero radial plane is the positive x-y plane, and azimuth angles range from 0° to 360° , with the positive z-axis at 90° .

The sample space is initially populated with gas molecules in a manner similar to that used in external flow. In this program, however, the gas molecules are initially in thermal with the walls locally. The selection of the density profile initially in the tube indeed poses a problem and is very critical in minimizing running time. The code therefore allows an arbitrary initial distribution specified by the user. The choices and procedures for selecting them are described in the next section. The key objective is to provide a distribution that is consistent with the input flux and will not have to be severely altered in magnitude to arrive at the steady state.

The molecular input distributions are generated from data produced by an external-flow run. This data consists of velocity samples of molecules impinging on the external body surface segment which is considered to be the orifice area of the inlet region. The form of the velocity samples is a series of horizontal S-shaped curves, one for each molecule type, in each of three directions, for both UNCOLLIDED (free-stream) and COLLIDED molecules. Each s-curve gives the fraction of molecules (impinging on the body surface segment) with a velocity \leq the given velocity. A typical curve for UNCOLLIDED molecules is given [v_M is the velocity corresponding to the center velocity of the distribution as computed in the external flow program]:



The information from these curves is used directly to generate the molecular distribution of the molecules entering

the inlet.

The population (pressure) inside the cavity region is selected initially as an input. The computer simulation proceeds until a steady profile inside the tube entrance region is generated. Since this does not necessarily require zero net flux, a series of runs with different cavity pressures is generated and a cross plot of flux versus pressure is the output for a single external flow input condition. The equilibrium solution if desired can be obtained from the zero flux point.

B. Scope of "Entrance Problem" for SUMS experiment

Within the free molecular regime the "entrance problem" has been examined by Hughes and deLeeuw⁽²⁾. In order to indicate the potential magnitude of the problem, Figure 20 shows the variation of both the surface pressure (p_s) and the chamber pressure (p_c) versus angle of attack under free molecular conditions at infinite speed ratio and a very long tube (conditions approached at high altitude during Space Shuttle re-entry). For the SUMS location, the local angle of attack β is approximately 28° giving a possible difference of a factor of 10 between the surface pressure (p_s) and the internal chamber (p_c) that directly affects the mass spectrometer reading. While this theoretical result is expected to be accurate at 140 Km and above, at lower altitudes the local flux distribution will have both a directed and a rather diffuse component (see Figure 18), and

also the effect of internal collisions within the tube will begin to play a role (see Figure 1).

Because the tube length to diameter ratio is very large (37) the time constant and therefore the computing time to reach equilibrium is very long. For this reason no attempt is made to simulate the problem all the way to the condition where the chamber pressure behind the tube is at the correct value to nearly balance the net flux. Instead a series of runs with different assumed "cavity" pressures are performed and the zero net flux (or a given small value for the dynamic condition) can be selected to interpolate the correct "cavity" pressure. This procedure if it proves generally successful can of course be automated within the code.

C. Preliminary Results of INTERNAL Code

Preliminary results for the SUMS entrance tube geometry are presented as a couple points on Figure 21 giving to the ratio p_c/p_s versus altitude.

Also shown are some theoretical curves for full continuum (Continuum), for continuum external flow but free molecular flow through the tube (No slip, $Kn_s \rightarrow \infty$), and for fully free molecular flow (Free Molecule). In addition, some points using a best fit Maxwellian to the external flow flux and free molecular internal flow from reference 2, are shown at 95 Km and 115 Km. These results are preliminary, but they do indicate the magnitude of the correction and the fact that no simple use of the currently available results can cover

the entire range of altitudes of interest.

IV. DISCUSSION

If we couple information from INTERNAL (Figure 21) to the results of the external flow computations (summarized as a plot of $p_s/1/2\rho U^2$ versus altitude in Figure 22), we can produce a preliminary estimate of the data reduction curve that could be coupled to the calibration of the mass spectrometer to deduce the $q = 1/2\rho U^2$ during Space Shuttle re-entry. Figure 23 is a cross-plot of the q/p_c versus p_c obtained from Figures 21 and 22 with the 87 and 105 Km results only estimated on the basis of interpolation of Figure 21. The establishment of such a data reduction curve for the nominal re-entry conditions, together with associated error bars as well sensitivities to wall temperature and angle of attack is necessary for the proper interpretation of data to be obtained by the SUMS instrument.

The data measured by the mass spectrometer in the SUMS experiment essentially provides collector currents of charged species of different masses. Calibration of the instrument can relate these to the overall pressure and composition at the entrance to the instrument being calibrated. Since the actual environment of the flux of molecules to be encountered under flight conditions cannot be simulated the calibration is performed with the incident flux essentially in equilibrium with the instrument outer walls (room temperature and no flow). The data reduction procedure must therefore relate the effective environment in the ground test simulation to the desired dynamic pressure q under the

flowing non-equilibrium condition and through the calibration to the instrument measurements.

As discussed above, capability now exists for calculating the surface flux distribution of molecules entering an opening at the SUMS location. The computational procedures for connecting that information to the pressure immediately behind the entrance tube has also been developed, although not fully exercised over the entire range of parameters applicable to the SUMS experiment. That pressure, in turn, can be directly related to the ground test environment used to produce the calibration curves for the mass spectrometer.

The primary objective of future work must be to provide a data reduction procedure that relates the spectrometer reading to the free stream dynamic pressure ($q = 1/2 \rho U^2$). With currently available procedures a relation between the pressure p_c and q such as the preliminary one shown on Figure 23 can be obtained using the best available information on the flight parameters such as velocity, angle of attack, tile temperature, surface conditions, molecular collision parameters, etc. This relation must then be combined with the calibration curve where the instrument readings are related either to p_c directly, or to a calibration pressure which can be related to p_c by conventional means. A single plot of q versus total measured collector current can thus be obtained from the combination of these results.

The relation between the dynamic pressure q and the calibration pressure p_c depends on many parameters of the

problem. Some of these, such as flight velocity, angle of attack, and tile surface temperature are expected to vary only slightly from their nominal values. Since the actual measured values of these quantities will be available on each individual flight, corrections to the data reduction relation should be evaluated in the form of sensitivity coefficients for small changes from the nominal. Studies to determine the effects of these parameters are necessary to establish the significant sensitivity coefficients that must be incorporated into the data reduction scheme.

The parameters such as surface accommodation, surface recombination, and free stream composition can also affect the results. In addition, modeling simplifications of both the geometry of the problem and the molecular collision phenomena can alter the quantitative value of the relation between q and p_c . Because of the unavailability of any in-flight measurements that could lead to an evaluation of these parameters, bounds on the uncertainties they produce should be studied. Sensitivity of the data reduction relation to the most significant of these can then produce bounds on uncertainties on the dynamic pressure q , due to reasonable variations. The final goal of future work is an algorithm for the evaluation of the dynamic pressure q , from the calibration pressure p_c , together with error bounds, due to the uncertainties associated with the external flow and the entrance problems.

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FIGURE 1. RELEVANT KNUDSEN NUMBERS VERSUS ALTITUDE

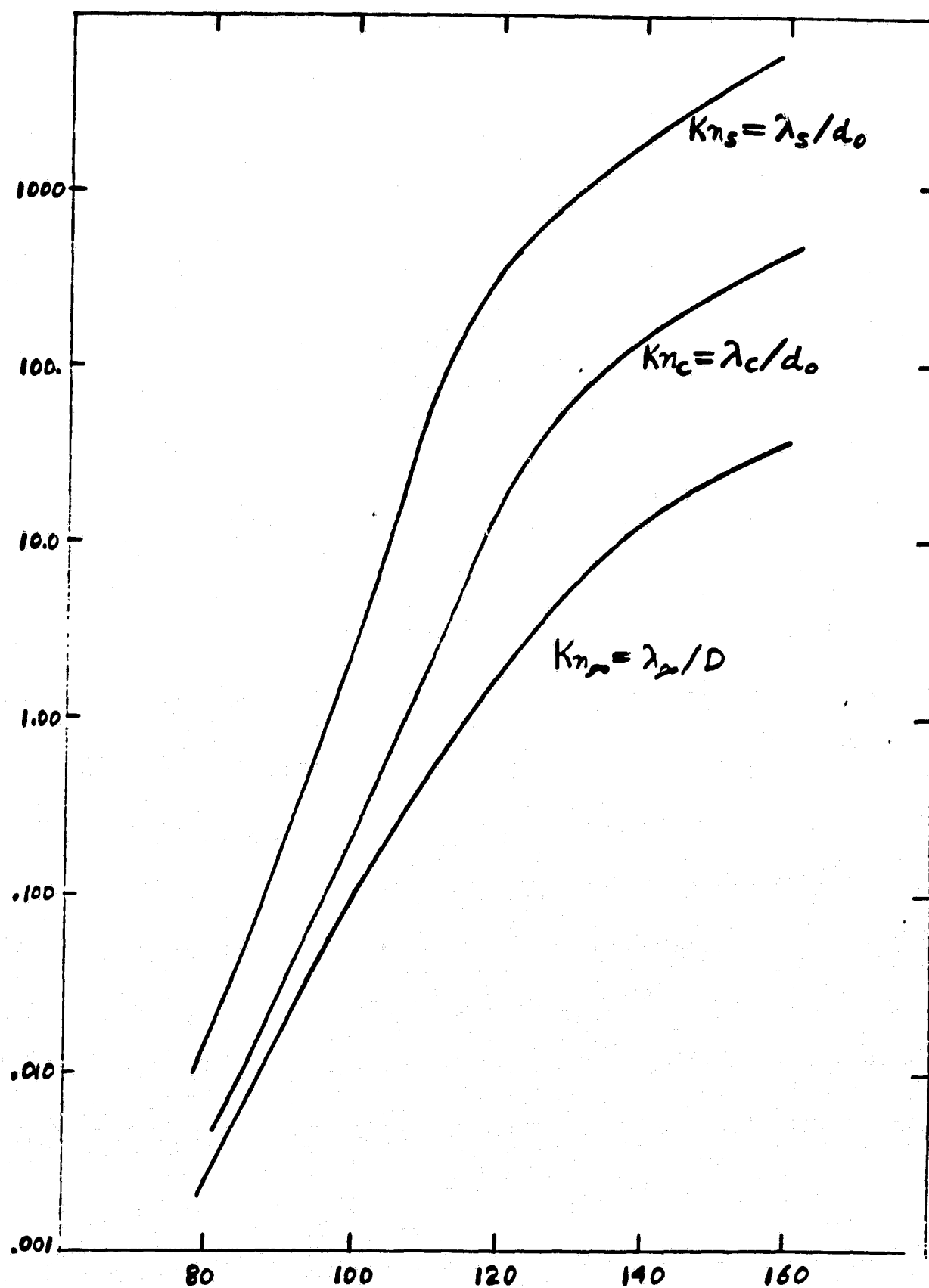


FIGURE 2. SYSTEM FLOW-CHART

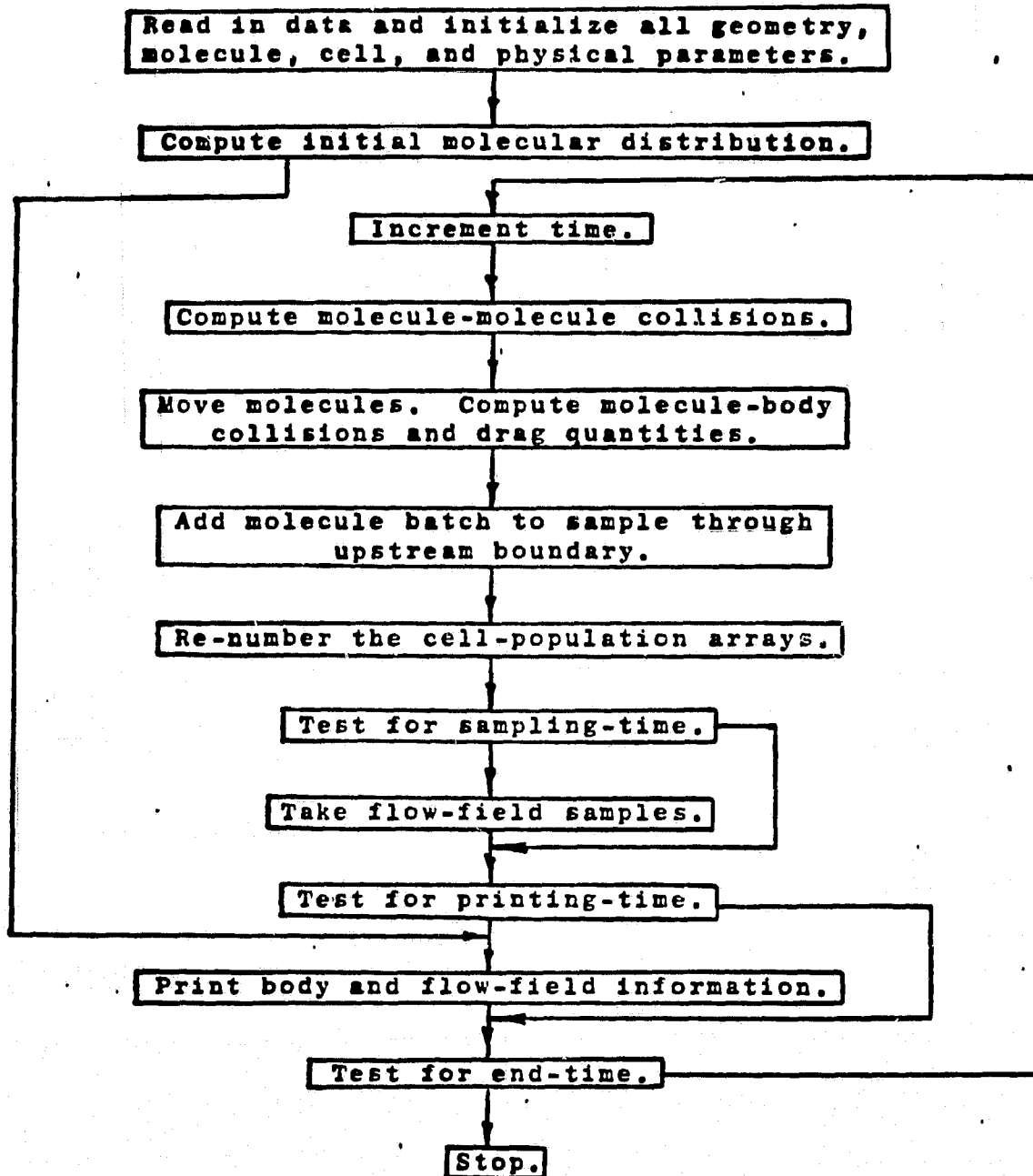
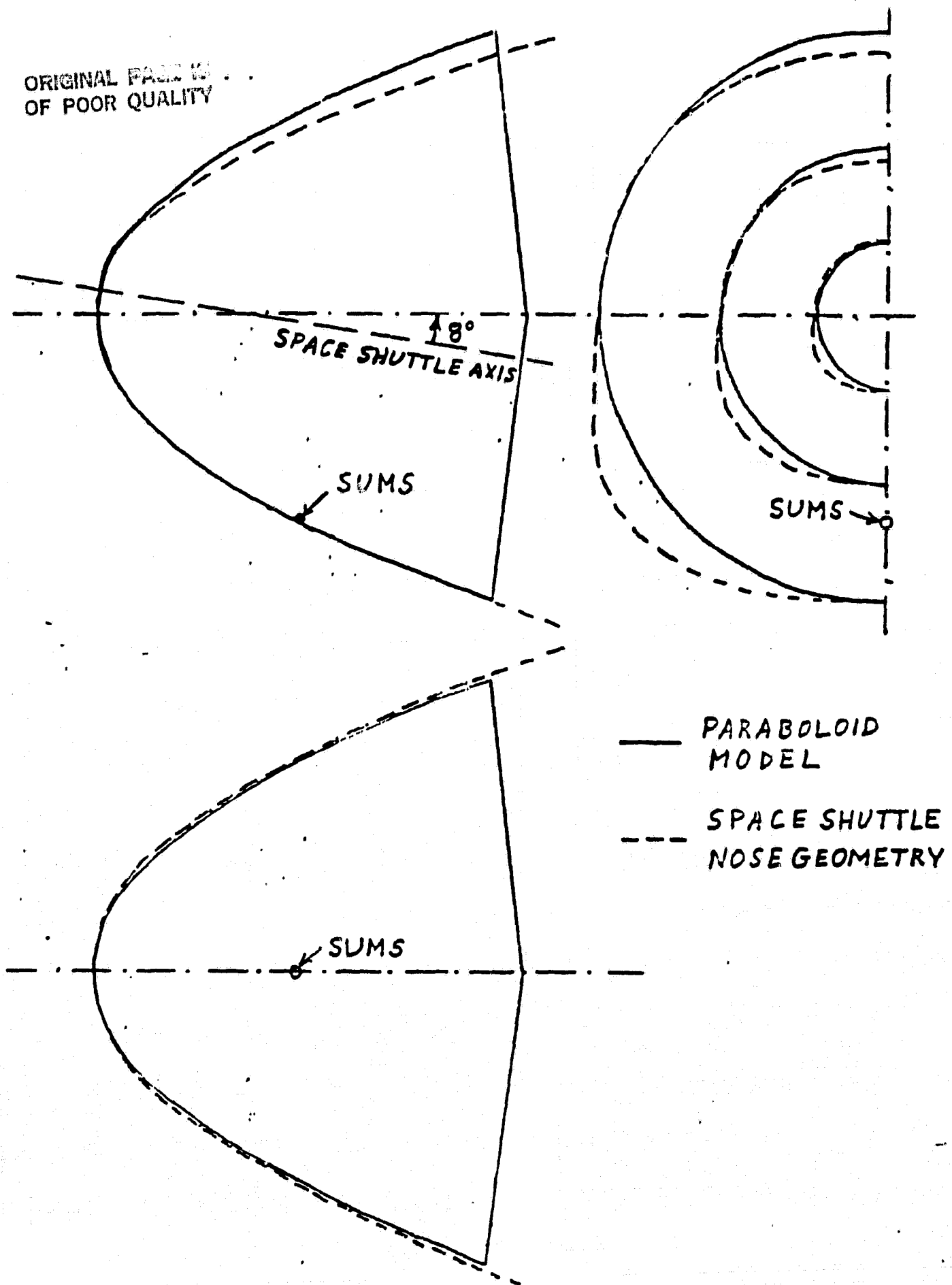


FIGURE 3. MODELING OF SPACE SHUTTLE NOSE (PARABOLA)



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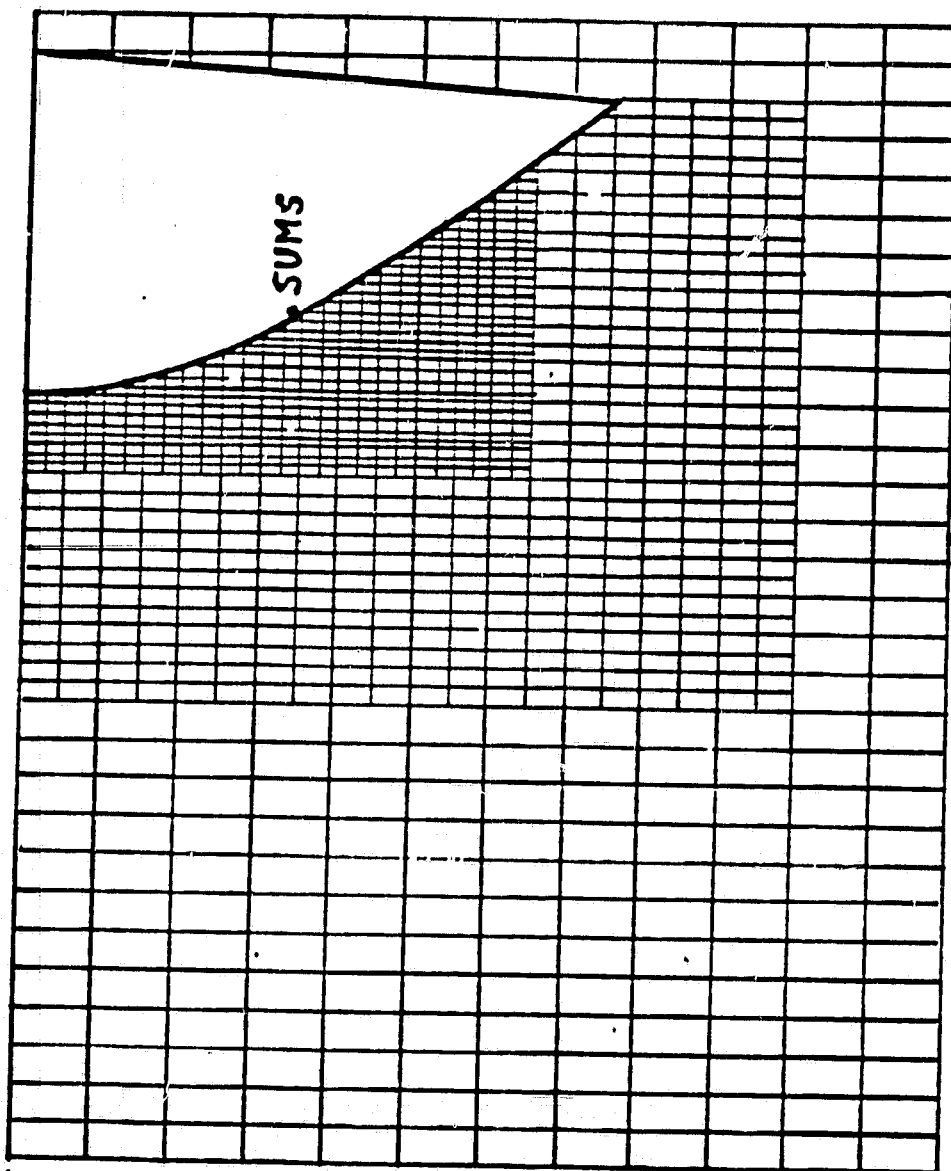


FIGURE 4. TYPICAL CELL CONFIGURATION FOR HYPERBOLA

FIGURE 5. MODELLING OF VISCOSITY-TEMPERATURE DEPENDENCE

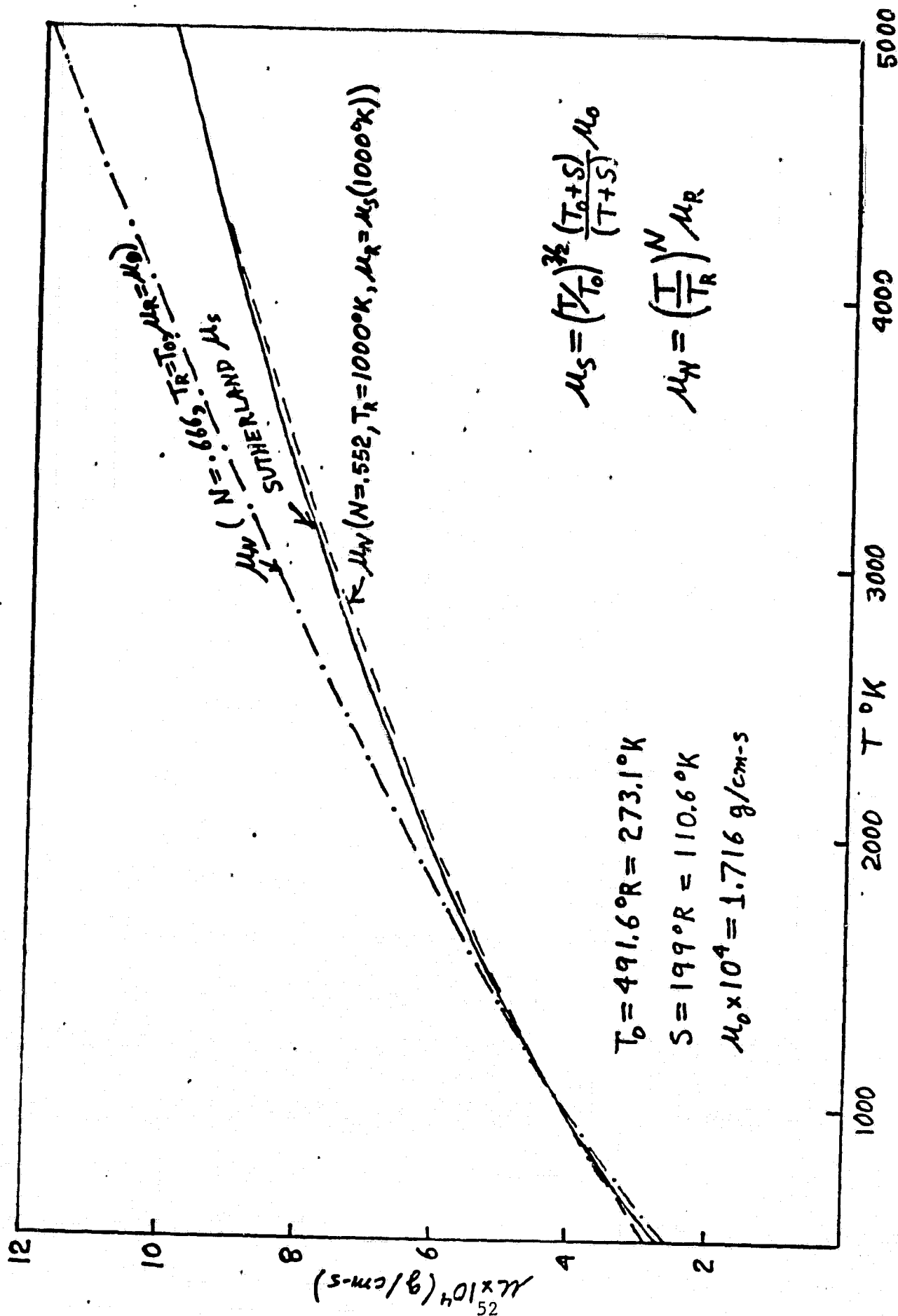
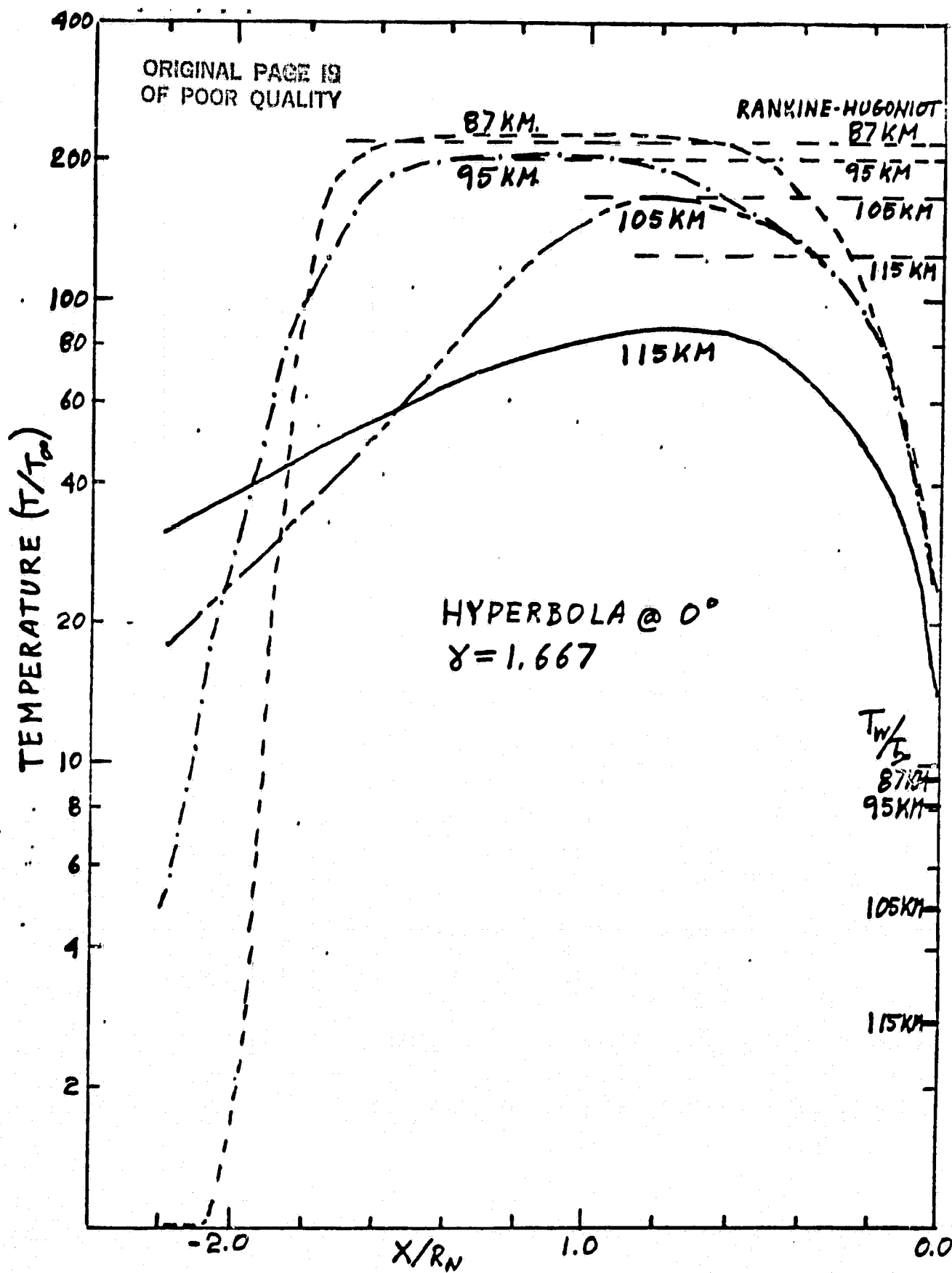


FIGURE 6. TEMPERATURE DISTRIBUTIONS (EFFECT OF ALTITUDE)



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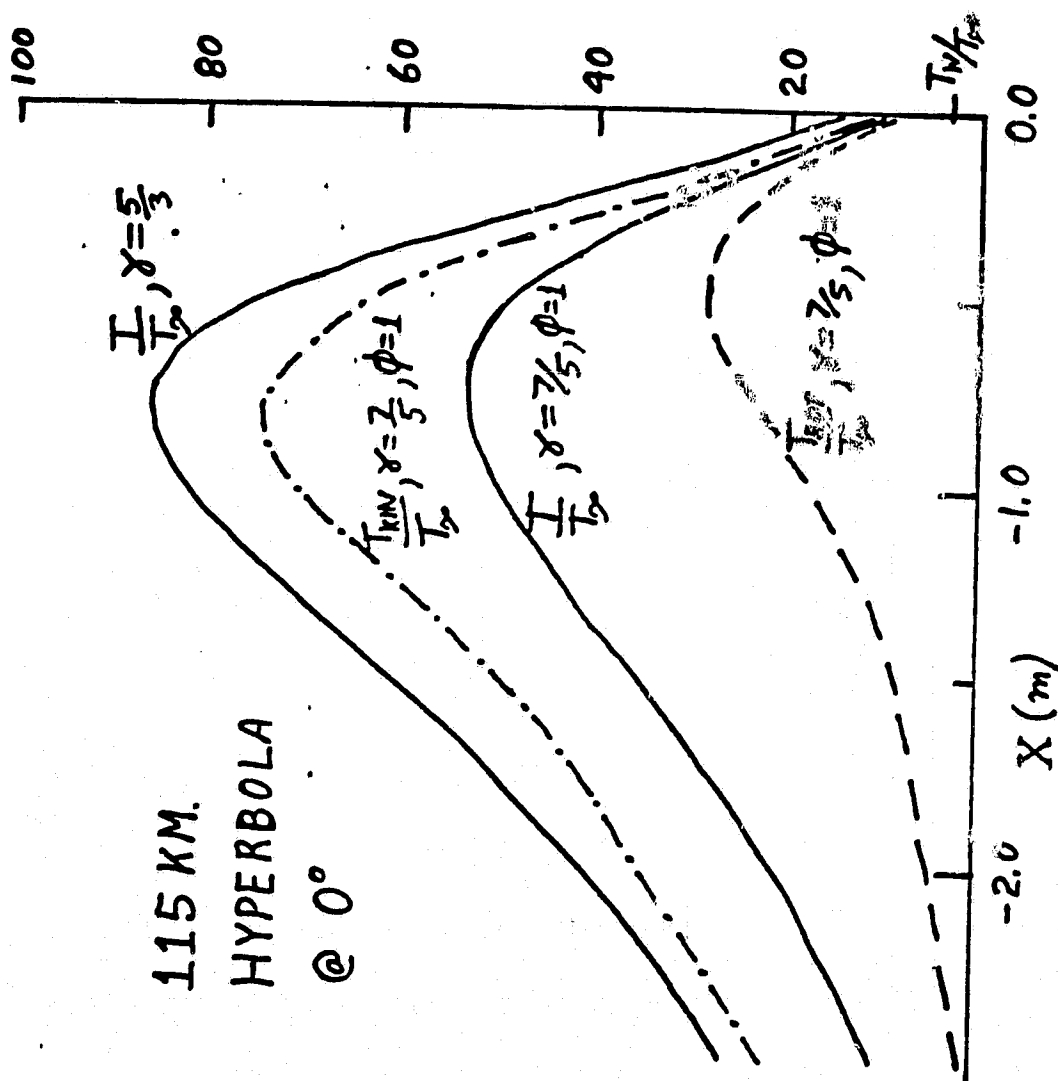
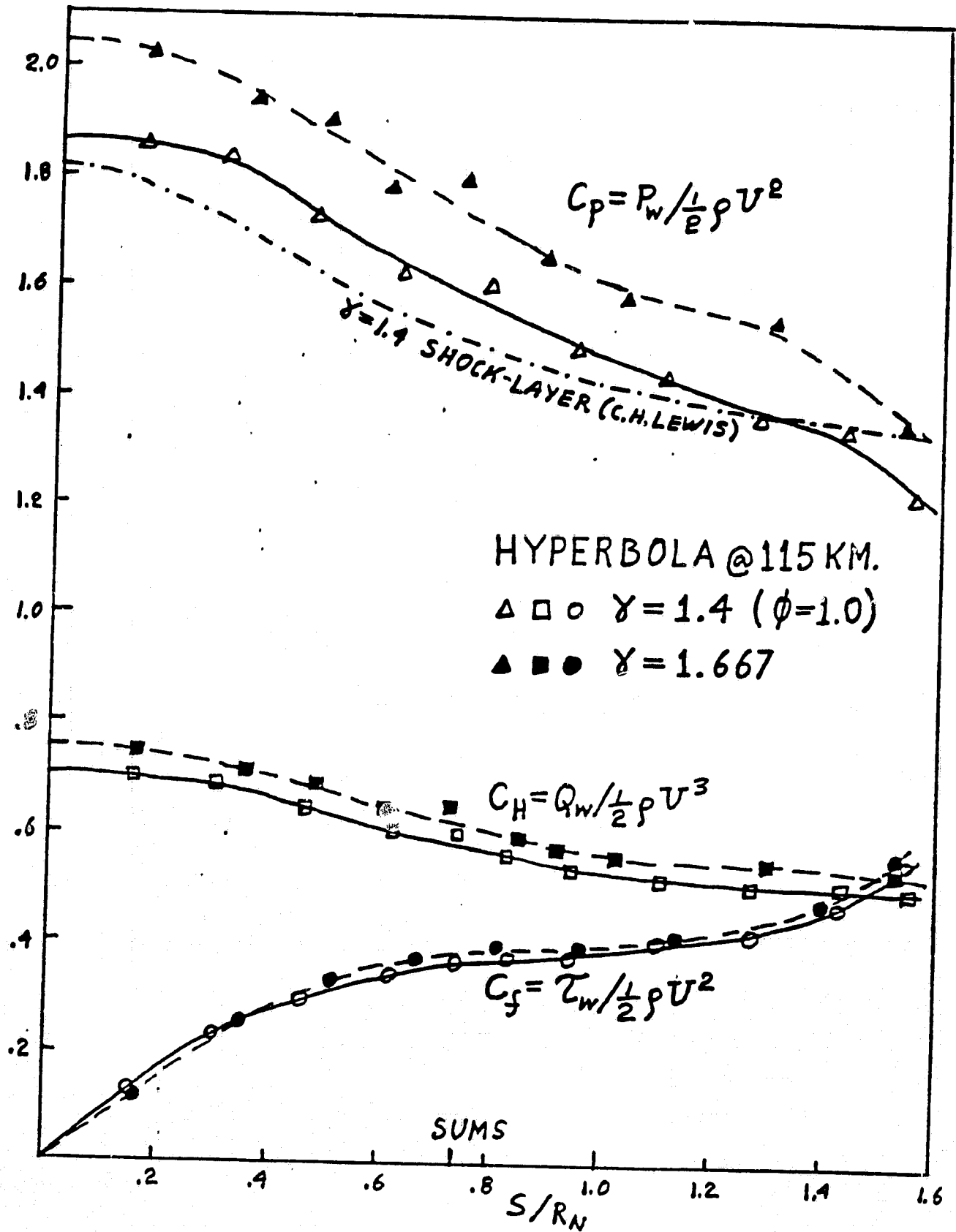


FIGURE 7. TEMPERATURE DISTRIBUTION ALONG SEPARATION STREAMLINE
(EFFECT OF γ)

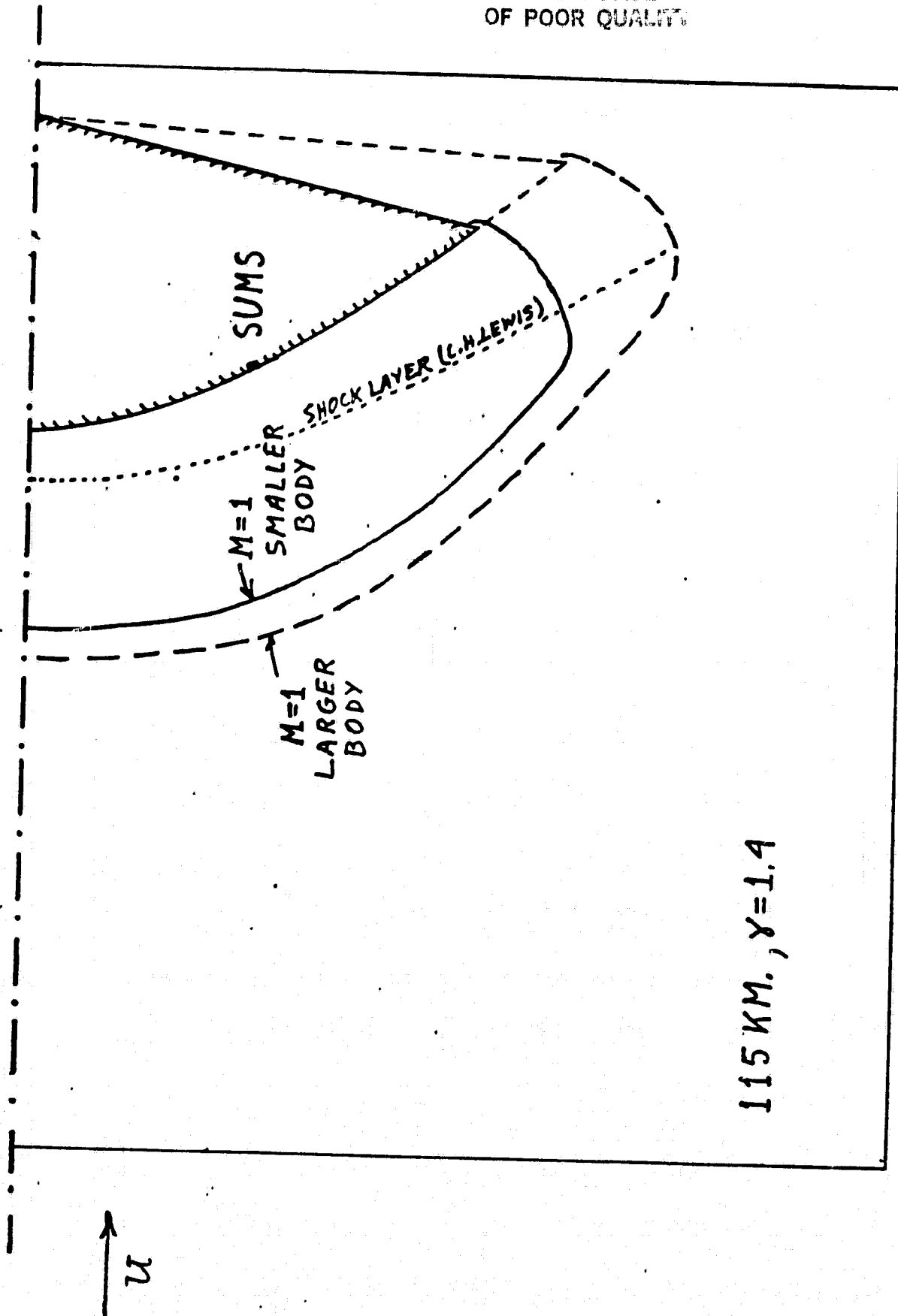
FIGURE 8. SURFACE FLUXES (EFFECT OF γ)



Normalized Distance from Stagnation Point

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FIGURE 9. SONIC LINES (EFFECT OF BODY SIZE)



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FIGURE 10. SONIC LINES (EFFECT OF GEOMETRIC MODEL)

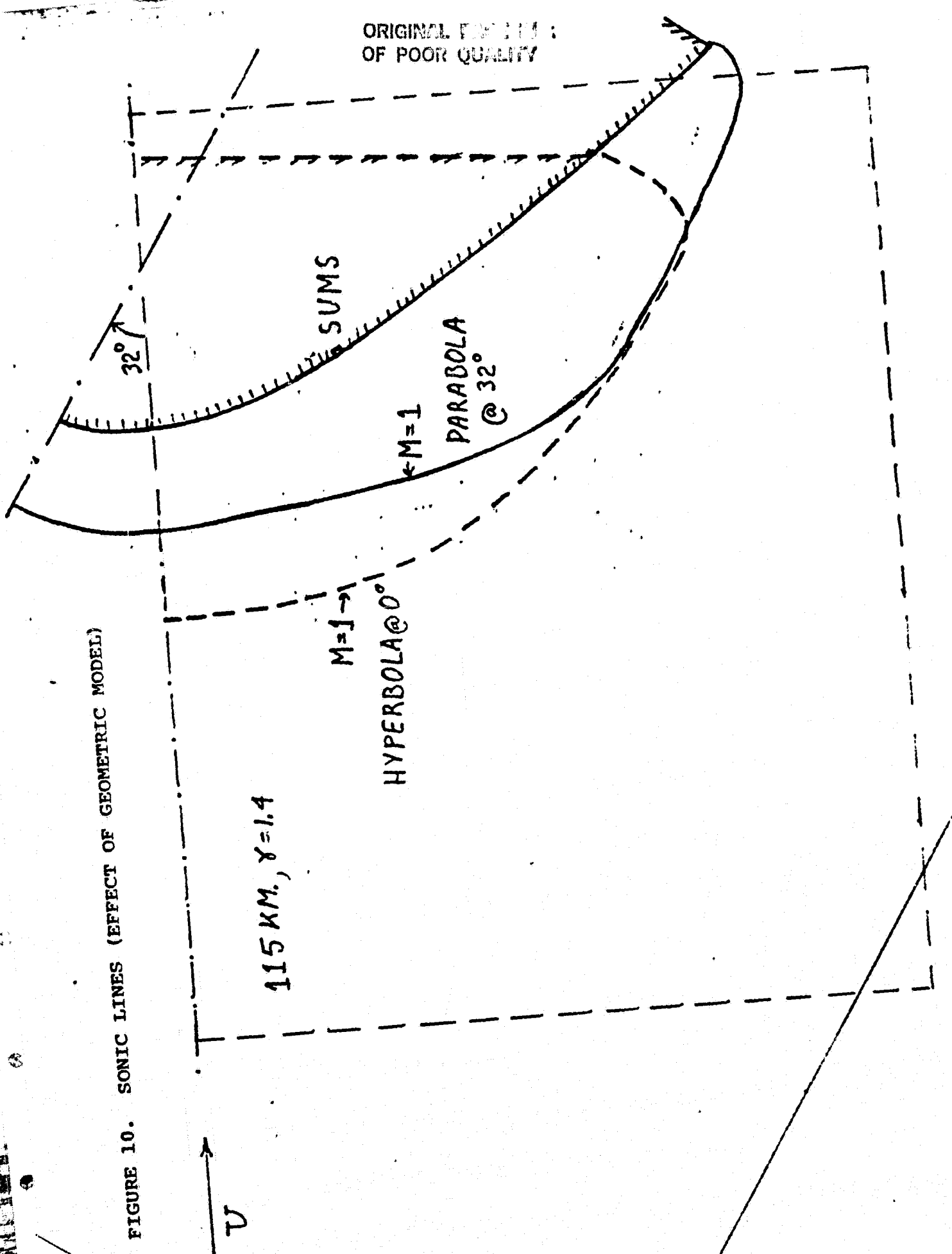
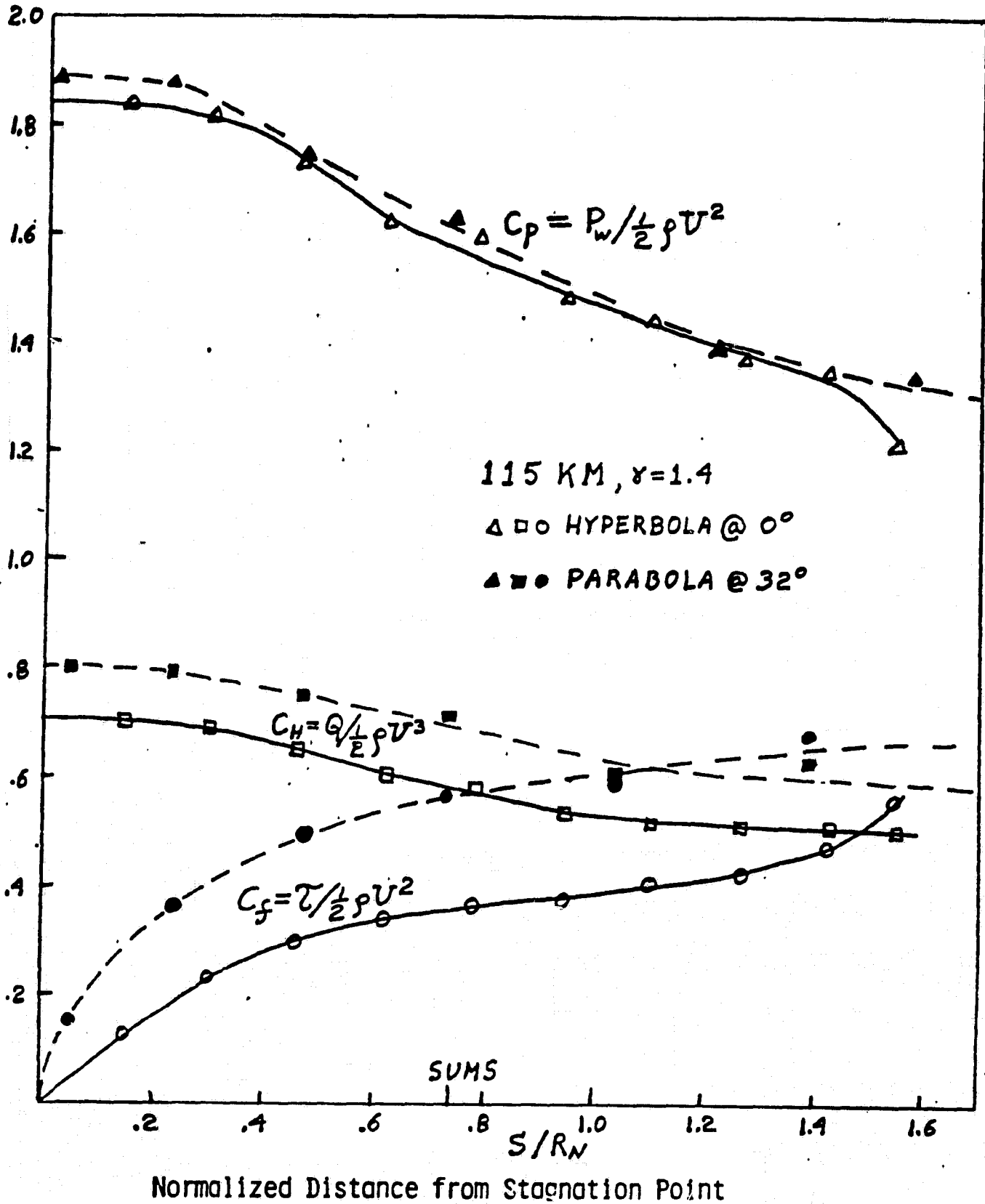


FIGURE 11. SURFACE FLUXES (EFFECT OF GEOMETRIC MODEL)



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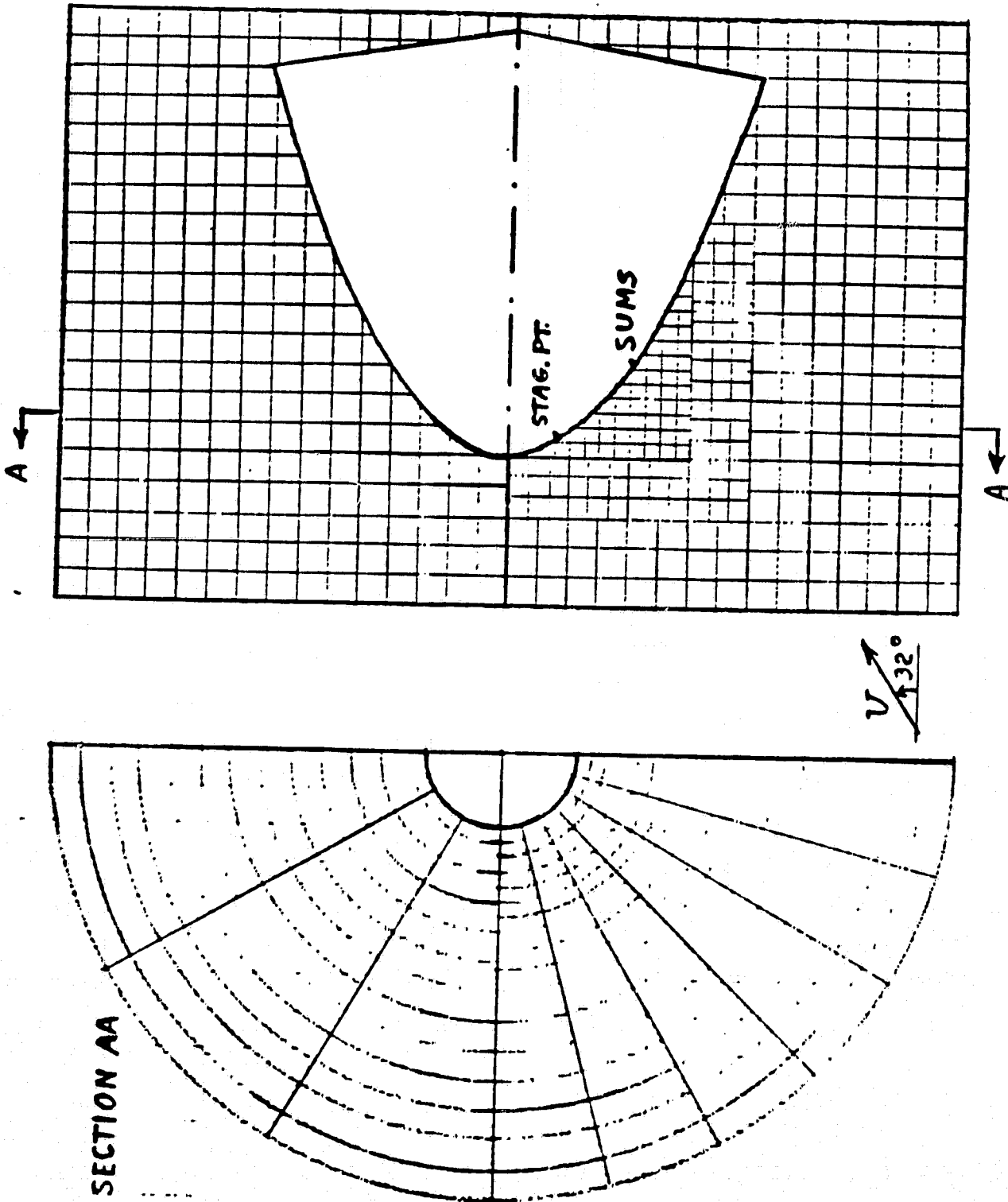


FIGURE 12. TYPICAL CELL CONFIGURATION (95 km run)

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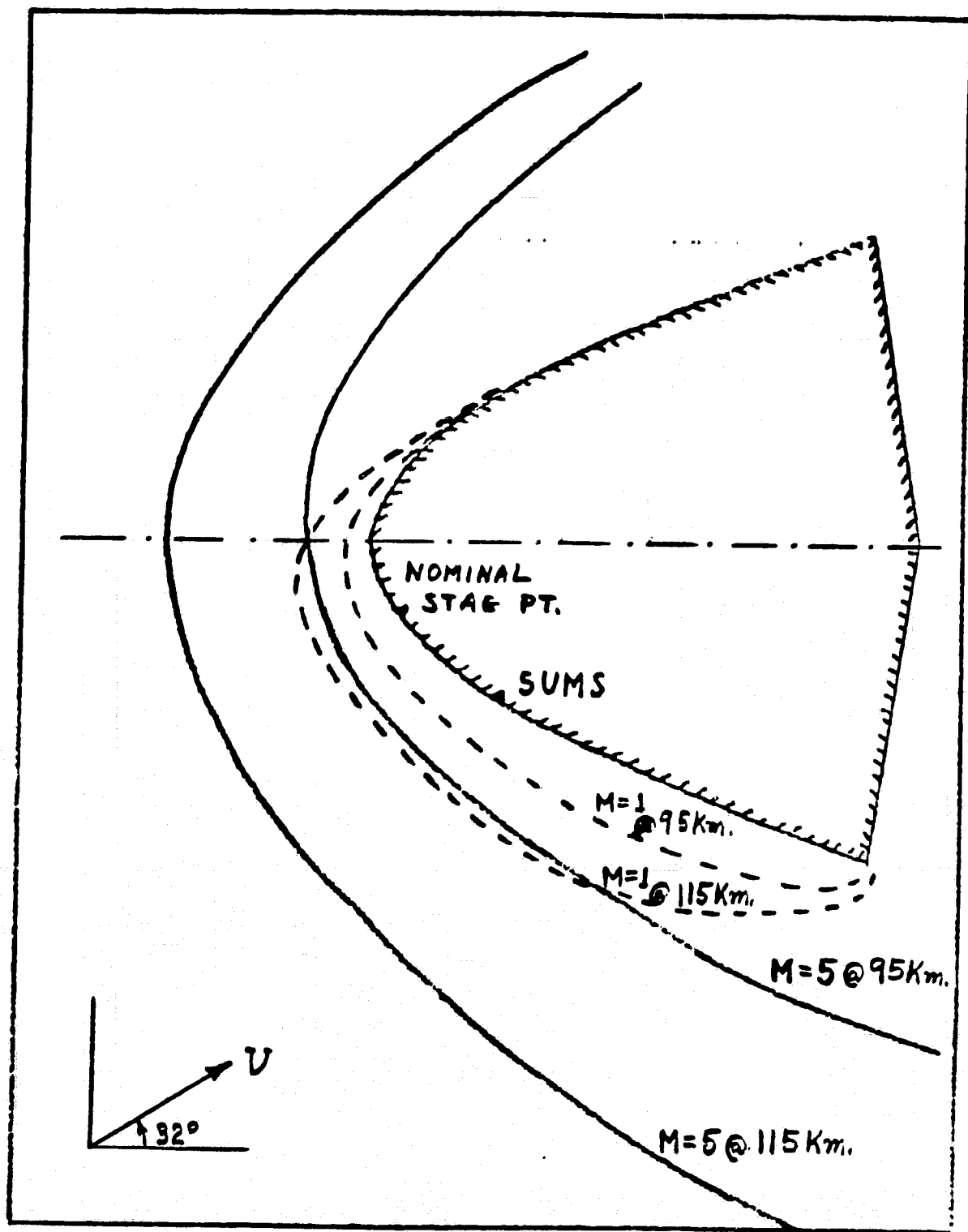


FIGURE 13. MACH NUMBER CONTOURS IN SYMMETRY PLANE AT 95 Km and 115 Km

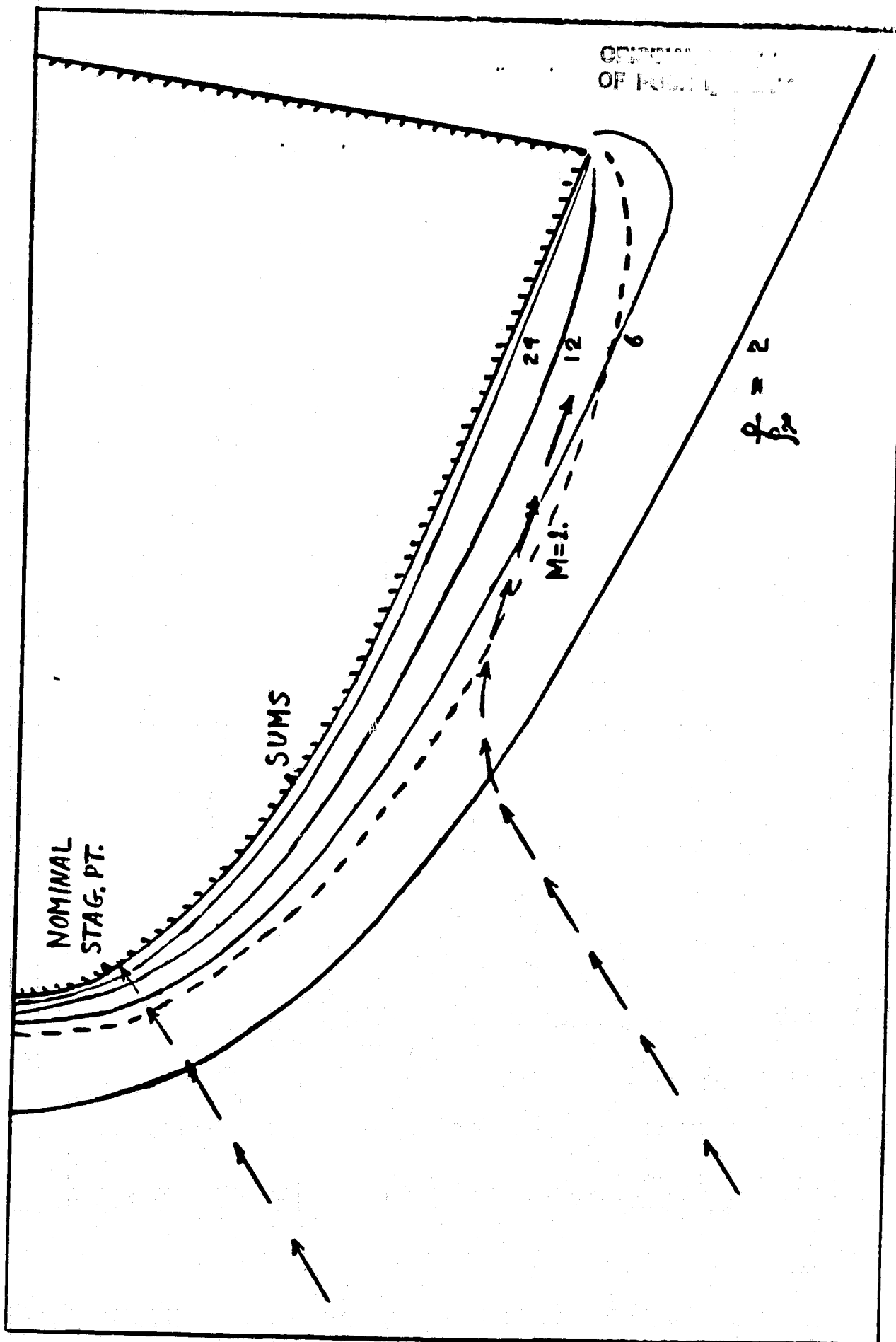
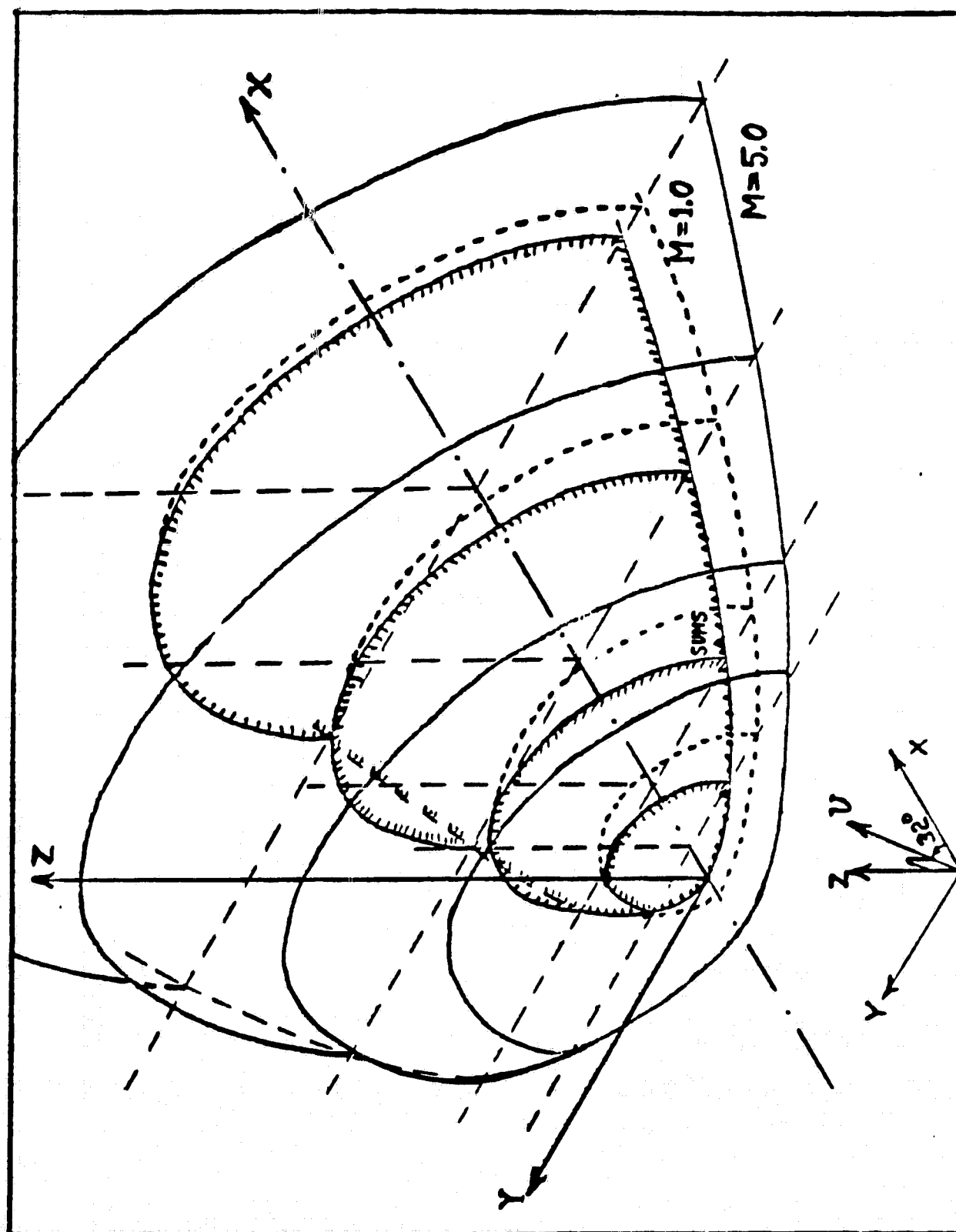


FIGURE 14. DENSITY CONTOURS ON WINDWARD SIDE AT 95 Km ALTITUDE

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ORIGINAL DESIGN OF EACH NUMBER CONTOURS
AT 95 Km ALTITUDE



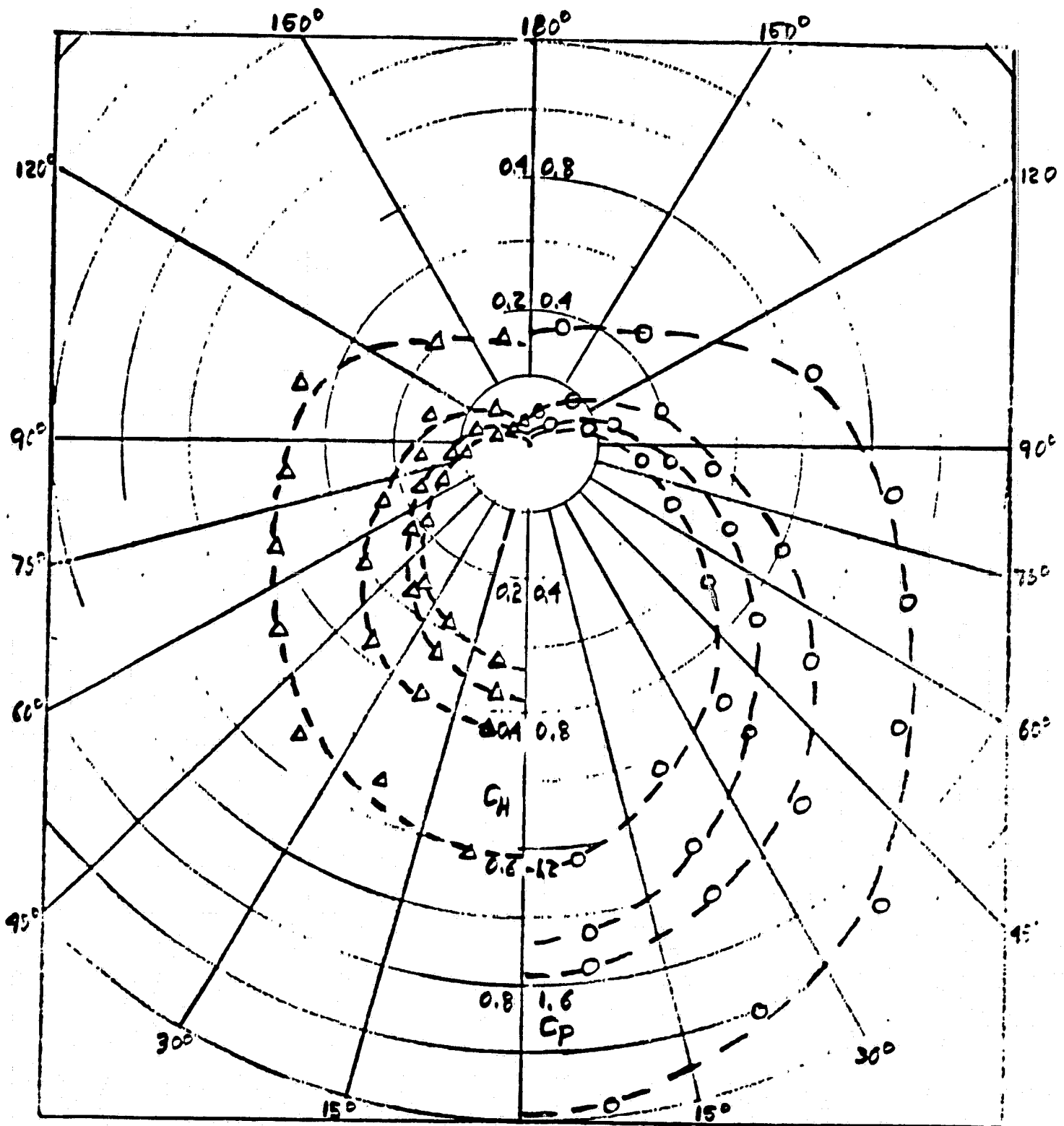


FIGURE 17. SURFACE PROPERTIES; PRESSURE C_P AND HEAT TRANSFER COEFFICIENT C_H VERSUS ANGLE IN CROSS-PLANE $x = 75$ cm (NEAR THE SUMS ENTRANCE) AT 95 Km ALTITUDE

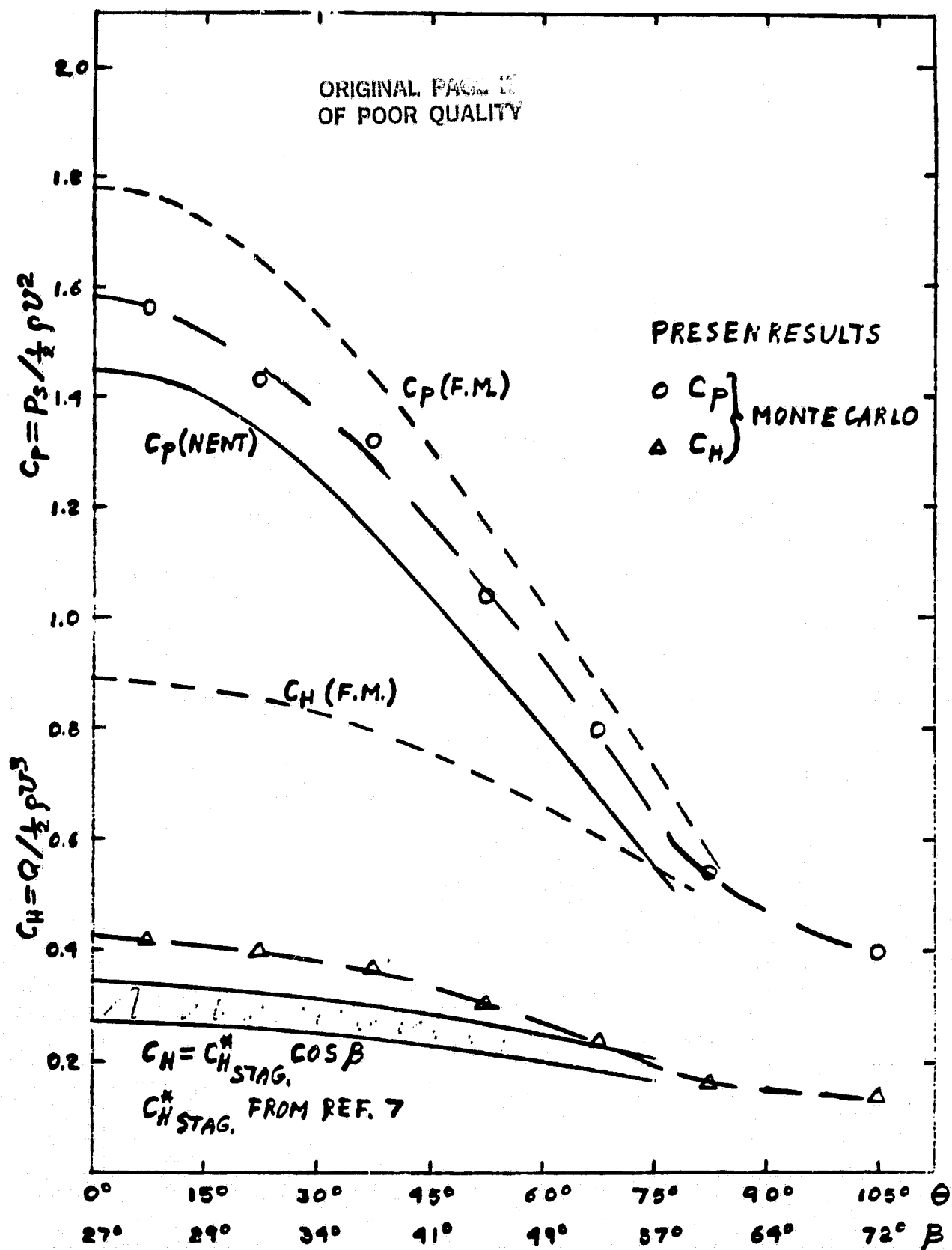


FIGURE 18. FLUX DISTRIBUTION FUNCTION AT SUMS INLET

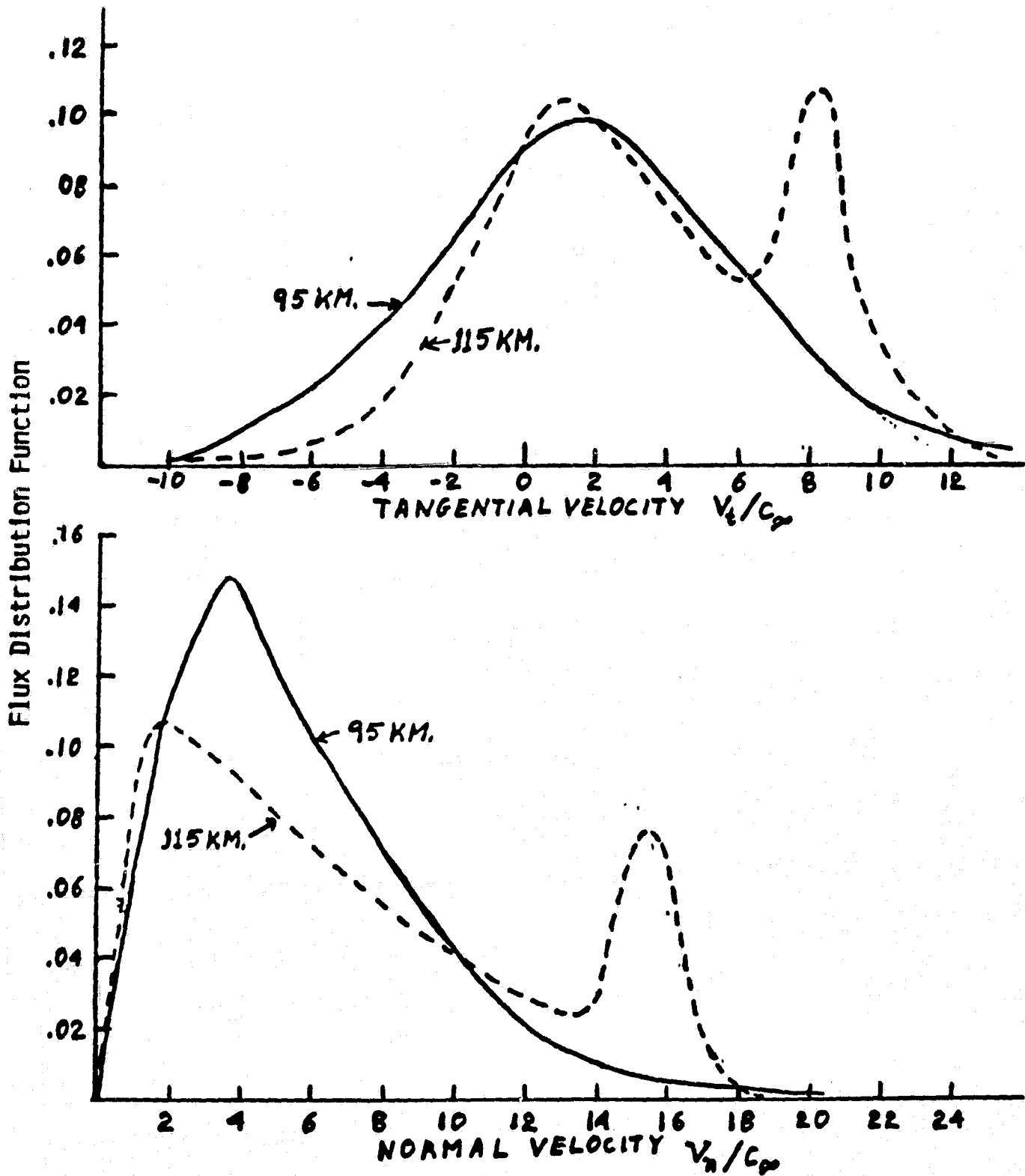


FIGURE 19. SUMS INLET GEOMETRY

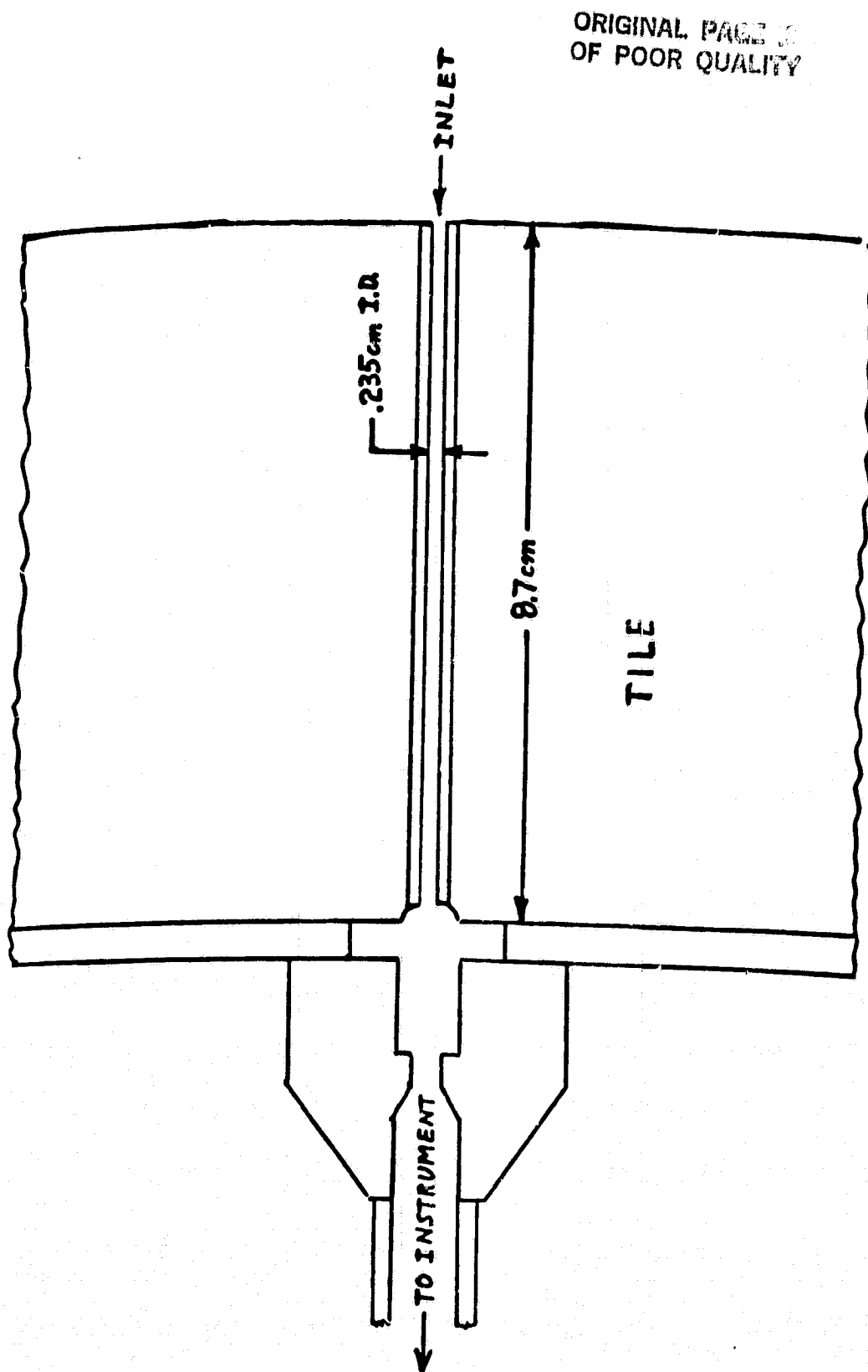
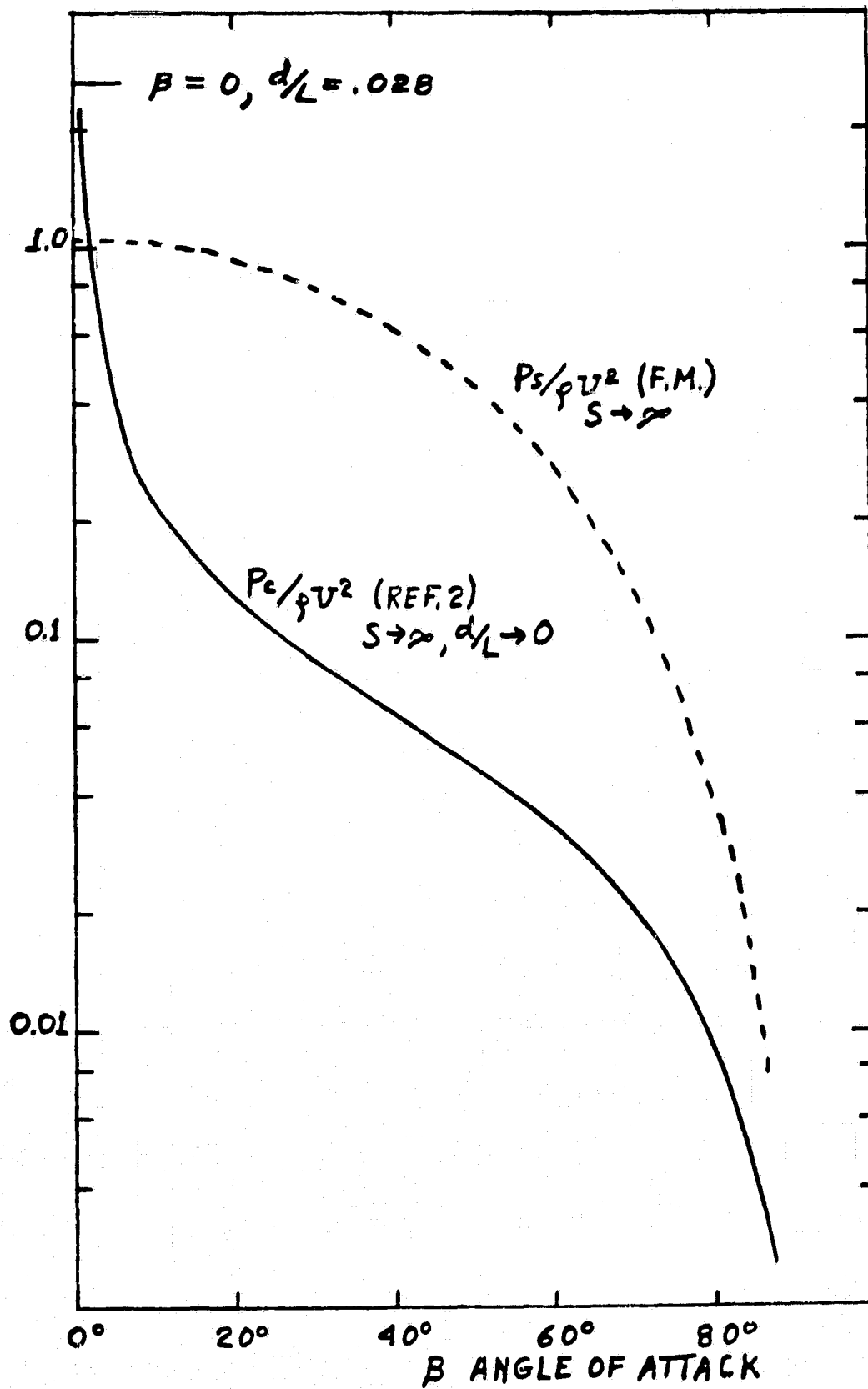


FIGURE 20. FREE MOLECULAR PROBE RESPONSE ($S \rightarrow \infty$, $d/L \rightarrow 0$)



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FIGURE 21. PRESSURE CORRECTION FACTOR (p_c/p_s)

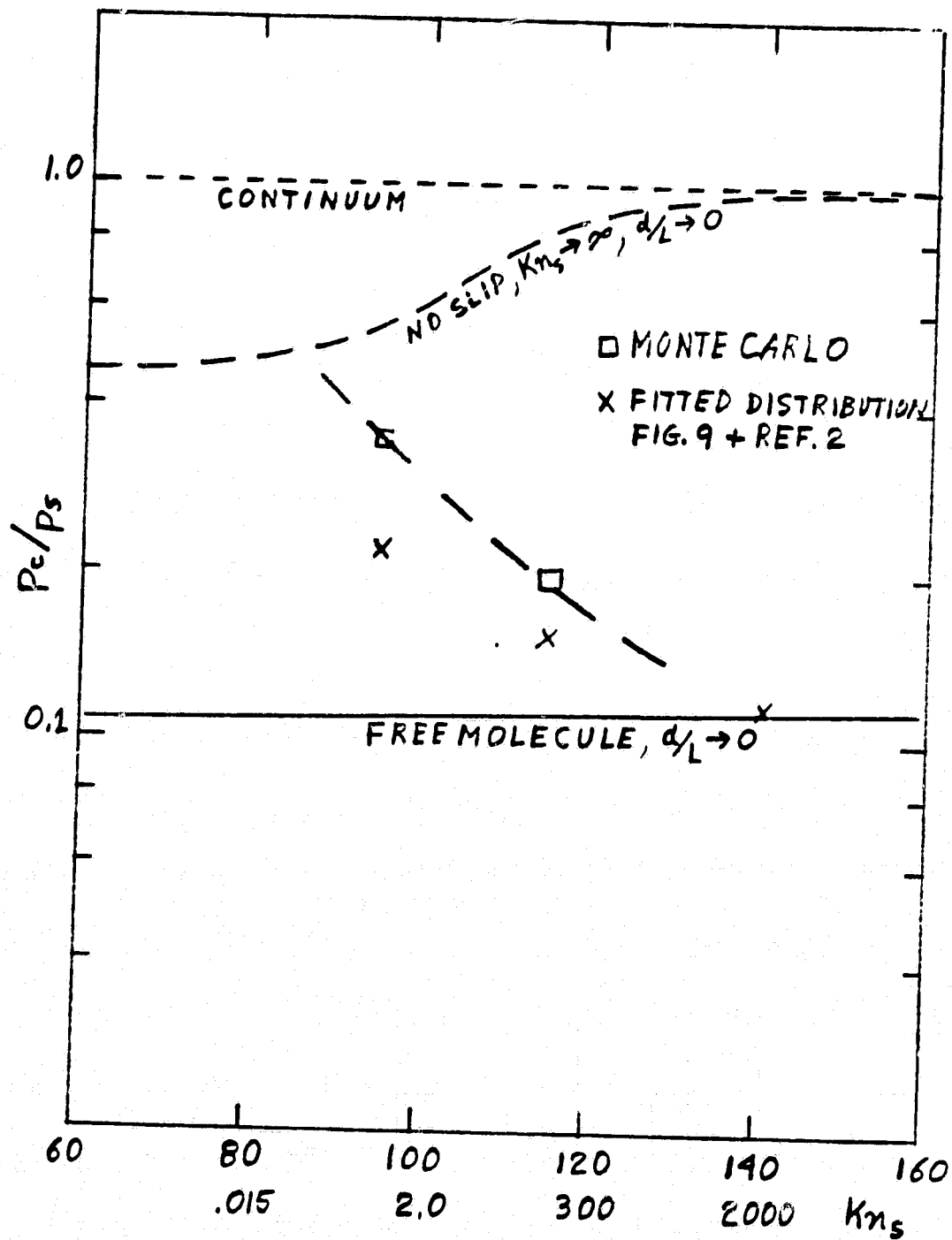


FIGURE 22. SURFACE PRESSURE AT SUMS INLET
VERSUS ALTITUDE

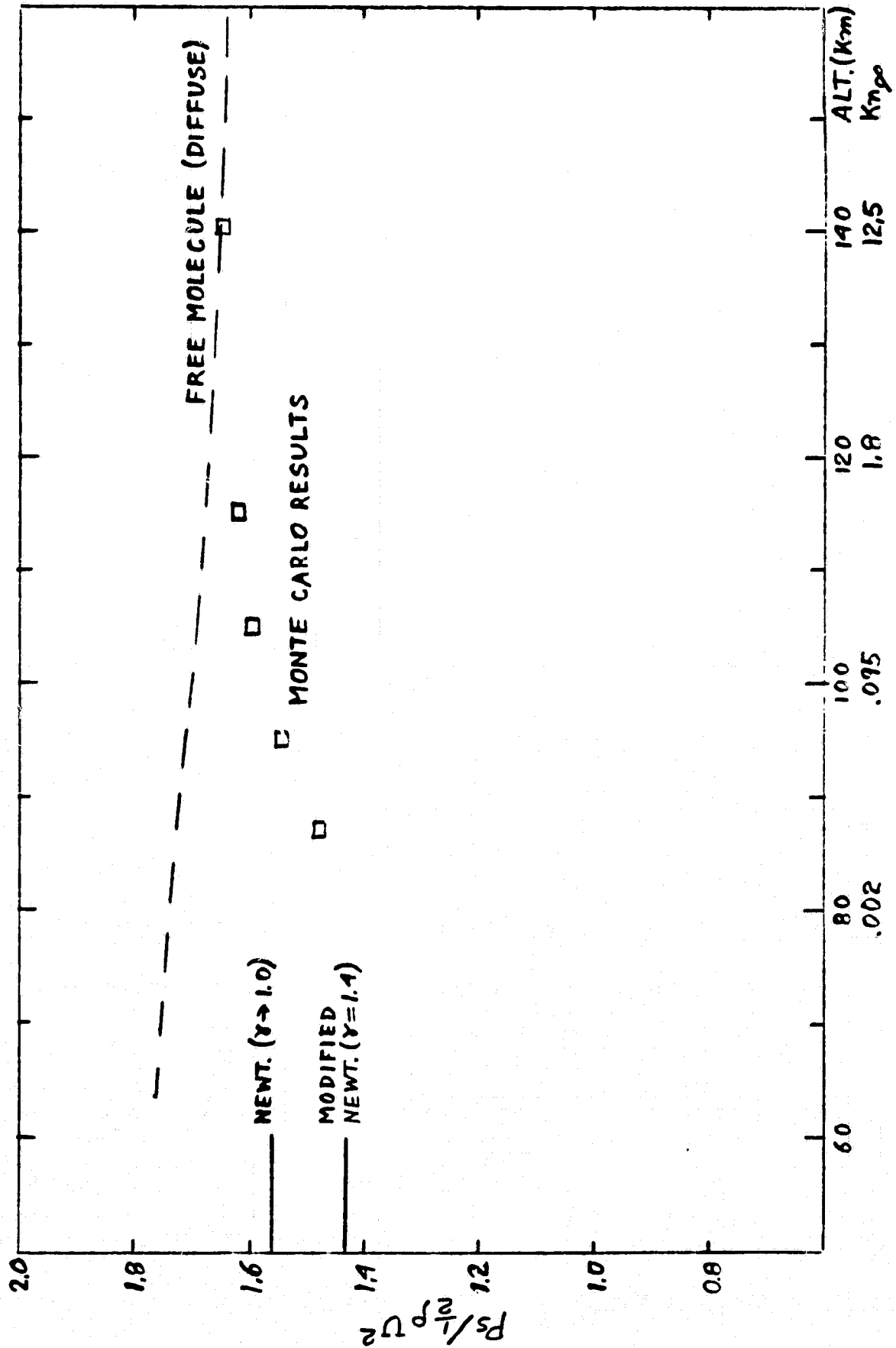
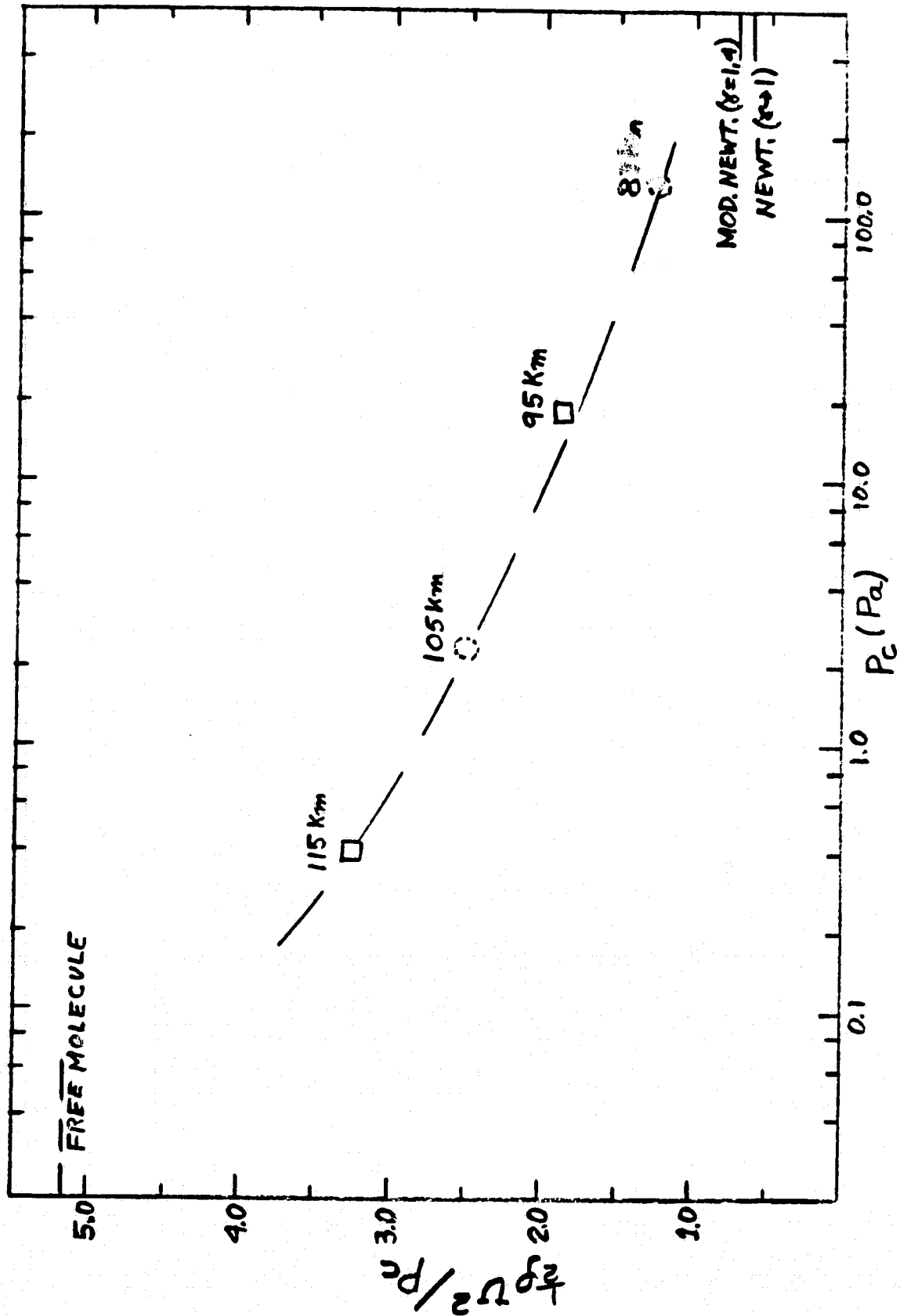


FIGURE 23. PRELIMINARY ESTIMATE OF DATA REDUCTION
RELATION q/p_c VERSUS p_c



APPENDIX A
COMPUTER CODE EXTERNAL


```
// JOB GKB 0367425.GKBSpace N=EXTCOMP REG=500 T=.8 P=150 C=0
// *FORMAT PR,DDNAME=SYSMSG,DEST=VM370
// EXEC PORTXCL,PARM.PORT='XREF'
// PORT.SYSIN DD *
```

MAIN PROGRAM FOR MONTE CARLO 3-D EXTERNAL FLOW CALCULATIONS
 OBJECTIVE OF THIS MAIN PROGRAM IS TO SET THE DIMENSIONS
 MAIN RUNNING PROGRAM IS *** RUN ***

FOLLOWING TWO CARDS HAVE TO BE ELIMINATED FOR NON IBM MACHINES

```
*****
INTEGER*2 LB,NBM,NBN,NB,NBP,NBT,NBS,NUMCEL
INTEGER*2 LM,LPF,LCOL,LKW
*****
```

THE NEXT CARD IS ASSOCIATED WITH PRINCETON RANDOM NUMBER GENERATOR

```
*****
COMMON/RANCON/NRAN(4)
*****
```

THE FOLLOWING DIMENSION STATEMENTS SET THE MAJOR ARRAY DIMENSIONS
 AND MUST BE CONSISTENT WITH THE FOLLOWING DATA CARD -----

NSP=NUMBER OF SPECIES - EXAMPLE BELOW NSP=1
 NMB=NUMBER OF BOXES WITHOUT COUNTING SUBDIVISIONS OR THOSE
 OCCUPIED BY THE BODY - EXAMPLE BELOW NMB=2600
 NMC=NUMBER OF FINAL CELLS - EXAMPLE BELOW NMC=1500
 NMP=MAX NUMBER OF MOLECULES OF EACH SPECIES ALLOWED IN PROGRAM.
 IF EXCEEDED, PROGRAM EITHER FAILS OR RESTARTS AT BEGINNING
 WITH NUMBER REDUCED BY 10% - EXAMPLE BELOW NMP=20000
 NPB=MAXIMUM NUMBER IN EACH CELL - EXAMPLE NPB=100

```
DIMENSION DBA(1,1500),NB(1,1500),NBP(1,1500),NBT(1,1500),LKW(1500)
DIMENSION TMP(1,1500),TMPA(1,1500),XV(1,1500),XVA(1,1500)
DIMENSION YV(1,1500),YVA(1,1500),ZV(1,1500),ZVA(1,1500),DB(1,1500)
DIMENSION TRP(1,1500),TRPA(1,1500),NBS(1,1500),NBM(1,1500)
DIMENSION NBN(1500),T(1,1,1500)
DIMENSION LB(20000),LM(1,20000),ER(1,20000)
DIMENSION LPF(1,20000),PAU(1,20000),PAV(1,20000),PAW(1,20000)
DIMENSION PAX(1,20000),PAY(1,20000),PAZ(1,20000),LCOL(1,20000)
DIMENSION FNB(2600),XC(2600),YC(2600),ZC(2600),NUMCEL(2600)
DATA NSP/1/,NMB/2600/,NMC/1500/,NMP/20000/,NPB/100/
```

```
2 FORMAT(/17X,'NORMAL TERMINATION OF THE PROGRAM')
NAMELIST/DIM/NSP,NMB,NMC,NMP,NPB,NRAN
```

INITIALIZATION OF RANDOM NUMBER GENERATOR - PRINCETON ROUTINE

```
NRAN(1)=0
NRAN(2)=0
NRAN(3)=0
NRAN(4)=0
```

PRINTOUT OF MAJOR ARRAY DIMENSIONS USED ABOVE

ORIGINAL PAGE IS
 OF POOR QUALITY

WRITE(6,DIM)

ORIGINAL PAGE 1
OF POOR QUALITY

CALL OF MAIN OPERATING PROGRAM WHICH REQUIRES INPUTS:
 &CONTRL,&TIMES,&FLOREF,&MOLEC,&SHAPES,&GEOM,&COUPLE
 THESE INPUTS ARE ALL CURRENTLY IN THE NAMELIST FORMAT
 AND MAY HAVE TO BE CHANGED IF THAT CONVENTION IS NOT AVAILABLE
 BRIEF DESCRIPTION OF THE PARAMETERS FOLLOWS

&CONTRL - ONE OCCURRENCE (NEW OR RESTART)

PARAMETER	DEFAULT	DEFINITION OR EXPLANATION
NAME	8 BLANKS	ANY ALPHANUMERIC NAME UP TO 8 CHARACTERS
TITLE	24 BLNKS	ANY ALPHANUMERIC TITLE UP TO 24 CHARACTERS
PERCNT	.001	ACCURACY IN INTEGRATION PROCEDURES
ICOPY	1	NUMBER OF ADDITIONAL COPIES OF OUTPUT
DUMP	.TRUE.	IF TRUE WILL CAUSE SYSTEM DUMP FOR ANY OF 12 PROGRAMMER DESIGNED ERROR HALTS.
DEBUG (1)	.FALSE.	IF TRUE WILL PRINT MESSAGE WHEN CELL POP. EXCEEDS MNB
DEBUG (2)	.FALSE.	IF TRUE WILL PRINT CPU TIME AROUND EACH PART OF LOOP
DEBUG (3)	.TRUE.	IF TRUE WILL PRINT CPU TIME REMAINING AT END OF LOOP
NEW	.TRUE.	IF TRUE - NEW RUN - IF FALSE - RESTART OF RUN
SAVE	.FALSE.	IF TRUE - SNAPSHOT SAVED ON TAPE(9) FOR RESTART
REDO	.FALSE.	IF TRUE PROGRAM WILL AUTOMATICALLY RESTART WITH 90% OF TOTAL IF TOTAL CELL POPULATION EXCEEDS MNM

&TIMES - ONE OCCURRENCE (NEW OR RESTART)

PARAMETER	DEFAULT	DEFINITION OR EXPLANATION
DTM	- - -	REAL NUMBER - FRACTION OF MEAN FREE TIME PER CYCLE
ITS	- - -	INTEGER - NUMBER OF CYCLES PER SAMPLE
ITP	- - -	INTEGER - NUMBER OF CYCLES BETWEEN PRINTOUTS
TST	- - -	INTEGER - ESTIMATE OF NUMBER OF CYCLES TO STEADY STATE
TLIM	- - -	INTEGER - TOTAL NUMBER OF CYCLES TO END OF RUN - WILL TERMINATE SOONER IF CPU TIME IS TO BE EXCEEDED

&FLOREF - ONE OCCURRENCE (NEW RUN ONLY)

PARAMETER	DEFAULT	DEFINITION
INM	- - -	INITIAL NUMBER OF MOLECULES INM<MNM< OR = NMP
MNM	- - -	MAXIMUM NUMBER OF MOLECULES PER SPECIES
MNB	- - -	MAXIMUM NUMBER PER CELL - DIAGNOSTIC ONLY
MSP	- - -	NUMBER OF MOLECULAR SPECIES (MAX. IS 3)
MET	0	IF 0 - DATA IS IN SI (METRIC) UNITS IF>0 - DATA IS IN ENGLISH UNITS
U	- - -	FLOW VELOCITY (M/SEC) OR (FT/SEC)
ANGLE	- - -	ANGLE OF ATTACK (DEGREES)
RNU	0.0	ARRAY GIVING MOLE FRACTIONS OF SPECIES IN FREE STREAM
RMA	0.0	ARRAY GIVING MOLECULAR WEIGHTS OF SPECIES ABOVE
TF	- - -	FREE STREAM TEMPERATURE (K OR R)
DENF	- - -	FREE STREAM NUMBER DENSITY (NUM/M**3 OR NUM/FT**3)

&MOLEC - ONE OCCURRENCE (NEW RUN ONLY)

PARAMETER	DEFAULT	DEFINITION
TRF	- - -	REFERENCE TEMPERATURE FOR MOLECULAR DATA
DIR	0.0	CROSS-SECTIONS AT REFERENCE TEMP. (MSPX MSP)
ETA	0.0	PARAMETERS IN DIFFUSION AND VISCOSITY LAW (MSPX MSP)
PHI	0.0	PARAMETERS FOR ROTATIONAL RELAXATION (MSPX MSP)
CHI	0.0	ROTATIONAL DEGREE OF FREEDOM PARAMETER (NROT/2 - 1)
ACR	.001	ACCURACY IN MOLECULAR COLLISION CALCULATIONS

ORIGINAL PAGE IS
OF POOR QUALITY

ESHAPE - ND+1 OCCURRENCES WHERE ND=NUMBER OF BODY SEGMENTS (NEW RUN)

PARAMETER DEFAULT DEFINITION

FIRST OCCURRENCE

BODY(1) 0.0 STARTING POINT OF BODY FROM FRONT OF CELLS (M OR FT)

BODY(I) I>1 - - NEED NOT BE SPECIFIED

SUBSEQUENT OCCURRENCES (ND)

BODY(1) - - - X COORDINATE FROM FRONT OF BODY OF THE DOWNSTREAM
EDGE OF THE CURRENT BODY SEGMENT

BODY(2) - - - TEMPERATURE OF THIS BODY SEGMENT

BODY(3) - - - SURFACE AREA/TOTAL CROSS-SECTIONAL AREA FOR SEGMENT.
IF 0.0 PROGRAM WILL COMPUTE THIS QUANTITY

BODY(4) - - - SWITCH - IF 0.0 THIS SEGMENT'S EQ. WILL APPEAR LATER
IF >0.0 THE EQ. OF THIS AND PRECEDING SEGMENTS IS
GIVEN BY BODY(6+2*MSP) TO BODY(9+2*MSP)

BODY(5) - - - SWITCH - IF NOT 0.0 THIS IS THE LAST SHAPES CARD

BODY(I) I EVEN ALPHA - ENERGY ACCOMODATION COEFFICIENT FOR SPECIES

BODY(J) J ODD SIGMA - TANGENTIAL ACCOMODATION COEFF. FOR SPECIES
I AND J < (6+2*MSP)

BODY(6+2*MSP) ORIGIN OF COORDINATES WITH RESPECT TO BODY START
FOR THE EQUATION OF THIS BODY SECTION

BODY(7+2*MSP) - COEFFICIENTS A,B,C IN THE EQUATION

BODY(9+2*MSP) R**2+A*X**2+B*X+C=0.0 FOR THIS BODY SECTION

EGEOM - ONE OCCURRENCE (NEW RUN ONLY)

PARAMETER DEFAULT DEFINITION

NWEDGE - - - TWO INTEGERS GIVING THE NUMBER OF WEDGES BELOW
AND ABOVE THE ANGLE THETAZ

THETAZ - - - ANGLE FROM NEG. Y AXIS DIVIDING DIFF. WEDGE SIZES

RMB - - - MAX. BODY RADIUS - IF 0. WILL BE COMPUTED BY PROGRAM

BW - - - WIDTH (DEL X) OF FIRST LEVEL CELLS (M OR FT)

BH - - - HEIGHT (DEL R) OF FIRST LEVEL CELLS (M OR FT)

NW - - - NUMBER OF FIRST LEVEL CELLS IN X DIRECTION

NH - - - NUMBER OF FIRST LEVEL CELLS IN RADIAL DIRECTION

NL 1 NUMBER OF LEVELS OF CELLS

NFA 0 NUMBER OF FIRST LEVEL CELLS AHEAD OF LEVEL 2

NCA 0 NUMBER OF FIRST LEVEL CELLS SUBDIVIDED INTO SECOND
LEVEL CELLS ALONG THE X DIRECTION

NHA 0 AS ABOVE BUT IN RADIAL DIRECTION

MW 0 NUMBER OF SECOND LEVEL CELLS IN THE X DIRECTION

MH 0 NUMBER OF SECOND LEVEL CELLS IN THE RADIAL DIRECTION

NFB 0 NUMBER OF SECOND LEVEL CELLS AHEAD OF LEVEL 3

NCB 0 NUMBER OF SECOND LEVEL CELLS SUBDIVIDED INTO THIRD
LEVEL CELLS ALONG THE X DIRECTION

NHB 0 AS ABOVE BUT IN RADIAL DIRECTION

LW 0 NUMBER OF THIRD LEVEL CELLS IN X DIRECTION

LH 0 NUMBER OF THIRD LEVEL CELLS IN RADIAL DIRECTION

LD - - - NUMBER OF FIRST LEVEL CELLS FROM AXIS IN R DIRECTION
FOR WEIGHTING FACTOR BOUNDARIES (5 INTEGERS)

LF - - - WEIGHTING FACTOR RATIOS AT BOUNDARIES ABOVE (5 INTEGERS)

ESCOUPLE - ONE OCCURRENCE (NEW RUN ONLY) - DISTRIBUTION FUNCTION

PARAMETER DEFAULT DEFINITION

NS - - - THE NUMBER OF BODY SEGMENTS FOR ACCUMULATING VELOCITY
DISTRIBUTION FUNCTION INFORMATION

```

C MS      - - -   ARRAY(NS) OF AXIAL SEGMENT NUMBERS
C IWS     - - -   ARRAY(NS) OF AZIMUTHAL WEDGE NUMBERS
C VEL     3.,3.,4. THERMAL VELOCITY SPREAD FOR THE UNCOLLIDED MOLECULES
C MJ      20      NUMBER OF DEL V REGIONS FOR SAMPLING VELOCITY SPACE
C SL      - - -   ARRAY GIVING LOWER BOUND ON THE VELOCITY SAMPLE OF
C          COLLIDED MOLECULES (MSP X NS X 3)
C DELS    - - -   THE RANGE (SL<V<SL+DELS) FOR SAMPLE OF COLLIDED
C          MOLECULES (MSP X NS X 3)

```

IF DISTRIBUTION FUNCTION INFORMATION IS NOT DESIRED USE:
 &COUPLE NS=0 &END

A SAMPLE INPUT DECK IS GIVEN BELOW:

```

&CONTRL NAME='SHUT', 'TLE ', TITLE='HYPE', 'RBOL', 'A AT', ' 95K', 'M M', 'ON. ',
  DEBUG=.F., .T., .T., NEW=.T., SAVE=.T., ICOPY=0, REDO=.T. &END
&TIMES DTM=.025, ITS=5, ITP=1000, TST=400, TLIM=1000 &END
&FLOREF INM=4500, MNM=20000, MNB=100, MSP=1, MET=0, U=7485.9, ANGLE=0.0, RNU=1., 2*0.,
  RMA=28.94, 0., 0., TF=195.51, DENF=2.52E+19 &END
&MOLEC TRF=1000, DIR=3.5E-19, ETA=.104, PHI=0.0, CHI=-1., ACR=.001 &END
&SHAPES BODY=1.00 &END
&SHAPES BODY=.0173, 1590., 3*0.0, 2*1.0 &END
&SHAPES BODY=.0672, 1590., 3*0.0, 2*1.0 &END
&SHAPES BODY=.1444, 1590., 3*0.0, 2*1.0 &END
&SHAPES BODY=.2432, 1590., 3*0.0, 2*1.0 &END
&SHAPES BODY=.3579, 1590., 3*0.0, 2*1.0 &END
&SHAPES BODY=.4842, 1590., 3*0.0, 2*1.0 &END
&SHAPES BODY=.6192, 1590., 3*0.0, 2*1.0 &END
&SHAPES BODY=.7405, 1590., 3*0.0, 2*1.0 &END
&SHAPES BODY=.7500, 1590., 0., 1., 0., 2*1., 0., -1.423278, -2.286, 0. &END
&SHAPES BODY=.9000, 1590., 0.0, 4*1.0, 0.9000, -111.82, 2*0.0 &END
&GEOM NWEDGE=1, 0, THETAZ=180., RMB=0.0, BW=.05, BH=.1, NW=40, NH=24, NL=2,
  NPA=15, NCA=20, NHA=16, MW=40, MH=32, LD=1, 2, 4, 8, 12, LPF=3, 2, 2, 2, 1 &END
&COUPLE NS=3, MS=1, 5, 8, IWS=1, 1, 1, VEL=2.5, 2.5, 2.0, MJ=20, SL=9*1., 18*-9.,
  DELS=27*20. &END

```

```

CALL RUN(NSP, NMB, NMC, NMP, NPB, DBA, NB, NBF, NBT, TMP, TMPA, XV, XVA, YV,
  1 YVA, ZV, ZVA, T, DB, FNB, XC, YC, ZC, NUMCEL, LM, LPF, PAU, PAV, PAW, PAX, PAY,
  2 PAZ, LCOL, TRP, TRPA, ER, LKW, NBS, LB, NBM, NBN)
WRITE(6, 2)
STOP
END

```

```

SUBROUTINE RUN(NSP, NMB, NMC, NMP, NPB, DBA, NB, NBF, NBT, TMP, TMPA, XV, XVA,
  1 YV, YVA, ZV, ZVA, T, DB, FNB, XC, YC, ZC, NUMCEL, LM, LPF, PAU, PAV, PAW, PAX,
  2 PAY, PAZ, LCOL, TRP, TRPA, ER, LKW, NBS, LB, NBM, NBN)
  MAIN RUNNING PROGRAM ** RUN *** CALLS ALL OTHER SUBROUTINES

```

```

INTEGER*2 LM, LPF, LCOL, LKW
INTEGER*2 LB, NBM, NBN, NB, NBF, NBT, NBS, NUMCEL
INTEGER PRT, SAMP, TST, TLIM, TIME, Q
LOGICAL DUMP, DEBUG(3), SAVE, NEW, REDO

```

```

REAL INTGRAL,LAM,MU,NU,JAY,KAY
DIMENSION BTA(3),C1(3),C2(3),C3(3),C7(3),C8(3),DFA(3),FL(3)
DIMENSION DELANG(2),FDN(3),HTI(3),HTR(3),JNT(3),KNM(3),NM(3),SR(3)
DIMENSION NAME(2),TITLE(6),NWEDGE(2),LD(5),LP(5),LWF(6),RLD(6)
DIMENSION RNU(3),RMA(3),WTM(3),CHI(3),DIR(3,3),DAM(3,3),PHI(3,3)
DIMENSION ETA(3,3),CN8(3,3),CNG(3),CMG(3),CN(3,3,3),CM(3,3,3)
DIMENSION CTI(3,3),CTR(3,3),CNI(3,3),CNR(3,3),LEV(3),SN(3),ST(3)
DIMENSION D1(3),D2(3),D3(3),D4(3),BODY(15),DBG1(3,3),LIMIT(10)
DIMENSION VL(3,3),DELV(3,3),SSA(2,3),SSB(2,3),VEL(3),COEFF(4,9)
DIMENSION XLIM(12),NTCF(3,3),MS(3),IWS(3),TANGN(3),NCOL(3,3)
DIMENSION FV(3,3,2,20,3),NTCV(3,3,2,20,3),SL(3,3,3),DELS(3,3,3)
DIMENSION XCB(18),XS(18),YCB(18),TB(18),ALPHA(3,18),SIGMA(3,18)
DIMENSION NTS(3,18,12),NTSF(3,18,12),UTL(3,18,12),UTT(3,18,12)
DIMENSION VTS(3,18,12),HTSI(3,18,12),HTS(3,18,12)
DIMENSION UTLI(3,18,12),UTTI(3,18,12),VTSI(3,18,12)
DIMENSION ENT(2,3,6,12),REM(2,3,6,12),ENTS(3,12),REMS(3,12)
DIMENSION PTH(3,12),THETA(12),DTH(12)
DIMENSION LB(NMP),NBN(NMC),NBM(NSP,NMC),LM(NSP,NMP)
DIMENSION ER(NSP,NMP),TRP(NSP,NMC),TRPA(NSP,NMC)
DIMENSION DBA(NSP,NMC),NB(NSP,NMC),NBF(NSP,NMC),NBT(NSP,NMC)
DIMENSION TMP(NSP,NMC),TMPA(NSP,NMC),XV(NSP,NMC),XVA(NSP,NMC)
DIMENSION YV(NSP,NMC),YVA(NSP,NMC),ZV(NSP,NMC),ZVA(NSP,NMC)
DIMENSION T(NSP,NSP,NMC),DB(NSP,NMC),LKW(NMC),NBS(NSP,NMC)
DIMENSION PNB(NMB),YC(NMB),YC(NMB),ZC(NMB),NUMCEL(NMB)

```

FORMATS

ORIGINAL FORMATS
OF POOR QUALITY

```

*****
1 FORMAT(1H1)
2 FORMAT(1H1/17X,'RARIFIED SUPERSONIC FLOW OF BINARY GAS',T74,'I')
3 FORMAT('+',103X,'COPY ',I2)
4 FORMAT(/17X,'FLOW THROUGH ALL THE BOUNDARIES'//)
5 FORMAT(3X,3I4,6E18.6/)
6 FORMAT(7X,2I4, E18.6,72X,E18.6)
30 FORMAT('1TIME = ',F6.3,60X,'RANDOM NUMBER GENERATOR HAS BEEN CALLED
1 ',I10,' TIMES')
31 FORMAT(' CPU TIME LEFT- ',F8.3)
32 FORMAT(7X,'-MOLECULES-'/3X,3I6)
33 FORMAT(' TIME = ',F8.3,5X,'COLLISION LOOP=',F8.3,5X,'MOVE LOOP = '
1,F8.3,5X,'TOTAL TIME = ',F8.3/21X,'2ND MOVE LOOP = ',F8.3,5X,
2'CLEANUP LOOP=',F8.3,4X,'PARTICLE NUMBERS = ',4I6)
34 FORMAT(9X,'-MOLECULAR COLLISIONS-'/3(3I14/))
35 FORMAT(2X,'-COLLISIONS WITH SURFACE-'/3X,3I8)
36 FORMAT(' MAXIMUM NUMBER OF MOLECULES SO FAR- ',I6//)
38 FORMAT(' EXCESS MOLECULES OCCURRED IN ',3A4)
40 FORMAT('/' SOMETHING IS WRONG WITH BOX NUMBERING IN RUN '/5I5,2E17.
17,3I5,2E17.7,5I5/2E17.7)
42 FORMAT('/' SOMETHING WRONG IN COMPUTING XSTART'/1X,8E16.8)
44 FORMAT(' NB('I2,',',I4,') POPULATION EXCEEDED ',I3,' IN MAIN AT TR
1TIME = ',F7.3)
50 FORMAT(///' SNAP SAVED ON TAPE')

*****
CPA=ELTIME(0)
CALL NOUNDF
LIMIT(4)=NMC
LIMIT(5)=NMP
LIMIT(6)=NPB
LIMIT(7)=NMB
LIMIT(10)=NSP
KAWLS=0
PI=3.141593
PIROOT=SQRT(PI)
MET=0
LARGE=0
NL=1
NPA=0
NCA=0
NHA=0
MW=0
MH=0
NFB=0
NCB=0
NHB=0
LW=0
LH=0

```

```

RUN0690
RUN0700
RUN0710
RUN0720
RUN0730
RUN0740
RUN0750
RUN0760
RUN0770
RUN0800
RUN0810
RUN0880
RUN0890
RUN0900
RUN0910
RUN0920
RUN0930
RUN0940
RUN0950
RUN0960
RUN0970
RUN0980

```

```

RUN1030
RUN1040
RUN1060
RUN1080
RUN1090
RUN1100
RUN1110
RUN1120
RUN1130
RUN1140
RUN1150
RUN1160
RUN1170
RUN1180

```

```

MJ=20
DUMP=.TRUE.
DEBUG(1)=.FALSE.
DEBUG(2)=.FALSE.
DEBUG(3)=.TRUE.
SAVE=.FALSE.
NEW=.TRUE.
REDO=.FALSE.
PERCNT=.001
ACR=.001
DO 58 I=1,15
58 BODY(I)=0.0
DO 60 I=1,3
RNU(I)=0.0
RMA(I)=0.0
CHI(I)=0.0
DO 59 J=1,18
ALPHA(I,J)=1.0
59 SIGMA(I,J)=1.0
DO 60 K=1,3
ETA(I,K)=0.0
PHI(I,K)=0.0
60 DIR(I,K)=0.0
VEL(1)=3.
VEL(2)=3.
VEL(3)=4.
WRITE(6,1)
READ(5,CONTRL)
WRITE(6,CONTRL)
IF(NEW) GO TO 103
REWIND 9
READ(9) DENF,U,XREF,TRF, KAWLS,NL,NW,NH,MW,MH,LW,LH,NXA,NXB,NCA
1 ,NCB,NFA,NFB,NHA,NHB,BW,BH,BWB,BHB,BWC,BHC,XLB,XLC,PI,NREG
2 ,S,SINANG,COSANG,AKN,NBX,RM,XR,ND,TIME,DTM,TI,ITS,ITP,TST
3 ,TLIM,RMA,RNU,DIR,XSTART,JNM,MNM,MNB,TR,BZC,CN7,DRF,PCF
4 ,FNA,HTF,INM,ITYPE,JTYPE,MJ,NAV,NMAX,NS,NWEDG,PRT,SAMP
5 ,BTA,C1,C2,C3,C7,C8,DAM,DFA,FL,DELANG,PDN,HTI,HTR,JNT,KNM
6 ,NM,WTM,C4,VRM,NCOL,LD,LF,LWF,RLD,CTI,CTR,CNI,CNR,LEV,SN
7 ,ST,D1,D2,D3,D4,SSA,SSB,MS,NSP,NMB,NMC,NMP,NPB,NRAN,VELR
8 ,IWS,TANGN,XLIM,COEFF,XCB,XS,YCB,TB,ALPHA,SIGMA,NTS,NTSF
9 ,UTL,UTT,VTS,HTS,HTSI,ENT,REM,ENTS,REMS,PTH,THETA,DTH,TMPAR
A ,DBA,NB,NBF,NBT,TMP,XV,XVA,YV,YVA,ZV,ZVA,T,DE,FNB,XC,YC,ZCR
B ,NUMCEL,PAU,PAV,PAW,PAX,PAY,PAZ,FV,NTCV,NTCF,LPF,LCOL,LM
C,ETA,PHI,CHI,CN,CM,CNG,CMG,CN8,TRP,TRPA,THETAZ,NWEDGE,MSP,ANGLE,TF
D,UTLI,UTTI,VTSI,ER,RMB,LKW,NBS,LB,NBM,NBN
REWIND 9
DTMO=DTM
READ(5,TIMES)
WRITE(6,TIMES)
IF(DTM.EQ.DTMO) GO TO 100
AIME=TIME*DTMO
TIME=AIME/DTM+0.1
DO 99 J=1,NSP
DO 99 L=1,NWEDG
ENTS(J,L)=ENTS(J,L)*DTM/DTMO

```

RUN1190

ORIGINAL PAGE 1
OF POOR QUALITY

RUN1230

RUN1260

RUN1240

RUN1250

RUN1300

RUN1310

RUN1320

RUN1330

RUN1340

RUN1360

RUN1370

RUN1380

RUN1390

RUN1400

RUN1410

RUN1420

RUN1430

RUN1440

RUN1460

RUN1470

RUN1480

RUN1490

RUN1500

ORIGINAL PAGE IS
OF POOR QUALITY

```

DO 98 K=1,6
DO 98 I=1,2
ENT(I,J,K,L)=ENT(I,J,K,L)*DTM/DTMO
98 CONTINUE
99 CONTINUE
100 IF(TI.GT.0.0) TST=TI/DTM
WRITE(6,2)
WRITE(6,4)
WRITE(6,5) ((I,J,L,(ENT(I,J,K,L),K=1,6),L=1,NWEDG),J=1,MSP),I=1,2)
WRITE(6,6) ((J,L,ENTS(J,L),PTH(J,L),L=1,NWEDG),J=1,MSP)
WRITE(6,2)
CALL PRINTA(THETAZ,NWEDGE,TITLE,NAME,XCB,YCB,TB,ALPHA,SIGMA,LD,LP,
1 XLIM,COEFF,LIMIT,MSP)
CALL PRINTB(PNA,MSP,PNB,LEV,LWF,NH,RLD,XLIM,XC,YC,ZC,NB,NUMCEL,LKW
1,NSP)
GO TO 280
103 READ(5,TIMES)
WRITE(6,TIMES)
READ(5,FLOREF)
WRITE(6,FLOREF)
READ(5,MOLEC)
WRITE(6,MOLEC)
IF(MSP.GT.LIMIT(10)) CALL DIAG(10,LIMIT(10),MSP)
CHIM=0.0
RMR=0.0
DMR=0.0
DO 115 M=1,MSP
RMR=RMR+RMA(M)*RNU(M)
CHIM=CHIM+CHI(M)*RNU(M)
DO 115 K=1,MSP
115 DMR=DMR+RNU(M)*RNU(K)*DIR(M,K)*(TRF/TF)**(ETA(M,K)/2.)
XREF=1./(DENF*DMR*1.414214)
VELR=SQRT(16628.64*TF/RMR)
IF(MET.NE.0) VELR=SQRT(99437.92*TF/RMR)
THR=XREF/VELR
S=U/VELR
XLIM(1)=0.
TR=0.
YR=0.
NREG=0
ND=0
READ(5,SHAPES)
WRITE(6,SHAPES)
XO=BODY(1)/XREF
104 READ(5,SHAPES)
WRITE(6,SHAPES)
ND=ND+1
IF(ND.GT.LIMIT(3)) CALL DIAG(3,LIMIT(3),ND)
XCB(ND)=BODY(1)/XREF+XO
TB(ND)=BODY(2)/TF
YCB(ND)=BODY(3)
DO 1104 M=1,MSP
ALPHA(M,ND)=BODY(4+2*M)
1104 SIGMA(M,ND)=BODY(5+2*M)
IF(TB(ND).GT.TR) TR=TB(ND)

```

RUN1540

RUN1550

RUN1590

RUN1610

RUN1620

RUN1630

RUN1640

RUN1650

RUN1670

RUN1680

RUN1690

RUN1770


```
IF (YCB(ND) .GT. YR) YR=YCB(ND) RUN1780
IF (BODY(4) .EQ. 0.0) GO TO 104
NREG=NREG+1 RUN1800
XLIM(NREG+2)=XCB(ND)
XT=BODY(6+2*MSP)/XREF+XO
A=BODY(7+2*MSP)
B=BODY(8+2*MSP)/XREF
C=BODY(9+2*MSP)/XREF**2
COEFF(1,NREG)=A
COEFF(2,NREG)=1.0
COEFF(3,NREG)=B-2.*A*XT
COEFF(4,NREG)=A*XT**2-B*XT+C
IF (BODY(5) .EQ. 0.0) GO TO 104
IF (NREG.GT.LIMIT(2)) CALL DIAG(2,LIMIT(2),NREG)
NSTEP=NREG+2
A=COEFF(1,1)
B=COEFF(3,1)
C=COEFF(4,1)
DISC=B*B-4.*A*C
IF (DISC.LT.0.) DISC=0.
DISC=SQRT(DISC)
IF (A.NE.0.) GO TO 108
XSTART=-C/B
GO TO 109
108 X1=.5/A*(-B+DISC)
X2=.5/A*(-B-DISC)
XSTART=AMIN1(X1,X2)
XMAX=AMAX1(X1,X2)
IF (XMAX.LT.XLIM(3)) XSTART=XMAX
IF (XSTART.GT.0.) GO TO 109
WRITE(6,42) XSTART,A,COEFF(2,1),B,C,DISC,X1,X2
IF (DUMP) CALL ABEND(1)
STOP
109 XLIM(2)=XSTART
A=COEFF(1,NREG)
B=COEFF(3,NREG)
C=COEFF(4,NREG)
DISC=B*B-4.*A*C
IF (DISC.LT.0.) DISC=0.
DISC=SQRT(DISC)
IF (A.NE.0.) GO TO 111
XLIM(NSTEP)=-C/B
GO TO 112
111 X1=.5/A*(-B+DISC)
X2=.5/A*(-B-DISC)
XLIM(NSTEP)=AMAX1(X1,X2)
XMIN=AMIN1(X1,X2)
IF (XMIN.GT.XLIM(NSTEP-1)) XLIM(NSTEP)=XMIN
IF (XLIM(NSTEP).GT.XLIM(NSTEP-1)) GO TO 112
WRITE(6,42) XLIM(NSTEP),A,COEFF(2,NREG),B,C,DISC,X1,X2
IF (DUMP) CALL ABEND(1)
STOP
112 AKN=1./(XLIM(NSTEP)-XSTART)
XCB(ND)=XLIM(NSTEP)
DO 260 N=1,3
```

ORIGINAL PAGE IS
OF POOR QUALITY

RUN1860
RUN1870
RUN1880
RUN1890
RUN1900
RUN1910
RUN1920
RUN1930
RUN1940
RUN1950
RUN1960
RUN1970
RUN1980
RUN1990

RUN2000
RUN2010
RUN2020
RUN2030
RUN2040
RUN2050
RUN2060
RUN2070
RUN2080
RUN2090
RUN2100
RUN2110
RUN2120
RUN2130
RUN2140
RUN2150
RUN2160

RUN2170
RUN2180

ORIGINAL PAGE 13
OF POOR QUALITY

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MS(N)=0
IWS(N)=0
DO 259 M=1,3
NCOL(N,M)=0
DO 258 K=1,3
SL(N,M,K)=0.0
258 DELS(N,M,K)=0.0
259 CONTINUE
260 CONTINUE
READ(5,GEOM)
READ(5,COUPLE)
WRITE(6,GEOM)
WRITE(6,COUPLE)
BW=BW/XREF
BH=BH/XREF
RMB=RMB/XREF
IF(RMB.GT.0.) GO TO 264
DO 262 K=1,NREG
XBEG=XLIM(K+1)
XEND=XLIM(K+2)
A=COEPP(1,K)
B=COEPP(2,K)
C=COEPP(3,K)
D=COEPP(4,K)
REND=SQRT(ABS((A*XEND**2+C*XEND+D)/B))
XPEAK=0.
IF(A.NE.0.) XPEAK=-.5*C/A
IF((XPEAK.LE.XBEG).OR.(XPEAK.GE.XEND)) GO TO 261
RTEMP=SQRT(ABS((A*XPEAK**2+C*XPEAK+D)/B))
IF(RTEMP.GT.REND) REND=RTEMP
261 IF(REND.GT.RMB) RMB=REND
262 CONTINUE
264 CONTINUE
NWEDG=NWEDGE(1)+NWEDGE(2)
IF(NWEDGE(2).EQ.0) THETAZ=180.0
IF(SAVE) REWIND 9
IF(NWEDG.GT.LIMIT(1)) CALL DIAG(1,LIMIT(1),NWEDG)
IF(MNM.GT.LIMIT(5)) CALL DIAG(5,LIMIT(5),MNM)
IF(MNB.GT.LIMIT(6)) CALL DIAG(6,LIMIT(6),MNB)
IF(NS.GT.LIMIT(8)) CALL DIAG(8,LIMIT(8),NS)
IF((NS.NE.0).AND.(MJ.GT.LIMIT(9))) CALL DIAG(9,LIMIT(9),MJ)
JNM=INM
DELANG(1)=THETAZ/NWEDGE(1)
DELANG(2)=0.0
IF(NWEDGE(2).NE.0) DELANG(2)=(180.-THETAZ)/NWEDGE(2)
SINANG=SIN(ANGLE/180.*PI)
COSANG=COS(ANGLE/180.*PI)
XR=BW*NW
XLIM(NSTEP+1)=XR
RM=BH*NH
VOL=PI*RM*RM*XR
NXA=NW*NH*NWEDG
NXB=MW*NH*NWEDGE(1)
NYC=LW*LH*NWEDGE(1)
IXA=NW*NH*NWEDGE(1)

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RUN2200

RUN2220

RUN2230

RUN2240

RUN2250

RUN2290

RUN2300

RUN2320

RUN2330

RUN2380

RUN2390

RUN2400

RUN2420

RUN2430

RUN2440

RUN2450

```

NBX=NXA+NXB+NXC
IF (NBX.GT.LIMIT(7)) CALL DIAG(7,LIMIT(7),RST)
BR=SQRT(TR)
DO 113 N=1,5
113 RLD(N)=BH*LD(N)
RLD(6)=RH
LWF(1)=1
RWFM=RLD(1)
B=RWFM*RWFM
C=B
DO 114 N=2,6
A=RLD(N)*RLD(N)
LWF(N)=LWF(N-1)*LF(N-1)
B=B+(A-C)/LWF(N)
D=RLD(N)/LWF(N)
IF (D.GT.RWFM) RWFM=D
114 C=A
INM=INM*RH*RH/B
DDN=INM/VOL
DO 140 MT=1,MSP
WTM(MT)=RMA(MT)/RMB
FDN(MT)=RNU(MT)*DDN
DFA(MT)=RNU(MT)
BTA(MT)=SQRT(WTM(MT))
SR(MT)=B*BTA(MT)
SN(MT)=SR(MT)*COSANG
ST(MT)=SR(MT)*SINANG
DO 117 K=1,MSP
DAM(K,MT)=DIR(K,MT)*(TRF/TF)**(ETA(K,MT)/2.)/DMR
CN8(K,MT)=DDN/DAM(K,MT)*1.414214
BT=AMIN1(BTA(K),BTA(MT))
VR1=S+3.*(1.+SQRT(TR))/BT
VR2=3.*SQRT((1.+2.*S**2/(5.+CHIM))*(1./WTM(K)+1./WTM(MT)))
CM(K,MT,1)=AMAX1(VR1,VR2)
CN(K,MT,1)=RAND(0)*CM(K,MT,1)
DF=PHI(K,MT)*(CHI(K)+CHI(MT)+2.)-1
DS=PHI(K,MT)*(2.-.5*ETA(K,MT))-1.0
DO 917 N=2,3
XPM=ACR**AMIN1(DF,DS)
IF ((DF.GT.0.).AND.(DS.GT.0.)) XPM=(DF/(DF+DS))**DF*(DS/(DF+DS))**DS
XPN=ACR**AMAX1(DF,DS)
IF ((DF.LT.0.).AND.(DS.LT.0.)) XPN=(DF/(DF+DS))**DF*(DS/(DF+DS))**DS
CM(K,MT,N)=XPM-XPN
CN(K,MT,N)=RAND(0)*CM(K,MT,N)
DF=CHI(K)
DS=CHI(MT)
917 CONTINUE
117 CONTINUE
ARG=SN(MT)
DO 119 NT=1,2
D=ERRF(ARG)
TENPA=EXP(-SN(MT)*SN(MT))/3.544908+0.5*ARG*D
TENPA=TENPA/BTA(MT)
TEMPC=0
DO 118 N=1,6

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ORIGINAL PAGE IS
OF POOR QUALITY

RUN2460
RUN2470
RUN2480
RUN2490
RUN2500
RUN2510
RUN2520
RUN2530
RUN2540
RUN2550
RUN2560
RUN2570
RUN2580
RUN2590
RUN2600
RUN2610
RUN2620
RUN2630

RUN2730
RUN2740
RUN2750

RUN2760
RUN2770
RUN2780
RUN2790
RUN2800
RUN2810
RUN2820

TEMPB=RLD(N)*RLD(N)	RUN2830
TEMPD=INM*TEMPA*DTM/(XR*RM*RM)*(TEMPB-TEMPC)/LWF(N)	RUN2840
DO 116 K=1,NWEDG	RUN2850
DELTH=DELANG(1)	RUN2860
IF(K.GT.NWEDGE(1)) DELTH=DELANG(2)	RUN2870
ENT(NT,MT,N,K)=TEMPD*DELTH/180.	RUN2880
116 REM(NT,MT,N,K)=0.	RUN2890
TEMPC=TEMPB	RUN2900
118 CONTINUE	RUN2910
SSR=ARG*ARG	RUN2920
SSA(NT,MT)=ARG+SQRT(SSR+2.)	RUN2930
SSB(NT,MT)=SSA(NT,MT)*(.25*SSA(NT,MT)-ARG)	RUN2940
119 ARG=-ARG	RUN2950
CHT=CHI(MT)	
IF(CHT.GT.0.) CMG(MT)=CHT**CHT*EXP(-CHT)	
IF(CHT.EQ.0.) CMG(MT)=1.0	
IF(CHT.LT.0.) CMG(MT)=ACR**CHT*EXP(-ACR)	
CNG(MT)=RAND(0)*CMG(MT)	
ANG=0.	RUN2970
N=0	RUN2960
ARG=ST(MT)	RUN2980
TEMPD=INM/PI*DTM/RM*2./(BTA(MT)*LWF(6))	RUN2990
DO 135 M=1,2	
I=NWEDGE(M)	RUN3010
IF(I.EQ.0) GO TO 135	
DO 134 K=1,I	RUN3020
N=N+1	RUN3030
AA=ANG	RUN3040
BB=DELANG(M)	RUN3050
CC=AA+BB	RUN3060
ANG=ANG+BB	RUN3070
THETA(N)=AA	RUN3080
DTH(N)=BB	RUN3090
AA=AA*PI/180.	RUN3100
CC=CC*PI/180.	RUN3110
ARGA=ARG*COS(AA)	RUN3120
ARGC=ARG*COB(CC)	RUN3130
ENTS(MT,N)=0.	RUN3140
REMS(MT,N)=0.	RUN3150
PTH(MT,N)=0.	RUN3160
TEMPA=0.	RUN3170
IF(ARGA.LE.-10.) GO TO 120	RUN3180
TEMPA=FNCTM(ARGA,PIROOT,L,COEFF)	RUN3190
PTH(MT,N)=2.*PIROOT*TEMPA	RUN3200
120 TEMPB=0.	RUN3210
IF(ARGC.LE.-10.) GO TO 125	RUN3220
TEMPB=FNCTM(ARGC,PIROOT,L,COEFF)	RUN3230
125 SUM1=TEMPA+TEMPB	RUN3240
TEMPC=.5*SUM1*(CC-AA)*TEMPD	RUN3250
IF(TEMPC.LT.1.E-06) GO TO 134	RUN3260
CALL SIMPSN(ARGA,ARGC,0,INTGRL,PERCNT,COEFF,PIROOT,SUM1,FNCTM)	RUN3270
ENTS(MT,N)=INTGRL*TEMPD*(CC-AA)	RUN3280
134 CONTINUE	RUN3290
135 CONTINUE	
140 CONTINUE	RUN3330

ORIGINAL PAGE IS
OF POOR QUALITY

XS(1)=.5*(XCB(1)-XSTART)*AKN	RUN3370
DO 155 N=2,ND	RUN3380
155 XS(N)=(.5*(XCB(N)+XCB(N-1))-XSTART)*AKN	RUN3390
IF(NS.EQ.0) GO TO 160	RUN3400
DO 159 I=1,NS	RUN3410
N=MS(I)	RUN3420
X=XS(N)/AKN+XSTART	RUN3430
J=0	RUN3440
157 J=J+1	RUN3450
IF(X.GT.XLIM(J+2)) GO TO 157	RUN3460
CALL HEIGHT(X,Y,J,COEFF,3)	RUN3470
TANGN(I)=(IWS(I)-.5)*DELANG(1)	RUN3480
IF(IWS(I)*DELANG(1).GT.THETAZ) TANGN(I)=THETAZ+(IWS(I)-NWEDGE(1)-.15)*DELANG(2)	RUN3490
Z=Y*SIN(TANGN(I)*PI/180.)	RUN3500
Y=-Y*COS(TANGN(I)*PI/180.)	RUN3510
CALL NORMAL(EYE,JAY,KAY,ONE,COEFF)	RUN3520
SNN=-S*(COSANG*EYE+SINANG*JAY)	RUN3530
ST1=S*(COSANG*ONE-EYE*SINANG*JAY/ONE)	RUN3540
ST2=-S*SINANG*KAY/ONE	RUN3550
DO 159 MT=1,MSP	RUN3560
VL(MT,1)=AMAX1(0.,SNN-VEL(1)/BTA(MT))	RUN3580
VL(MT,2)=ST1-VEL(2)/BTA(MT)	RUN3590
VL(MT,3)=ST2-VEL(3)/BTA(MT)	RUN3600
DELV(MT,1)=SNN+VEL(1)/BTA(MT)-VL(MT,1)	RUN3610
DELV(MT,2)=2.*VEL(2)/BTA(MT)	RUN3620
DELV(MT,3)=2.*VEL(3)/BTA(MT)	RUN3630
AMJ=MJ-1	RUN3640
DO 159 K=1,3	RUN3650
DO 159 J=1,MJ	RUN3660
FV(MT,I,1,J,K)=VL(MT,K)+(J-1)/AMJ*DELV(MT,K)	RUN3670
PV(MT,I,2,J,K)=SL(MT,I,K)+(J-1)/AMJ*DELS(MT,I,K)	RUN3680
159 CONTINUE	RUN3690
160 CONTINUE	RUN3700
IF(YR.GT.0.) GO TO 169	RUN3710
A=XSTART	RUN3720
B=XCB(1)	RUN3730
L=1	RUN3740
SUM1=FNCTN(A,PIROOT,L,COEFF)+FNCTN(B,PIROOT,L,COEFF)	RUN3750
CALL SIMPSN(A,B,L,INTGRL,PERCNT,COEFF,PIROOT,SUM1,FNCTN)	RUN3760
INTGRL=INTGRL/(ABS(COEFF(2,L))*RMB*RMB)*(B-A)	RUN3770
YCB(1)=INTGRL	RUN3780
DO 168 N=2,ND	RUN3790
A=XCB(N-1)	RUN3800
B=XCB(N)	RUN3810
IF(XLIM(L+2).GE.B) GO TO 167	RUN3820
L=L+1	RUN3830
167 SUM1=FNCTN(A,PIROOT,L,COEFF)+FNCTN(B,PIROOT,L,COEFF)	RUN3840
CALL SIMPSN(A,B,L,INTGRL,PERCNT,COEFF,PIROOT,SUM1,FNCTN)	RUN3850
INTGRL=INTGRL/(ABS(COEFF(2,L))*RMB*RMB)*(B-A)	RUN3860
168 YCB(N)=INTGRL	RUN3870
169 LV=2	
IF(NWEDGE(2).EQ.0) LV=1	
CALL CELL(THETAZ,LV,BW,BH,NW,NH,0.,0,0,DELANG,NWEDGE,XC,YC,ZC,FNB)	
IF(NL.LT.2) GO TO 170	RUN3890

ORIGINAL PAGE IS
OF POOR QUALITY

CALL ZERO (NW, NH, NFA, NCA, NHA, 0, NWEDGE (1), FNB)	RUN3900
XLB=BW*NFA	RUN3910
BWB=BW*NCA/MW	RUN3920
BHB=BH*NHA/MH	RUN3930
CALL CELL (THETAZ, 1, BWB, BHB, MW, MH, XLB, NXA, 0, DELANG, NWEDGE, XC, YC, ZC, 1 FNB)	RUN3940
IF (NL.LT.3) GO TO 170	RUN3950
CALL ZERO (MW, MH, NFB, NCB, NHB, NXA, NWEDGE (1), FNB)	RUN3960
XLC=XLB+BWB*NFB	RUN3970
BWC=BWB*NCB/LW	RUN3980
BHC=BHB*NHB/LH	RUN3990
CALL CELL (THETAZ, 1, BWC, BHC, LW, LH, XLC, NXA, NXB, DELANG, NWEDGE, XC, YC, ZC, 1C, FNB)	RUN4000
170 CALL SBTRCT (1, IXA, NXA, BW, BH, DELANG, XC, YC, FNB, XLIM, COEFF)	RUN4010
LEV (1)=NBX+1	RUN4020
LEV (2)=NBX+1	RUN4030
IF (NL.LT.2) GO TO 190	RUN4040
NI=NXA+1	RUN4050
NF=NXA+NXB	RUN4060
LEV (1)=NI	RUN4070
CALL SBTRCT (NI, NF, NF, BWB, BHB, DELANG, XC, YC, FNB, XLIM, COEFF)	RUN4080
IF (NL.LT.3) GO TO 190	RUN4090
NI=NF+1	RUN4100
NF=NBX	RUN4110
LEV (2)=NI	RUN4120
CALL SBTRCT (NI, NF, NF, BWC, BHC, DELANG, XC, YC, FNB, XLIM, COEFF)	RUN4130
190 FNA=0.0	RUN4140
M=0	RUN4150
DO 210 N=1, NBX	RUN4160
NUMCEL (N)=0	RUN4170
IF (FNB (N).LE.0.) GO TO 210	RUN4180
M=M+1	RUN4190
NUMCEL (N)=M	RUN4200
FNA=FNA+FNB (N)	RUN4210
DYC=BH/2.	
IF (N.GT.NXA) DYC=BHB/2.	
IF (N.GT.NXA+NXB) DYC=BHC/2.	
YTC=YC (N)+DYC	
DO 200 LA=1, 6	
IF (YTC.LE.RLD (LA)) GO TO 201	
200 CONTINUE	
201 LKW (M)=LWP (LA)	
210 CONTINUE	
NPX=M	RUN4220
IF (NPX.GT.LIMIT (4)) CALL DIAG (4, LIMIT (4), NPX)	RUN4230
220 TIME=0	RUN4240
LARGE=0	RUN4250
SAMP=0	RUN4260
PRT=0	RUN4270
NAV=0	RUN4290
AIME=0.	RUN4300
TI=-1.	RUN4310
NMAX=0	RUN4380
DO 250 MT=1, 3	
C1 (MT)=RAND (0)	RUN4410

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OF POOR QUALITY

ORIGINAL PAGE IS
OF POOR QUALITY

C2(MT)=RAND(0)	RUN4420
C3(MT)=RAND(0)	RUN4430
C7(MT)=RAND(0)	RUN4440
C8(MT)=RAND(0)	RUN4450
D1(MT)=RAND(0)	RUN4460
D2(MT)=RAND(0)	RUN4470
D3(MT)=RAND(0)	RUN4480
D4(MT)=RAND(0)	RUN4490
FL(MT)=0.	RUN4500
HTL(MT)=0.	RUN4510
HTR(MT)=0.	RUN4520
JNT(MT)=0	RUN4530
NM(MT)=0	RUN4540
DO 230 N=1,3	RUN4550
CTI(MT,N)=0.	RUN4560
CTR(MT,N)=0.	RUN4570
CNI(MT,N)=0.	RUN4580
230 CNR(MT,N)=0.	RUN4590
DO 240 N=1,ND	RUN4600
DO 240 K=1,NWEDG	RUN4610
NTS(MT,N,K)=0	RUN4620
NTSF(MT,N,K)=0	RUN4630
HTSI(MT,N,K)=0.	RUN4640
UTLI(MT,N,K)=0.	
UTTI(MT,N,K)=0.	
VTSI(MT,N,K)=0.	
UTL(MT,N,K)=0.	RUN4650
UTT(MT,N,K)=0.	RUN4660
VTS(MT,N,K)=0.	RUN4670
240 HTS(MT,N,K)=0.	RUN4680
DO 250 I=1,3	
NTCP(MT,I)=0	RUN4820
DO 250 L=1,2	RUN4830
DO 250 K=1,3	RUN4840
DO 250 J=1,MJ	RUN4850
NTCV(MT,I,L,J,K)=0	RUN4860
250 CONTINUE	RUN4870
DO 245 N=1,NPX	RUN4690
DO 245 MT=1,NSP	
NB(MT,N)=0	RUN4700
NBP(MT,N)=0	RUN4710
NBS(MT,N)=0	
NBT(MT,N)=0	RUN4720
DBA(MT,N)=0.	RUN4730
XVA(MT,N)=0.	RUN4740
YVA(MT,N)=0.	RUN4750
ZVA(MT,N)=0.	RUN4760
TMPA(MT,N)=0.	RUN4790
TRPA(MT,N)=0.0	
DO 245 NN=1,NSP	
T(MT,NN,N)=0.0	
FND=DDN	
245 CONTINUE	RUN4800
DRF=2./((FND*S*S*RMB*RMB*PI)	RUN4930
FCF=1./((FND*S*RMB*RMB*PI)	RUN4940

ORIGINAL PAGE 11
OF POOR QUALITY

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HTF=.5*DRP/S
C9=Rand(0)*RWPM
LL=FNA/VOL*INM
WRITE(6,2)
WRITE(6,4)
WRITE(6,5) ((I,J,L,(ENT(I,J,K,L),K=1,6),L=1,NWEDG),J=1,MSP),I=1,2)
WRITE(6,6) ((J,L,ENTS(J,L),PTH(J,L),L=1,NWEDG),J=1,MSP)
WRITE(6,2)
CALL PRINTA(THETAZ,NWEDGE,TITLE,NAME,XCB,YCB,TB,ALPHA,SIGMA,LD,LP,
1XLIM,COEFF,LIMIT,MSP)
CALL GAS(NWEDG,THETAZ,DELANG,NWEDGE,BTA,C1,DFA,NM,RLD,LWF,FNB,DB,N
1B,NBF,LPF,PAU,PAV,PAW,PAX,PAY,PAZ,XLIM,COEFF,LM,LIMIT(4),LIMIT(6),
2LARGE,MNM,MNB,DEBUG(1),LCOL,NUMCEL,MSP,ER,CHI,CNG,CMG,NSP,LB,NBM,
3NBN)
CPUTYM=TFIND(0)
IF(LARGE.NE.0) GO TO 345
DO 265 I=1,MSP
265 IF(NM(I).GT.NMAX) NMAX=NM(I)
CALL PRINTE(FNA,MSP,FNB,LEV,LWF,NM,RLD,XLIM,XC,YC,ZC,NB,NUMCEL,LKW
1,NSP)
IF(DEBUG(2)) WRITE(6,1)
CALL ACCUM(NMC,NPB,FNB,NB,PAU,PAV,PAW,ER,TMP,TRP,XV,YV,ZV,LM,MSP,
1NSP,LPF,NBF,NBM)
CPA=ELTIME(0)
CPI=CPA
GO TO 340
280 TIME=TIME+1
285 LARGE=0
CPI=ELTIME(0)
AIME=TIME*DTM
IF(DEBUG(1)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA,(NM(I),I=1,3),NMAX
PRT=PRT+1
SAMP=SAMP+1
CALL COLIDE(CN,CH,WTN,DB,DBA,NB,NCOL,LCOL,PAU,PAV,PAW,ER,T,LM,MSP,
1LIMIT(4),LIMIT(6),NUMCEL,ETA,PHI,CHI,CN8,NSP,LPF,LKW,NBF,NBM)
KNM(1)=0
KNM(2)=0
KNM(3)=0
CPC=ELTIME(0)
IF(DEBUG(1)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA,(NM(I),I=1,3),NMAX
CALL MOVE(0,AKN,MJ,NS,NWEDG,THETAZ,XSTART,LIMIT(3),LIMIT(1),LIMIT
1(8),LIMIT(9),DELANG,NWEDGE,BTA,C2,C3,DFA,PL,HTI,HTR,JNT,KNM,NM,XCB
2,XLIM,MS,IWS,NTCF,NTCV,FV,CTI,CTR,CNI,CNR,ALPHA,SIGMA,COEFF,HTS,HTR
3SI,NTS,NTSP,UTL,UTT,VTS,PAU,PAV,PAW,PAX,PAY,PAZ,LPF,LCOL,TB,MSP,ER
4,CHI,CNG,CMG,NSP,UTLI,UTTI,VTSTI)
KNM(1)=NM(1)
KNM(2)=NM(2)
KNM(3)=NM(3)
CPM=ELTIME(0)
IF(DEBUG(1)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA,(NM(I),I=1,3),NMAX
CALL FLOW(NWEDG,MNM,LARGE,BTA,C1,C7,C8,D1,D2,D3,D4,DTH,NM,SN,ST,THRU
1ETA,LWF,RLD,PTH,ENTS,REMS,SSA,SSB,PAU,PAV,PAW,PAX,PAY,PAZ,LPF,ENT,
2REM,LCOL,MSP,ER,CHI,CNG,CMG,NSP)
IF(LARGE.NE.0) GO TO 345
CPB=ELTIME(0)

```

RUN4960
RUN4970RUN5000
RUN5010

RUN5060

RUN5130
RUN5140RUN5170
RUN5180
RUN5190RUN5220
RUN5230RUN5280
RUN5290

RUN5330


```

IF (DEBUG (1)) WRITE (6, 33) AIME, CPC, CPM, CPI, CPB, CPA, (NM (I), I=1, 3), NMAX
CALL MOVE (1, AKN, MJ, NS, NWEDG, THETAZ, XSTART, LIMIT (3), LIMIT (1), LIMIT
1 (8), LIMIT (9), DELANG, NWEDGE, BTA, C2, C3, DFA, FL, HTI, HTR, JNT, KNM, NM, XCBRUN5350
2, XLIM, MS, IWS, NTCF, NTCV, FV, CTI, CTR, CNI, CNR, ALPHA, SIGMA, COEPP, HTS, HTRUN5360
3 SI, NTS, NTSP, UTL, UTT, VTS, PAU, PAV, PAW, PAX, PAY, PAZ, LPP, LCOL, TB, MSP, ER
4, CHI, CMG, CMG, NSP, UTLI, UTTI, VTSI)
DO 330 MT=1, MSP
CPB=CPB+ELTIME (0)
IF (DEBUG (1)) WRITE (6, 33) AIME, CPC, CPM, CPI, CPB, CPA, (NM (I), I=1, 3), NMAX
M=0
DO 290 N=1, NBX
IF (NUMCEL (N).EQ.0) GO TO 290
M=M+1
NB (MT, M)=0
NBF (MT, M)=0
290 CONTINUE
NG=NM (MT)
N=0
295 N=N+1
IF (N.GT.NG) GO TO 310
X=PAX (MT, N)
Y=PAY (MT, N)
Z=PAZ (MT, N)
R=SQRT (Y*Y+Z*Z)
ARG=Y/R
TANG=180.* (1.-ARCCOS (ARG)/PI)
IWDGE=TANG/DELANG (1)
IF ((IWDGE.GE.NWEDGE (1)).AND. (DELANG (2).NE.0.)) IWDGE=(TANG-THETAZ)
1/DELANG (2)+NWEDGE (1)
IF (IWDGE.LT.0) IWDGE=0
IF (IWDGE.GE.NWEDGE) IWDGE=NWEDGE-1
L=X/BW+1.
IF (L.GT.NW) L=NW
M=R/BH
IF (M.GE.NH) M=NH-1
K=(IWDGE*NH+M)*NW+L
IF (K.LE.NXA) GO TO 296
WRITE (6, 40) L, M, K, MT, N, X, Y, NWEDGE, NWEDG, TANG, Z, IWDGE, NH, NW
IF (DUMP) CALL ABEND (4)
STOP
296 KW=0
IF (NL.EQ.1) GO TO 300
IF (IWDGE.GE.NWEDGE (1)) GO TO 300
IF (PNB (K).GT.0.) GO TO 300
L=(X-XLB)/BWB+1.
IF (L.GT.MW) L=MW
M=P*BHB
IF (M.GE.MH) M=MH-1
K=(IWDGE*MH+M)*MW+L+NXA
IF (K.LE.NXA+NXB) GO TO 297
WRITE (6, 40) L, M, K, MT, N, X, Y, NWEDGE, NWEDG, TANG, Z, IWDGE, MH, MW, NXA
IF (DUMP) CALL ABEND (5)
STOP
297 IF (NL.EQ.2) GO TO 300
IF (PNB (K).GT.0.) GO TO 300

```

RUN5410

RUN5420

RUN5440

RUN5450

RUN5460

RUN5470

RUN5480

RUN5490

RUN5500

RUN5510

RUN5520

RUN5530

RUN5540

RUN5550

RUN5560

RUN5570

RUN5580

RUN5600

RUN5610

RUN5620

RUN5630

RUN5640

RUN5650

RUN5660

RUN5670

RUN5680

RUN5690

RUN5700

RUN5710

RUN5730

RUN5740

RUN5750

RUN5760

RUN5770

RUN5780

RUN5790

RUN5800

RUN5810

RUN5820

RUN5830

RUN5840

ORIGINAL PAGE IS
OF POOR QUALITY

ORIGINAL PROGRAM
OF POOR QUALITY

L=(X-XLC)/BWC+1.	RUN5850
IF (L.GT.LW) L=LW	RUN5860
M=R/BHC	RUN5870
IF (M.GE.LH) M=LH-1	RUN5880
K=(IWDGE*LN+M)*LW+L+NXA+NXB	RUN5890
IF (K.LE.NBX) GO TO 300	RUN5900
WRITE(6,40) L,M,K,MT,N,X,Y,NWEDGE,NWEDG,TANG,Z,IWDGE,LH,LW,NXA,NXB	RUN5910
IF (DUMP) CALL ABEND(6)	RUN5920
STOP	RUN5930
300 LA=0	RUN5940
IF (NUMCEL(K).EQ.0) GO TO 306	
301 LA=LA+1	RUN5960
IF (R.GT.RLD(LA)) GO TO 301	RUN5970
KW=LWF(LA)	RUN5980
KKW=LPF(MT,N)	RUN5990
IF (KW.EQ.KKW) GO TO 305	RUN6000
IF (KW.LT.KKW) GO TO 302	RUN6010
M=KW/KKW	RUN6020
A=RAND(0)	RUN6030
B=M	RUN6040
B=1./B	RUN6050
IF (A.LT.B) GO TO 305	RUN6060
306 PAX(MT,N)=PAX(MT,NG)	RUN6070
PAY(MT,N)=PAY(MT,NG)	RUN6080
PAZ(MT,N)=PAZ(MT,NG)	RUN6090
PAU(MT,N)=PAU(MT,NG)	RUN6100
PAV(MT,N)=PAV(MT,NG)	RUN6110
PAW(MT,N)=PAW(MT,NG)	RUN6120
ER(MT,N)=ER(MT,NG)	
LPF(MT,N)=LPF(MT,NG)	RUN6130
LCOL(MT,N)=LCOL(MT,NG)	RUN6140
N=N-1	RUN6150
NM(MT)=NM(MT)-1	RUN6160
NG=NM(MT)	RUN6170
GO TO 295	RUN6180
302 M=KKW/KW-1	RUN6190
IF ((NM(MT)+M).LE.MNM) GO TO 307	
LARGE=3	RUN6230
GO TO 345	RUN6240
307 CONTINUE	
DO 304 L=1,M	RUN6200
NM(MT)=NM(MT)+1	RUN6210
NG=NM(MT)	
PAX(MT,NG)=PAX(MT,N)	RUN6260
PAY(MT,NG)=PAY(MT,N)	RUN6270
PAZ(MT,NG)=PAZ(MT,N)	RUN6280
PAU(MT,NG)=PAU(MT,N)	RUN6290
PAV(MT,NG)=PAV(MT,N)	RUN6300
PAW(MT,NG)=PAW(MT,N)	RUN6310
ER(MT,NG)=ER(MT,N)	
LCOL(MT,NG)=LCOL(MT,N)	RUN6320
304 LPF(MT,NG)=KW	RUN6330
305 LPF(MT,N)=KW	RUN6340
Q=NUMCEL(K)	RUN6350
J=NB(MT,Q)+1	RUN6360

```

IF (J.LE.MNB) GO TO 308
IF (DEBUG(1)) WRITE(6,44) MT,K,MNB,AIME
308 NB(MT,Q)=J
NBF(MT,Q)=NBF(MT,Q)+KW
LB(N)=Q
GO TO 295
310 CONTINUE
NBM(MT,1)=0
DO 320 N=1,NBX
M=NUMCEL(N)
IF (M.EQ.0) GO TO 320
A=NBF(MT,M)
DB(MT,M)=A*DFA(MT)/FNB(N)
NBM(MT,M+1)=NBM(MT,M)+NB(MT,M)
NBN(M)=NBM(MT,M)
320 CONTINUE
IF (NM(MT).GT.NMAX) NMAX=NM(MT)
DO 325 N=1,NG
Q=LB(N)
NBN(Q)=NBN(Q)+1
NA=NBN(Q)
325 LM(MT,NA)=N
330 CONTINUE
IF (SAMP.LT.ITS) GO TO 335
CALL ACCUM(NMC,NPB,FNB,NB,PAU,PAV,PAW,ER,TMP,TRP,XV,YV,ZV,LM,MSP,
1 NSP,LPF,NBF,NBM)
SAMP=0
IF (TIME.LE.TST) GO TO 335
CALL AVRGE(FNB,DB,DBA,NB,NBT,XV,YV,ZV,XVA,YVA,ZVA,TMP,TMPA,TRP,TRP
1A,MSP,NSP,NBF,NBS)
NAV=NAV+1
335 CPA=ELTIME(0)
CPI=CPC+CPM+CPB+CPA
CPJ=2.*CPI+5.
340 CPUTYM=TFIND(0)
IF (DEBUG(2)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA,(NM(I),I=1,3),NMAX
IF ((TIME.GE.TLIM).OR.(CPUTYM.LE.CPJ)) GO TO 345
IF (PRT.LT.ITP) GO TO 280
PRT=0
345 WRITE(6,30) AIME,KAWLS
IF (DEBUG(3)) WRITE(6,31) CPUTYM
WRITE(6,32) (NM(I),I=1,3)
WRITE(6,34) ((NCOL(I,J),J=1,3),I=1,3)
WRITE(6,35) (JNT(I),I=1,3)
IF (LARGE.NE.0) GO TO 360
WRITE(6,36) NMAX
IF (.NOT.SAVE) GO TO 355
IF (PRT.NE.0.AND.CPUTYM.GT.CPJ.AND.TIME.LT.TLIM) GO TO 355
WRITE(9) DENF,U,XREF,TRF, KAWLS,NL,NW,NH,MW,MH,LW,LH,NXA,NXB,NCA
1 ,NCB,NFA,NFB,NHA,NHB,BW,BH,BWB,BHB,BWC,BHC,XLB,XLC,PI,NREG
2 ,S,SINANG,COSANG,AKN,NBX,RM,XR,ND,TIME,DTM,TI,ITS,ITP,TST
3 ,TLIM,RMA,RNU,DIR,XSTART,JNM,MNM,MNB,TR,BZC,CN7,DRF,FCF
4 ,FNA,HTF,INM,ITYPE,JTYPE,MJ,NAV,NMAX,NS,NWEDG,PRT,SAMP
5 ,BTA,C1,C2,C3,C7,C8,DAM,DFA,FL,DELANG,FDN,HTI,HTR,JNT,KNM
6 ,NM,WTM,C4,VRM,NCOL,LD,LF,LWF,RLD,CTI,CTR,CNI,CNR,LEV,SN

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RUN6370

RUN6400

RUN6410

RUN6430

ORIGINAL PAGE 13
OF POOR QUALITY

RUN6450

RUN6480

RUN6490

RUN6500

RUN6510

RUN6520

RUN6570

RUN6600

RUN6650

RUN6660

RUN6670

RUN6740

RUN6750

RUN6900

RUN6930

RUN6940

RUN6950

RUN6960

RUN6970

RUN6980

```

7      ,ST,D1,D2,D3,D4,SSA,SSB,MS,NSP,NMB,NMC,NMP,NPB,NRAN,VELR
8      ,IWS,TANGN,XLIM,COEFF,XCB,XS,YCB,TB,ALPHA,SIGMA,NTS,NTSF RUN7000
9      ,UTL,UTT,VTS,HTS,HTSI,ENT,REM,ENTS,REMS,PTH,THETA,DTH,TMPARUN7010
A      ,DBA,NB,NBP,NBT,TMP,XV,XVA,YV,YVA,ZV,ZVA,T,DB,FNB,XC,YC,ZCRUN7020
B      ,NUMCEL,PAU,PAV,PAW,PAX,PAY,PAZ,PV,NTCV,NTCF,LPP,LCOL,LM RUN7030
C,ETA,PHI,CHI,CN,CN,CNG,CMG,CN8,TRP,TRPA,THETAZ,NWEDGE,MSP,ANGLE,TF
D,UTLI,UTTI,VTST,ER,RNB,LKW,NBS,LB,NBN,NBN
REWIND 9 RUN7040
WRITE(6,50) RUN7050
355 CONTINUE
IF (TIME.LE.TST) GO TO 350
DT=AIME-TI RUN6770
CALL PRINT1(DT,COSANG,SINANG,RMA,RNU,DRF,PCF,HTP,PL,HTI,HTR,CTI,
1CTR,CNI,CNR)
CALL PRINT2(AKN,XSTART,DT,RNU,RMA,DRF,PCF,HTP,UTLI,UTTI,VTST,HTSI,
1DELANG,NWEDGE,XS,XCB,YCB,HTS,NTS,NTSF,UTL,UTT,VTS,LIMIT(3),
2LIMIT(1),MSP)
IF (NS.NE.0) CALL PRINT3(MSP,MJ,NS,NWEDG,LIMIT(3),LIMIT(1),
1LIMIT(8),LIMIT(9),RMA,XS,IWS,MS,TANGN,NTSF,NTCF,NTCV,PV)
CALL PRINT4(MSP,CHI,RNU,NSP,TRPA,NUMCEL,FDN,WTM,DBA,NBS,TMPA,XVA,
1YVA,ZVA,1,NBT,XC,YC,ZC,LEV,LKW)
GO TO 353 RUN6860
350 CONTINUE
CALL PRINT4(MSP,CHI,RNU,NSP,TRP,NUMCEL,FDN,WTM,DB,NB,TMP,XV,YV,ZV,
10,NBP,XC,YC,ZC,LEV,LKW)
353 IF (DEBUG(2)) WRITE(6,1)
IF ((TIME.LT.TLIM).AND.(CPUTYM.GT.CPJ)) GO TO 280
IF (IC.EQ.ICOPY) RETURN
IC=IC+1 RUN7080
WRITE(6,2)
WRITE(6,4)
WRITE(6,5) ((I,J,L,(ENT(I,J,K,L),K=1,6),L=1,NWEDG),J=1,MSP),I=1,2)
WRITE(6,6) ((J,L,ENTS(J,L),PTH(J,L),L=1,NWEDG),J=1,MSP)
WRITE(6,2) RUN7120
WRITE(6,3) IC RUN7130
CALL PRINTA(THETAZ,NWEDGE,TITLE,NAME,XCB,YCB,TB,ALPHA,SIGMA,LD,LF,RUN7140
1XLIM,COEFF,LIMIT,MSP)
CALL PRINTB(FNA,MSP,FNB,LEV,LWF,NM,RLD,XLIM,XC,YC,ZC,NB,NUMCEL,LKW
1,NSP)
SAVE=.FALSE. RUN7180
GO TO 345 RUN7190
360 WRITE(6,38) (DBG1(I,LARGE),I=1,3) RUN7200
IF (REDO) GO TO 364
IF (DUMP) CALL ABEND(9)
STOP RUN7220
364 CONTINUE RUN7230
IF (NEW) GO TO 365
READ(9) DENF,U,XREF,TRF, KAWLS,NL,NW,NH,MW,MH,LW,LH,NXA,NXB,NCA
1,NCB,NFA,NFB,NHA,NHB,BW,BH,BWB,BHB,BWC,BHC,XLB,XLC,PI,NREG
2,S,SINANG,COSANG,AKN,NBX,RM,XR,ND,TIME,DTM,TI,ITS,ITP,TST
3,TLIM,RMA,RNU,DIR,XSTART,JNM,MNM,MNB,TR,BZC,CN7,DRF,PCF
4,FNA,HTP,INM,ITYPE,MJ,NAV,NMAX,NS,NWEDG,PRT,SAMP
5,BTA,C1,C2,C3,C7,C8,DAM,DFA,PL,DELANG,FDN,HTI,HTR,JNT,KNM
6,NM,WTM,C4,VRM,NCOL,LD,LF,LWF,RLD,CTI,CTR,CNI,CNR,LEV,SN
7,ST,D1,D2,D3,D4,SSA,SSB,MS,NSP,NMB,NMC,NMP,NPB,NRAN,VELR

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8      ,IWS,TANGN,XLIM,COEFF,XCB,XS,YCB,TB,ALPHA,SIGMA,NTS,NTSF
9      ,UTL,UTT,VTS,HTS,HTSI,ENT,REM,ENTS,REMS,PTH,THETA,DTH,TMPA
A      ,DBA,NB,NBF,NBT,TMP,XV,XVA,YV,YVA,ZV,ZVA,T,DB,FNB,XC,YC,ZC
B      ,NUMCEL,PAU,PAV,PAW,PAX,PAY,PAZ,PV,NTCV,NTCF,LPP,LCOL,LN
C,ETA,PHI,CHI,CN,CN,CNG,CMG,CN8,TRP,TRPA,THETAZ,NWEDGE,MSP,ANGLE,TF
D,UTLI,UTTI,VTST,ZR,RMB,LKW,NBS,LB,NBM,NBN

```

REWIND 9

365 JNM=9*JNM/10

RUN7260

ANM=INM

INM=9*ANM/10

DDN=.9*DDN

DRF=DRF/.9

PCF=PCF/.9

HTF=HTF/.9

DO 370 MM=1,MSP

FDN(MM)=FDN(MM)*INM/ANM

DO 366 KK=1,MSP

366 CN8(KK,MM)=CN8(KK,MM)*.9

DO 370 LT=1,NWEDG

ENTS(MM,LT)=ENTS(MM,LT)*INM/ANM

REMS(MM,LT)=0.0

DO 370 NK=1,2

DO 370 NJ=1,6

ENT(NK,MM,NJ,LT)=ENT(NK,MM,NJ,LT)*INM/ANM

370 REM(NK,MM,NJ,LT)=0.0

IF(NEW) GO TO 220

TST=TIME+TST

TI=-1.

PRT=ITP

WRITE(6,2)

WRITE(6,4)

WRITE(6,5) ((I,J,L,(ENT(I,J,K,L),K=1,6),L=1,NWEDG),J=1,MSP),I=1,2)

WRITE(6,6) ((J,L,ENTS(J,L),PTH(J,L),L=1,NWEDG),J=1,MSP)

WRITE(6,2)

IF((LARGE.EQ.2).OR.(LARGE.EQ.3)) GO TO 280

REDO=.FALSE.

GO TO 360

END

RUN7410

SUBROUTINE DIAG(N,ITEST,NUM)

DIAG010

REAL*8 PARAM(10) NWEDGE', '

NREG', '

ND', '

NPX', '

1 MNM', '

MNB', '

NBX', '

NS', '

MJ', '

MSP'/

DIAG040

DIAG050

DIAG060

DIAG070

DIAG080

DIAG090

DIAG100

DIAG110

DIAG120

DIAG130

DIAG140

DIAG150

DIAG160

FORMATS

32 FORMAT(9X,'ENT,REM,ENTS,REMS,PTH,THETA,DTH')

42 FORMAT(///5X,43H ARRAY DIMENSIONS ARE ABOUT TO BE VIOLATED./)

44 FORMAT(5X,18H MAXIMUM VALUE IS 15,20H, WHEREAS YOU INPUT 15,3H

1A8,1H))

56 FORMAT(/5X,78H IF YOU DESIRE TO USE THIS VALUE, THE FOLLOWING ARRADIAG130

1YS MUST BE RE-DIMENSIONED./)

62 FORMAT(9X,'HTS,HTSI,NTS,NTSF,UTL,UTT,VTS')

64 FORMAT(9X,'XLIM,COEFF'//11X,'NOTE THAT THE XLIM ARRAY MUST BE DINEDIAG160

```

1 NSIONED TO 3 MORE THAN THE COEFF ARRAY.')
66 FORMAT(9X,'XCB,XS,YCB,TB,ALPHA,SIGMA')
68 FORMAT(9X,'DBA,NB,NBF,NBT,TMP,TMPA,XV,XVA,YV,YVA,ZV,ZVA,T,DB')
70 FORMAT(9X,'LPP,PAU,PAV,PAW,PAX,PAY,PAZ,LCOL')
72 FORMAT(8X,3H LM)
74 FORMAT(//5X,76H IF YOU CHANGE THE ARRAY DIMENSIONS, ALSO CHANGE THDIAG220
1E 'LIMIT' DATA STATEMENT.)
75 FORMAT(9X,'ALL ARRAYS ASSOCIATED WITH SPECIES')
76 FORMAT(9X,'PNB,XC,YC,ZC,NUMCEL')
78 FORMAT(9X,'FV,NTCV,NTCP,MS,IWS,SL,DELS,TANGN')
80 FORMAT(9X,'FV,NTCV')

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DIAG170
DIAG180
DIAG190
DIAG200
DIAG210
DIAG220
DIAG230
DIAG240
DIAG250
DIAG260
DIAG270

```

WRITE(6,42)
WRITE(6,44) ITEST,NUM,PARAM(N)
WRITE(6,56)
GO TO (1,2,3,4,5,6,7,8,9,10),N
1 WRITE(6,62)
WRITE(6,32)
GO TO 11
2 WRITE(6,64)
GO TO 11
3 WRITE(6,66)
WRITE(6,62)
GO TO 11
4 WRITE(6,68)
GO TO 11
5 WRITE(6,70)
GO TO 11
6 WRITE(6,72)
GO TO 11
7 WRITE(6,76)
GO TO 11
8 WRITE(6,78)
GO TO 11
9 WRITE(6,80)
GO TO 11
10 WRITE(6,75)
11 WRITE(6,74)
STOP
END

```

DIAG280
DIAG290
DIAG300
DIAG310
DIAG320

DIAG340
DIAG350

DIAG370

DIAG390
DIAG400

DIAG420

DIAG440

DIAG460

DIAG480

DIAG500

DIAG520

DIAG540

DIAG550

ORIGINAL PAGE
OF POOR QUALITY

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SUBROUTINE PRINTA(THETAZ,NWEDGE,TITLE,NAME,XCB,YCB,TB,ALPHA,SIGMA,PRA0010
1LD,LF,XLIM,COEFF,LIMIT,MSP)
INTEGER TST,TLM,TIME
LOGICAL SAVE,NEW
DIMENSION NWEDGE(2),LIMIT(1),TITLE(6),NAME(2),XCB(1),YCB(1),TB(1)
DIMENSION ALPHA(3,1),SIGMA(3,1),LD(1),LF(1),XLIM(1),COEFF(4,1)
DIMENSION RNU(3),RMA(3),CHI(3),DIR(3,3),PHI(3,3),ETA(3,3)
DIMENSION WTM(3),DAM(3,3),VELS(3),XSP(3)
COMMON /FIRST/NL,NW,NH,MW,MH,LW,LH,NXA,NXB,NCA,NCB,NPA,NPB,NHA,NHB
COMMON /SECND/BW,BH
COMMON /THIRD/PI,NREG,S,SINANG,COSANG,AKN
COMMON /FIFTH/ND,TIME,DTM,TI,ITS,ITP,TST,TLM,RMA,RNU,DIE

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PRA0030

PRA0040

PRA0060

PRA0070

PRA0080

PRA0090

PRA0100

COMMON /SIXTH/RMB, XSTART, JNN, MNN, MNB, NEW, SAVE, PERCNT, NSR, TR
COMMON/EIGHT/DENF, U, TF, ANGLE, TRF, CHI, PHI, ETA, WTH, DAM, VELR, XREF
DATA NOT/'NOT '/

PRA0110

PRA0120

PRA0130

PRA0140

PRA0150

PRA0160

PRA0170

FORMATS

ORIGINAL PAGE 10
OF POOR QUALITY

1 FORMAT(16X,40(' '),T74,'I'//9X,'3-D',I2,'-FLUID PROGRAM - ')
2 FORMAT('+',31X,A4)
3 FORMAT('+',35X, 'A RESTART OF A PREVIOUS RUN',T74,'I'/12X,2
1A4,' - ',6A4,' - ',I2,' REGIONS',T74,'I',16(/T74,'I'))
4 FORMAT(7X,'FRONT OF BODY =' ,E12.4,' XSTART MAX HEIGHT =' ,E12.4,'
1RMB',T74,'I'/7X,'X-LIMIT',T37,'BODY COEFFICIENTS',T74,'I')
6 FORMAT(5F14.6,3X,'I')
10 FORMAT(1X,72(' '))
12 FORMAT(/14X,'PARAMETERS OF SEGMENTS FOR BODY COLLISIONS',T96,'I'/
18X,'X-COORD. TEMP. ALPHA1 ALPHA2 ALPHA3 SIGMA1 SIGMA2
2 SIGMA3 AREAS',T96,'I')
14 FORMAT(4X,E12.4,7F9.4,E12.4,T96,'I')
16 FORMAT(23X,'WEIGHTING FACTORS'/1X,10I6,T96,'I')
17 FORMAT(///25X,'ARRAY STORAGE USFD'/5X,I6,' *',10I6,T96,'I')
18 FORMAT(1H1/17X,'LENGTH OF CELL IN MEAN-FREE-PATHS = ',F12.4,' BW'
A,T76,'I'
1/17X,'HEIGHT OF CELL IN MEAN-FREE-PATHS = ',F12.4,' BH',T76,'I'
2/16X,'NUMBER OF L1 CELLS ALONG FLOW AXIS =' ,I13,' NW',T76,'I'
3/17X,'NUMBER OF L1 CELLS IN RADIAL DIR. =' ,I13,' NH',T76,'I'
4/21X,'NUMBER OF LEVELS OF CELL SIZE =' ,I13,' NL',T76,'I')
20 FORMAT(11X,'NUMBER OF L1 CELLS IN FRONT OF L2 CELLS =' ,I13,' NFA'
1,T76,'I'/15X,'NUMBER OF AXIAL SUBDIVIDED L1 CELLS =' ,I13,' NCA',T
276,'I'/14X,'NUMBER OF RADIAL SUBDIVIDED L1 CELLS =' ,I13,' NHA',T7
36,'I'/16X,'NUMBER OF L2 CELLS ALONG FLOW AXIS =' ,I13,' MW',T76,'I'
4/17X,'NUMBER OF L2 CELLS IN RADIAL DIR. =' ,I13,' MH',T76,'I')
22 FORMAT(11X,'NUMBER OF L2 CELLS IN FRONT OF L3 CELLS =' ,I13,' NFB'
1,T76,'I'/15X,'NUMBER OF AXIAL SUBDIVIDED L2 CELLS =' ,I13,' NCB',T
276,'I'/14X,'NUMBER OF RADIAL SUBDIVIDED L2 CELLS =' ,I13,' NHB',T7
36,'I'/16X,'NUMBER OF L3 CELLS ALONG FLOW AXIS =' ,I13,' LW',T76,'I'
4/17X,'NUMBER OF L3 CELLS IN RADIAL DIR. =' ,I13,' LH',T76,'I')
23 FORMAT(3X,'NUMBER OF AZIMUTHAL WEDGES IN LOWER ',I3,' DEGREES =' ,IPRA0470
113,' NWEDGE1 I'/3X,'NUMBER OF AZIMUTHAL WEDGES IN UPPER ',I3,' DE
2GREES =' ,I13,' NWEDGE2 I'/)
24 FORMAT(16X,'BASIC TIME INTERVAL FOR COLLISIONS =' ,E13.4,' DTM
1 I'/8X,'TIME INTERVAL FOR SAMPLING FLOW FIELD INFO =' ,E13.4,' DTS
2 I'/24X,'TIME INTERVAL FOR PRINTING =' ,E13.4,' DTP I'/9X,
3'TIME TO STEADY-STATE CONDITIONS (ASSUMED) =' ,E13.4,' TST I'/
419X,'TIME AT WHICH RUN IS TERMINATED =' ,E13.4,' TLIM I'/)
26 FORMAT(9X,'INITIAL NUMBER OF MOLECULES - EITHER TYPE =' ,I13,' INM
1 I'/9X,'MAXIMUM NUMBER OF MOLECULES - EITHER TYPE =' ,I13,' MN
2M I'/1X,'MAX NUMBER OF MOLECULES IN ANY CELL - EITHER TYPE =' ,PRA0570
3I13,' MNB I')
27 FORMAT(/22X,'VELOCITY OF FREE STREAM FLOW =' ,E13.4,' U',T76,'I'/1
19X,'SPEED RATIO OF FREE STREAM FLOW =' ,E13.4,' S',T76,'I'/19X,'MAC
AH NUMBER OF FREE STREAM FLOW =' ,E13.4,' M',T76,'I'/19X,'SPECIFIC H
BEAT RATIO (CALCULATED)=' ,E13.4,' GAMMA',T76,'I'/ 35X,'ANG
2LE OF ATTACK =' ,F13.4,' ANGLE I'/16X,'NUMBER DENSITY OF FREE ST

```

3 REAM FLOW =',E13.4,' N',T76,'I'/19X,'TEMPERATURE OF FREE STREAM FL
4 OW =',F13.4,' TF',T76,'I'/16X,'MOLE FRACTIONS OF FREE STREAM FLOW
5 =',3E13.4,' RNU I'/16X,'MOLECULAR WEIGHTS OF SPECIES ABOVE =',3F13
6.4,' RNA I')
28 FORMAT(/10X,'REFERENCE TEMPERATURE FOR MOLECULAR DATA =',F12.4,'
1 TRF',T90,'I'/14X,'CROSS-SECTION',26X,'TEMP EXPONENT',T90,'I'/3(3X,
23E12.4,3X,3F12.6,T90,'I'/)/5X,'CHI/2-1',11X,'ROTATIONAL PARAMETER
3 PHI',T90,'I'/3(F12.4,5X,3F12.6,T90,'I'/))
29 FORMAT(9X,'DATA SAVED ON TAPE 9')
30 FORMAT(/31X,'REF MOLECULAR SPEED =',E13.4,' VELR',T76,'I'/20X,'SP
1 ECIES FREE STREAM MOLECULAR SPEEDS',T76,'I'/14X,3E16.6,T76,'I'/26X
2,'REFERENCE MEAN FREE PATH =',E13.4,' XREF',T76,'I'/26X,'SPECIES M
3 EAN FREE PATHS',T76,'I'/14X,3E16.6,T76,'I'/11X,'LONGITUDINAL KNUDS
4 EN NUMBER (CALCULATED) =',E13.4,' AKN',T76,'I'/13X,'TRANSVERSE KNUD
5 SEN NUMBER (CALCULATED) =',E13.4,' AKT',T76,'I')

-----
IARRAY=708+LIMIT(3)*(32+56*LIMIT(1))+20*LIMIT(2)+LIMIT(4)*(120+4*LIMIT(6))+56*LIMIT(5)+LIMIT(8)*(68+96*LIMIT(9))+20*LIMIT(7)+224*LIMIT(10)
1 LIMIT(6))+56*LIMIT(5)+LIMIT(8)*(68+96*LIMIT(9))+20*LIMIT(7)+224*LIMIT(10)
2 IT(1)
WRITE(6,1) MSP
IF(NEW) WRITE(6,2) NOT
WRITE(6,3) NAME,TITLE,NREG
WRITE(6,4) XSTART,RMB
DO 100 I=1,NREG
100 WRITE(6,6) XLIM(I+2),(COEFF(J,I),J=1,4)
WRITE(6,10)
WRITE(6,12)
DO 110 I=1,ND
110 WRITE(6,14) XCB(I),TB(I),(ALPHA(J,I),J=1,3),(SIGMA(J,I),J=1,3),YCB
1(I)
WRITE(6,10)
WRITE(6,16) (LD(N),LF(N),N=1,5)
WRITE(6,17) IARRAY,(LIMIT(I),I=1,10)
WRITE(6,18) BW,BH,NW,NH,NL
IF(NL.GT.1) WRITE(6,20) NFA,NCA,NHA,MW,MH
IF(NL.GT.2) WRITE(6,22) NFB,NCB,NHB,LW,LH
IETAZ=THETAZ
JETAZ=180-IETAZ
WRITE(6,23) IETAZ,NWEDGE(1),JETAZ,NWEDGE(2)
DTS=DTM*ITS
DTP=DTM*ITP
AST=DTM*TST
ALIM=DTM*TLIM
CHT=0.0
DO 120 J=1,MSP
120 CHT=CHT+CHI(J)*RNU(J)
GAMMA=(7.+2.*CHT)/(5.+2.*CHT)
AM=S*SQRT(2./GAMMA)
WRITE(6,24) DTM,DTS,DTP,AST,ALIM
WRITE(6,26) JNM,MNM,MNB
WRITE(6,27) U,S,AM,GAMMA,ANGLE,DENF,TF,(RNU(I),I=1,3),(RNA(I),I=1,3)
1)
WRITE(6,28) TRF,((DIR(I,K),K=1,3),(ETA(I,K),K=1,3),I=1,3),(CHI(I),

```

PRA0670

PRA0680

PRA0690

PRA0700

PRA0710

PRA0720

PRA0730

PRA0750

PRA0760

PRA0770

PRA0780

PRA0790

PRA0800

PRA0810

PRA0820

PRA0850

PRA0860

PRA0880

PRA0890

PRA0900

PRA0910

PRA0920

PRA0930

PRA0940

PRA0950

PRA0960

PRA0970

PRA0980

PRA0990

PRA0970

PRA0980

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PRA0980

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1 (PHI (I,K),K=1,3),I=1,3)
DO 210 I=1,3
VELS(I)=0.0
210 XSP(I)=0.0
DO 220 J=1,MSP
VELS(J)=VELR/SQRT(WTM(J))
XT=0.0
DO 215 M=1,MSP
215 XT=XT+RNU(M)*DAM(J,M)*SQRT(1.+WTM(J)/WTM(M))
220 XSP(J)=1.414214*XREF/XT
AKT=1./RMB
WRITE(6,30) VELR,(VELS(I),I=1,3),XREF,(XSP(I),I=1,3),AKN,AKT
IF(SAVE) WRITE(6,29)
RETURN
END

```

PRA1040
PRA1050
PRA1060

```

SUBROUTINE PRINTB(PNA,MSP,PNB,LEV,LWP,NM,RLD,XLIM,XC,YC,ZC,NB,
1 NUMCEL,LKW,N)
INTEGER*2 LKW,NB,NUMCEL
DIMENSION PNB(1),LEV(1),LWP(1),NM(1),RLD(1),XLIM(1),XC(1)
DIMENSION YC(1),ZC(1),NB(N,1),NUMCEL(1),LKW(1)
COMMON /FIRST/NL
COMMON /THIRD/PI,NREG,S,SINANG,COSANG,AKN
COMMON /FOUR/NBX
1 FORMAT(1H1)
2 FORMAT(2X,'-----CELL GEOMETRY-----')
1-----'/2X,'BOX LEVEL POSITION OF CENTER VOLUME WPRB0100
2 EIGHTING POPULATION'/2X,'NUM.',12X,'X',7X,'Y THETA',12X,'FACTOPRB0110
3R ',15X,' CELL#')
3 FORMAT(1X,I4,I5,3X,2P8.3,P7.1,E12.3,2X,I2,4X,3I5,3X,I4)
4 FORMAT(2X,'-----TOTALS-----',E12.4,8X,3I5)
WRITE(6,1)
WRITE(6,2)
DO 200 I=1,NBX
IF(NUMCEL(I).EQ.0) GO TO 200
X=(XC(I)-XLIM(2))*AKN
Y=YC(I)*AKN
LEVEL=1
IF(NL.LT.2) GO TO 120
IF(I.LT.LEV(1)) GO TO 120
LEVEL=2
IF(NL.LT.3) GO TO 120
IF(I.LT.LEV(2)) GO TO 120
LEVEL=3
120 CONTINUE
J=NUMCEL(I)
M1=NB(1,J)
M2=0
IF(MSP.GE.2) M2=NB(2,J)
M3=0
IF(MSP.GE.3) M3=NB(3,J)
140 WRITE(6,3) I,LEVEL,X,Y,ZC(I),PNB(I),LKW(J),M1,M2,M3,J
200 CONTINUE
NM2=0
IF(MSP.GE.2) NM2=NM(2)

```

PRB0050
PRB0060
PRB0070
PRB0080
PRB0090
PRB0100
PRB0110
PRB0150
PRB0170
PRB0180
PRB0190
PRB0200
PRB0210
PRB0220
PRB0230
PRB0240
PRB0250
PRB0260
PRB0320
PRB0370

```

NM3=0
IF (MSP.GE.3) NM3=NM(3)
WRITE(6,4) PNA,NM(1),NM2,NM3
RETURN
END

```

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PRB0390
PRB0400

```

SUBROUTINE SIMPSN(A,B,L,INTGRL,PERCNT,COEFF,PIROOT,SUM1,FUN)
REAL INTGRL
DIMENSION COEFF(4,1)

```

SIMP010
SIMP020
SIMP030
SIMP040

THE PURPOSE OF THIS SUBROUTINE IS TO PERFORM A SIMPSON'S RULE
INTEGRATION.

SIMP050
SIMP060
SIMP070

FORMAT

SIMP080
SIMP090
SIMP100

4 FORMAT(/30H TOO MANY ITERATIONS. TEST IS E15.7,14H, INTEGRAL IS E15.7,14H IN THE RANGE E15.7,4H TO E15.7)

SIMP110
SIMP120
SIMP130

```

PREV=0.0
N=4
N1=3
G=.5*(B-A)+A
SUME=FUN(G,PIROOT,L,COEFF)
SUMO=0.0
K=0

```

SIMP140
SIMP150
SIMP160

```

235 SUME=SUME+SUMO
N1=N1+K
K=N1+1

```

SIMP170
SIMP180
SIMP190

```

IF (K.LT.5000) GO TO 237
WRITE(6,4) TEST,PREV,A,B
STOP

```

SIMP200
SIMP210
SIMP220

```

237 SUMO=0.0
DO 240 I=1,N1,2
G=(I*(B-A))/K+A

```

SIMP230
SIMP240
SIMP250

```

240 SUMO=SUMO+FUN(G,PIROOT,L,COEFF)
INTGRL=(SUM1+4.*SUMO+2.*SUME)/(3.*K)
TEST=ABS(2.-4.*PREV/(INTGRL+PREV))
PREV=INTGRL

```

SIMP260
SIMP270
SIMP280

```

IF (TEST.GT.PERCNT) GO TO 235
RETURN
END

```

SIMP290
SIMP300
SIMP310

```

FUNCTION FNCTN(X,PIROOT,L,COEFF)
DIMENSION COEFF(4,1)

```

SIMP320
SIMP330
SIMP340

SIMP350
SIMP360
SIMP370

FNCN010
FNCN020
FNCN030

THE PURPOSE OF THIS FUNCTION IS TO EVALUATE THE INTEGRAND USED IN
THE SIMPSON-S RULE INTEGRATION ROUTINE.

FNCN040
FNCN050
FNCN060

```

A=COEFF(1,L)
B=COEFF(2,L)
C=COEFF(3,L)
D=COEFF(4,L)
AA=4.*(A-B)*(A*X+C)*X+C*C-4.*B*D

```

FNCN070
FNCN080
FNCN090
FNCN100
FNCN110

```

IF(AA.LT.0.) AA=0.
FNCTN=SQRT(AA)
RETURN
END

```

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FNCH120
FNCH130
FNCH140
FNCH150

```

SUBROUTINE HEIGHT(X,R,L,COEFF,I)
DIMENSION COEFF(4,1),DBG2(2,3)
DATA DBG2/'SBTR','CT ','GAS ',' ','RUN',' '/

```

HGHT010
HGHT020
HGHT030

THE PURPOSE OF THIS SUBROUTINE IS TO COMPUTE THE R-COORDINATE OF
THE BODY CURVE AT ANY GIVEN X-COORDINATE, IN THE X-R PLANE.

HGHT040
HGHT050
HGHT060
HGHT070

```

2 FORMAT(//' MESSAGE ',I2,E17.8,3X,2A4)
ARG=-((COEFF(1,L)*X+COEFF(3,L))*X+COEFF(4,L))/COEFF(2,L)
IF(ARG.GE.0.) GO TO 100
WRITE(6,2) L,X,(DBG2(M,I),M=1,2)
ARG=0.
100 R=SQRT(ARG)
RETURN
END

```

HGHT080
HGHT090
HGHT100
HGHT110
HGHT120
HGHT130
HGHT140
HGHT150

```

FUNCTION FNCTM(ARG,PIROOT,L,COEFF)
DIMENSION COEFF(4,1)
D=0.
IF(ABS(ARG).LT.10.) D=EXP(-ARG*ARG)/PIROOT
E=0.
IF(ARG.GT.-10.) E=ARG*ERRF(ARG)
TEM=.5*(D+E)
FNCTM=TEM
RETURN
END

```

FNCH010
FNCH020
FNCH030
FNCH040
FNCH050
FNCH060
FNCH070
FNCH080
FNCH090
FNCH100

```

SUBROUTINE CELL(TH,LEV,A,B,K,KH,XO,I,J,DELANG,NWEDGE,XC,YC,ZC,FNB)
DIMENSION DELANG(1),NWEDGE(1),XC(1),YC(1),ZC(1),FNB(1)
COMMON /THIRD/PI

```

CELL010

THE PURPOSE OF THIS SUBROUTINE IS TO

1. COMPUTE THE VOLUME OF EACH CELL (ALL 3 POSSIBLE LEVELS)
AND STORE THE RESULT IN THE ARRAY CALLED 'FNB'.
CELL070
CELL080
2. COMPUTE THE X, R, AND THETA COORDINATES OF THE CENTER OF
EACH CELL (ALL 3 POSSIBLE LEVELS) AND STORE THE RESULTS IN
ARRAYS CALLED 'XC', 'YC', AND 'ZC'.
CELL090
CELL100
CELL110

CELL120
CELL130

```

INDEX=I+J
ZO=0.
DO 120 MT=1,LEV
ICNT=NWEDGE(MT)
ANGLE=DELANG(MT)
FACTOR=ANGLE/180.*PI*B*B*A
Z=ZO-.5*ANGLE
DO 110 L=1,ICNT
Z=Z+ANGLE
Y=-.5*B

```

CELL140
CELL150
CELL160
CELL170
CELL180
CELL190
CELL200
CELL210
CELL220
CELL230

DO 110 N=1,KH	CELL240
X=XO-.5*A	CELL250
Y=Y+B	CELL260
DO 110 M=1,K	CELL270
X=X+A	CELL280
INDEX=INDEX+1	CELL290
XC(INDEX)=X	CELL300
YC(INDEX)=Y	CELL310
ZC(INDEX)=Z	CELL320
110 FNB(INDEX)=FACTOR*(2*N-1)	CELL330
ZO=TH	CELL340
120 CONTINUE	CELL350
RETURN	CELL360
END	CELL370
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SUBROUTINE ZERO(NWIDE,NHI,NBEG,NLONG,MUP,NAREA,ICNT,FNB)	ZERO010
DIMENSION FNB(1)	ZERO020

THIS SUBROUTINE SETS THE SIZES TO ZERO OF THOSE CELLS WHICH ARE TO BE SUBDIVIDED INTO SMALLER CELLS.	ZERO030
	ZERO040
	ZERO050
	ZERO060
	ZERO070

NGO=NBEG+1	ZERO080
NSTOP=NBEG+NLONG	ZERO090
DO 110 N=NGO,NSTOP	ZERO100
DO 110 M=1,MUP	ZERO110
DO 110 L=1,ICNT	ZERO120
INDEX=NWIDE*(NHI*(L-1)+M-1)+N+NAREA	ZERO130
110 FNB(INDEX)=0.0	ZERO140
RETURN	ZERO150
END	ZERO160
	ZERO170
SUBROUTINE SBTRCT(NGO,NTMP,NSTOP,BWIDTH,BHITE,DELANG,XC,YC,FNB,XLSBCT010 1IM,COEFF)	SBCT020
DIMENSION DELANG(1),XC(1),YC(1),FNB(1),XLIM(1),COEFF(4,1)	SBCT030
COMMON /THIRD/PI,NREG	SBCT040

THIS SUBROUTINE SUBTRACTS FROM EACH CELL SIZE ('FNB' ARRAY) THAT PORTION OCCUPIED BY THE BODY.	SBCT050
	SBCT060
	SBCT070
	SBCT080
	SBCT090

FACTOR=PI*DELANG(1)/180.	SBCT100
DO 150 N=NGO,NSTOP	SBCT110
IF(FNB(N).LE.0.) GO TO 150	SBCT120
DFNB=.005*FNB(N)	SBCT130
SLICE=.01*BWIDTH	SBCT140
IF(N.GT.NTMP) FACTOR=PI*DELANG(2)/180.	SBCT150
X=XC(N)-.5*(BWIDTH+SLICE)	SBCT160
YBOT=YC(N)-.5*BHITE	SBCT170
YTOP=YBOT+BHITE	SBCT180
DO 130 M=1,100	SBCT190
X=X+SLICE	SBCT200
IF(X.LE.XLIM(2)) GO TO 130	SBCT210
	SBCT220

DO 120 L=1,NREG	SBCT230
IF(X.LT.XLIM(L+2)) GO TO 125	SBCT240
120 CONTINUE	SBCT250
GO TO 130	SBCT260
125 CALL HEIGHT(X,YBODY,L,COEFF,1)	SBCT270
IF(YBODY.LE.YBOT) GO TO 130	SBCT280
YTEMP=YTOP	SBCT290
IF(YBODY.LT.YTOP) YTEMP=YBODY	SBCT300
FNB(N)=FNB(N)-SLICE*(YTEMP*YTEMP-YBOT*YBOT)*FACTOR	SBCT310
130 CONTINUE	SBCT320
IF(FNB(N).LT.DFNB) FNB(N)=0.	SBCT330
150 CONTINUE	SBCT340
RETURN	SBCT350
END	SBCT360

```

SUBROUTINE IMPACT(RM,G1,G2,G3,ET,EI,PHI,CHI,ETA,XM,CIM)
COMMON/THIRD/PI
IF(PHI.EQ.0.) GO TO 20
IF(CHI.EQ.0.) GO TO 20
DF=PHI*CHI-1.
DS=PHI*(2.-.5*ETA)-1.
E=ET+EI
10 X=Rand(0)
IF(X.EQ.0.0) GO TO 10
XT=X**DF*(1.-X)**DS
IF(XT.GT.XM) GO TO 15
CIM=CIM+XT
IF(CIM.LT.XM) GO TO 10
CIM=CIM-XM
15 ET=(1.-PHI)*ET+(1.-X)*PHI*E
EI=(1.-PHI)*EI+X*PHI*E
20 GP=SQRT(ET/RM)
EP=2.*PI*Rand(0)
CSX=2.*Rand(0)-1.
SSX=SQRT(1.-CSX**2)
G1=GP*CSX
G2=GP*SSX*COS(EP)
G3=GP*SSX*SIN(EP)
RETURN
END

```

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SUBROUTINE GAS(NWEDG,THETAZ,DELANG,NWEDGE,BTA,C1,DPA,NM,RLD,LWF,FNGAS0010
1B,DB,NB,NBF,LPF,PAU,PAV,PAW,PAX,PAY,PAZ,XLIM,COEFF,LM,I2,I3,LARGE,GAS0020
2MMH,MNB,DEBUG1,LCOL,NUMCEL,IP,ER,CHI,CNG,CMG,I,LB,NBM,NBN)
INTEGER*2 LM(I,1),LPF(I,1),LCOL(I,1),LB(1),NBM(I,1),NBN(1)
INTEGER*2 NB,NBF,NUMCEL
INTEGER Q
LOGICAL DUMP,DEBUG1
DIMENSION DELANG(2),NWEDGE(2),NUMCEL(1)
DIMENSION BTA(1),C1(1),DPA(1),NM(1),RLD(1),LWF(1),FNB(1),CHI(1)
DIMENSION DB(I,1),NB(I,1),NBF(I,1),PAU(I,1),PAV(I,1),PAW(I,1)
DIMENSION PAX(I,1),PAY(I,1),PAZ(I,1),ER(I,1),COEFF(4,1),XLIM(1)
DIMENSION CNG(1),CMG(1)
COMMON /FIRST/NL,NW,NH,MW,MH,LW,LH,NXA,NXB
COMMON /SECND/BW,BH,BWB,BHB,BWC,BHC,XLB,XLC

```

GAS0050
GAS0060
GAS0070

GAS0110
GAS0120

CM0000 /T0100/11,0000,3,01000,00000
CM0000 /P0000/11,00,10,0100,00,0000,LL

GAS0130

GAS0140

GAS0150

GAS0160

GAS0170

GAS0180

GAS0190

GAS0200

GAS0210

GAS0220

GAS0230

GAS0240

GAS0250

GAS0260

GAS0270

GAS0280

GAS0290

GAS0300

GAS0310

GAS0320

GAS0330

GAS0340

GAS0350

GAS0360

GAS0370

GAS0380

GAS0390

GAS0400

GAS0410

GAS0420

GAS0430

GAS0440

GAS0450

GAS0460

GAS0470

GAS0480

GAS0490

GAS0500

GAS0510

GAS0520

GAS0530

GAS0540

GAS0550

GAS0560

GAS0570

THE PURPOSE OF THIS SUBROUTINE IS TO

1. COMPUTE THE INITIAL VELOCITY OF EACH MOLECULE.

THE INITIAL VELOCITIES ARE 'PAX', 'PAY', AND 'PAZ'.

2. COMPUTE THE INITIAL POSITIONS OF EACH MOLECULE.

THE INITIAL POSITIONS ARE 'PAX', 'PAY', AND 'PAZ'.

3. COMPUTE AN INITIAL SET OF VALUES FOR THE CELL POPULATIONS-

'PAX', 'PAY', AND 'PAZ'.

4. COMPUTE AN INITIAL SET OF VALUES FOR THE CELL POPULATIONS-

'PAX', 'PAY', AND 'PAZ'.

5. COMPUTE AN INITIAL SET OF VALUES FOR THE CELL POPULATIONS-

'PAX', 'PAY', AND 'PAZ'.

6. COMPUTE AN INITIAL SET OF VALUES FOR THE CELL POPULATIONS-

'PAX'.

7. COMPUTE AN INITIAL SET OF VALUES FOR THE CELL POPULATIONS-

'PAX', 'PAY', AND 'PAZ'.

8. COMPUTE AN INITIAL SET OF VALUES FOR THE CELL POPULATIONS-

'PAX', 'PAY', AND 'PAZ'.

9. COMPUTE AN INITIAL SET OF VALUES FOR THE CELL POPULATIONS-

'PAX', 'PAY', AND 'PAZ'.

10. COMPUTE AN INITIAL SET OF VALUES FOR THE CELL POPULATIONS-

'PAX', 'PAY', AND 'PAZ'.

11. COMPUTE AN INITIAL SET OF VALUES FOR THE CELL POPULATIONS-

'PAX', 'PAY', AND 'PAZ'.

12. COMPUTE AN INITIAL SET OF VALUES FOR THE CELL POPULATIONS-

'PAX', 'PAY', AND 'PAZ'.

13. COMPUTE AN INITIAL SET OF VALUES FOR THE CELL POPULATIONS-

'PAX', 'PAY', AND 'PAZ'.

14. COMPUTE AN INITIAL SET OF VALUES FOR THE CELL POPULATIONS-

'PAX', 'PAY', AND 'PAZ'.

15. COMPUTE AN INITIAL SET OF VALUES FOR THE CELL POPULATIONS-

'PAX', 'PAY', AND 'PAZ'.

16. COMPUTE AN INITIAL SET OF VALUES FOR THE CELL POPULATIONS-

'PAX', 'PAY', AND 'PAZ'.

17. COMPUTE AN INITIAL SET OF VALUES FOR THE CELL POPULATIONS-

'PAX', 'PAY', AND 'PAZ'.

18. COMPUTE AN INITIAL SET OF VALUES FOR THE CELL POPULATIONS-

'PAX', 'PAY', AND 'PAZ'.

19. COMPUTE AN INITIAL SET OF VALUES FOR THE CELL POPULATIONS-

'PAX', 'PAY', AND 'PAZ'.

20. COMPUTE AN INITIAL SET OF VALUES FOR THE CELL POPULATIONS-

'PAX', 'PAY', AND 'PAZ'.

21. COMPUTE AN INITIAL SET OF VALUES FOR THE CELL POPULATIONS-

'PAX', 'PAY', AND 'PAZ'.

22. COMPUTE AN INITIAL SET OF VALUES FOR THE CELL POPULATIONS-

'PAX', 'PAY', AND 'PAZ'.

23. COMPUTE AN INITIAL SET OF VALUES FOR THE CELL POPULATIONS-

'PAX', 'PAY', AND 'PAZ'.

24. COMPUTE AN INITIAL SET OF VALUES FOR THE CELL POPULATIONS-

'PAX', 'PAY', AND 'PAZ'.

25. COMPUTE AN INITIAL SET OF VALUES FOR THE CELL POPULATIONS-

'PAX', 'PAY', AND 'PAZ'.

ORIGINAL PAGE IN
OF POOR QUALITY

125 IF (CHI(NT).LE.-1.) GO TO 130
X=9.*RAND(0)
IF (X.EQ.0.0) GO TO 125
IT=X*CHI(NT)*EXP(-X)
IF (IT.GE.CNG(NT)) GO TO 126
CNG(NT)=CNG(NT)+IT
IF (CNG(NT).LT.CNG(NT)) GO TO 125
CNG(NT)=CNG(NT)-CNG(NT)
126 ER(NT,N)=X
130 X=XR*RAND(0)
140 R=RM*RAND(0)
LA=0
150 LA=LA+1
IF (R.GT.RLD(LA)) GO TO 150
A=LWF(LA)

C9=C9+R/A
 IF (C9.LT.RWPM) GO TO 140
 C9=C9-RWPM
 IF (X.LE.XLIN(2)) GO TO 159
 DO 152 L=1,NREG
 IF (X.LT.XLIN(L+2)) GO TO 154
 152 CONTINUE
 GO TO 159
 154 CALL HEIGHT(X,YBODY,L,COEFF,2)
 IF (R.LT.YBODY) GO TO 130
 159 PAX(MT,N)=X
 D=PI*RAND(0)
 PAY(MT,N)=R*COS(D)
 PAZ(MT,N)=R*SIN(D)
 TANG=180.*(1.-D/PI)
 IWDGE=TANG/DELANG(1)
 IF ((IWDGE.GE.NWEDGE(1)).AND.(DELANG(2).NE.0.)) IWDGE=(TANG-THETAZ)
 1/DELANG(2)+NWEDGE(1)
 IF (IWDGE.GE.NWEDGE) IWDGE=NWEDGE-1
 L=X/BW+1.
 IF (L.GT.NW) L=NW
 M=R/BH
 IF (M.GE.NH) M=NH-1
 K=(IWDGE*NH+M)*NW+L
 IF (K.LE.NXA) GO TO 160
 WRITE(6,2) L,M,K,MT,N,DELANG,NWEDGE,NWEDGE,D,TANG,IWDGE,NH,NW
 IF (DUMP) CALL ABEND(11)
 STOP
 160 IF (NL.EQ.1) GO TO 162
 IF (IWDGE.GE.NWEDGE(1)) GO TO 162
 IF (PNB(K).GT.0.) GO TO 162
 L=(X-XLB)/BWB+1.
 IF (L.GT.MW) L=MW
 M=R/BHB
 IF (M.GE.MH) M=MH-1
 K=(IWDGE*MH+M)*MW+L+NXA
 IF (K.LE.NXA+NXB) GO TO 161
 WRITE(6,2) L,M,K,MT,N,DELANG,NWEDGE,NWEDGE,D,TANG,IWDGE,MH,MW,NXA
 IF (DUMP) CALL ABEND(12)
 STOP
 161 IF (NL.EQ.2) GO TO 162
 L=(X-XLC)/BWC+1.
 IF (PNB(K).GT.0.) GO TO 162
 IF (L.GT.LW) L=LW
 M=R/BHC
 IF (M.GE.LH) M=LH-1
 K=(IWDGE*LH+M)*LW+L+NXA+NXB
 IF (K.LE.NBX) GO TO 164
 WRITE(6,2) L,M,K,MT,N,DELANG,NWEDGE,NWEDGE,D,TANG,IWDGE,LH,LW,NXA,NXB
 1B
 IF (DUMP) CALL ABEND(13)
 STOP
 162 IF (PNB(K).GT.0.) GO TO 164
 WRITE(6,3) L,M,K,MT,N,PNB(K)
 IF (DUMP) CALL ABEND(14)

ORIGINAL PAGE 15
OF POOR QUALITY

GAS0580
 GAS0590
 GAS0600
 GAS0610
 GAS0620
 GAS0630
 GAS0640
 GAS0650
 GAS0660
 GAS0670
 GAS0680
 GAS0690
 GAS0700
 GAS0710
 GAS0720
 GAS0730
 GAS0750
 GAS0760
 GAS0770
 GAS0780
 GAS0790
 GAS0800
 GAS0810
 GAS0820
 GAS0830
 GAS0840
 GAS0850
 GAS0870
 GAS0880
 GAS0890
 GAS0900
 GAS0910
 GAS0920
 GAS0930
 GAS0940
 GAS0950
 GAS0960
 GAS0970
 GAS0990
 GAS0980
 GAS1000
 GAS1010
 GAS1020
 GAS1030
 GAS1040
 GAS1050
 GAS1060
 GAS1070
 GAS1080
 GAS1090
 GAS1100
 GAS1110

ORIGINAL PAGE IS
OF POOR QUALITY

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STOP
164 Q=NUMCEL (K)
IF (Q.GT.0) GO TO 165
WRITE (6,2) Q,L,M,K,N,DELANG,NWEDGE,NWEDG,D,TANG,IWDGE,LH,LW,NXA,NXB
1,X,R
IF (DUMP) CALL ABEND(15)
STOP
165 J=NB (MT,Q) +1
KW=LWP (LA)
LPP (MT,N)=KW
LCOL (MT,N)=0
LLC=LLC+KW
IF (J.LE.MNB) GO TO 166
IF (DEBUG1) WRITE (6,4) MT,Q,MNB
GO TO 167
166 NB (MT,Q) =J
LB (N)=Q
NBF (MT,Q)=NBF (MT,Q)+KW
167 IF (LLC.LT.LL) GO TO 110
NM (MT)=N
NBM (MT,1)=0
DO 170 Q=1,NBX
N=NUMCEL (Q)
IF (N.EQ.0) GO TO 170
A=NBF (MT,N)
DB (MT,N)=A*DPA (MT)/FNB (Q)
NBM (MT,N+1)=NBM (MT,N)+NB (MT,N)
NBN (N)=NBM (MT,N)
170 CONTINUE
NG=NM (MT)
DO 175 N=1,NG
Q=LB (N)
NBN (Q)=NBN (Q)+1
NA=NBN (Q)
175 LM (MT,NA)=N
180 CONTINUE
RETURN
190 LARGE=1
RETURN
END

```

GAS1120
GAS1130
GAS1140
GAS1150
GAS1160
GAS1170
GAS1180
GAS1190
GAS1200
GAS1210
GAS1220
GAS1230
GAS1240
GAS1250
GAS1260
GAS1270

GAS1290
GAS1300
GAS1310

GAS1330

GAS1360
GAS1370

GAS1380

GAS1390
GAS1400
GAS1410
GAS1420
GAS1430

SUBROUTINE FLOW (NWEDG,MNM,LARGE,BTA,C1,C7,C8,D1,D2,D3,D4,DTH,NM,SNFLO0010
1,ST,THETA,LWP,RLD,PTH,ENTS,REMS,SSA,SSB,PAU,PAV,PAW,PAX,PAY,PAZ,LPPLO0020
2P,ENT,REM,LCOL,IP,ER,CHI,CNG,CMG,I)

INTEGER*2 LPP,LCOL FLO0040
DIMENSION BTA (1),D1 (1),D2 (1),D3 (1),D4 (1),DTH (1),NM (1),SN (1),ST (1) FLO0050
DIMENSION C1 (1),C7 (1),C8 (1),LWP (1),RLD (1),THETA (1),SSA (2,1) FLO0060
DIMENSION SSB (2,1),PAU (1,1),PAV (1,1),PAW (1,1),PAX (1,1),PAY (1,1)
DIMENSION PAZ (1,1),LPP (1,1),ENT (2,3,6,1),REM (2,3,6,1),LCOL (1,1)
DIMENSION ENTS (3,1),REMS (3,1),PTH (3,1),ER (1,1),CHI (1),CNG (1),
1CMG (1)

COMMON /THIRD/PI FLO0100
COMMON /FORTH/NBX,RM,XR FLO0110

FLO0120

THE PURPOSE OF THIS SUBROUTINE IS TO ADD A NEW BATCH OF MOLECULES FLO0130

TO THE SAMPLE THROUGH THE UPSTREAM BOUNDARY.

		FLO0140
-----		FLO0150
DO 370 MT=1,IP		
ARG=SN(MT)		FLO0170
XGO=0.		FLO0180
E=1.		FLO0190
DO 180 NT=1,2		FLO0200
SM=AMAX1(0.,ARG-4.)		FLO0210
SSM=AMAX1(0.,ARG)		FLO0220
TEMPC=0.		FLO0230
DO 170 LA=1,6		FLO0240
TEMPB=RLD(LA)*RLD(LA)		FLO0250
AY = TEMPB-TEMPC		FLO0260
C = TEMPC		FLO0270
TEMPC=TEMPB		FLO0280
DO 170 K=1,NWEDG		FLO0290
AM=ENT(NT,MT,LA,K)+REM(NT,MT,LA,K)		FLO0300
M=AM		FLO0310
AMM=M		FLO0320
REM(NT,MT,LA,K)=AM-AMM		FLO0330
IF(M.EQ.0) GO TO 170		FLO0340
DY =AY/AMM		FLO0350
DO 160 N=1,M		FLO0360
IF(NM(MT).GE.MNM) GO TO 380		
NM(MT)=NM(MT)+1		FLO0370
NMX=NM(MT)		FLO0390
R = SQRT(C + DY*(N+RAND(0) -1.))		FLO0400
D=(THETA(K)+RAND(0)*DTH(K))*PI/180.		FLO0410
PAY(MT,NMX)=R*COS(D)		FLO0420
PAZ(MT,NMX)=R*SIN(D)		FLO0430
LPP(MT,NMX)=LWP(LA)		FLO0440
LCOL(MT,NMX)=0		FLO0450
130 V=SM+RAND(0)*(SSM+4.-SM)		FLO0460
C1(MT)=C1(MT)+2.*V*EXP(SSB(NT,MT)+2.*ARG*V-V*V)/SSA(NT,MT)		FLO0470
IF(C1(MT).LT.1.) GO TO 130		FLO0480
C1(MT)=C1(MT)-1.		FLO0490
PAU(MT,NMX)=E*V/BTA(MT)		FLO0500
140 V=8.*RAND(0)-4.		FLO0510
C7(MT)=C7(MT)+EXP(-V*V)		FLO0520
IF(C7(MT).LT.1.) GO TO 140		FLO0530
C7(MT)=C7(MT)-1.		FLO0540
PAV(MT,NMX)=(V+ST(MT))/BTA(MT)		FLO0550
150 V=8.*RAND(0)-4.		FLO0560
C8(MT)=C8(MT)+EXP(-V*V)		FLO0570
IF(C8(MT).LT.1.) GO TO 150		FLO0580
C8(MT)=C8(MT)-1.		FLO0590
PAW(MT,NMX)=V/BTA(MT)		FLO0600
ER(MT,NMX)=0.0		
IF(CHI(MT).LE.-1.) GO TO 160		
125 X=9.*RAND(0)		
IF(X.LE.0.0) GO TO 125		
XT=X**CHI(MT)*EXP(-X)		
IF(XT.GE.CMG(MT)) GO TO 126		
CNG(MT)=CNG(MT)+XT		
IF(CNG(MT).LT.CMG(MT)) GO TO 125		

ORIGINAL PAGE 13
OF POOR QUALITY

CNG(MT)=CNG(MT)-CMG(MT)	
126 ER(MT,NMX)=X	
160 PAX(MT,NMX)=XGO	
170 CONTINUE	FLO0610
ARG=-ARG	FLO0620
XGO=XR	FLO0630
E=-1.	FLO0640
180 CONTINUE	FLO0650
DO 370 K=1,NWEDG	FLO0660
AM=ENTS(MT,K)+REMS(MT,K)	FLO0670
M=AM	FLO0680
AMM=M	FLO0690
REMS(MT,K)=AM-AMM	FLO0700
IF(M.EQ.0) GO TO 370	FLO0710
DX=XR/AMM	FLO0720
DO 365 N=1,M	FLO0730
IF(NM(MT).GE.MNM) GO TO 380	FLO0740
NM(MT)=NM(MT)+1	
NMX=NMX(MT)	FLO0750
PAX(MT,NMX)=(N-1.+RAND(0))*DX	FLO0770
330 TH=(THETA(K)+RAND(0)*DTH(K))*PI/180.	FLO0780
A=COS(TH)	FLO0790
B=SIN(TH)	FLO0800
SM=ST(MT)*A	FLO0810
C=0.	FLO0820
IF(ABS(SM).LT.10.) C=EXP(-SM*SM)	FLO0830
D=0.	FLO0840
IF(SM.GT.-10.) D=SQRT(PI)*SM*ERRF(SM)	FLO0850
D1(MT)=D1(MT)+(C+D)/PTH(MT,K)	FLO0860
IF(D1(MT).LT.1.) GO TO 330	FLO0870
D1(MT)=D1(MT)-1.	FLO0880
PAY(MT,NMX)=-RM*A	FLO0890
PAZ(MT,NMX)=RM*B	FLO0900
VNM=.5*SM+SQRT(.25*SM*SM+.5)	FLO0910
VM=AMAX1(0.,SM-4.)	FLO0920
340 V=VM+RAND(0)*(SM+4.-VM)	FLO0930
D2(MT)=D2(MT)+V*EXP(VNM*(VNM-2.*SM)-V*(V-2.*SM))/VNM	FLO0940
IF(D2(MT).LT.1.) GO TO 340	FLO0950
D2(MT)=D2(MT)-1.	FLO0960
VN=V	FLO0970
350 V=8.*RAND(0)-4.	FLO0980
D3(MT)=D3(MT)+EXP(-V*V)	FLO0990
IF(D3(MT).LT.1.) GO TO 350	FLO1000
D3(MT)=D3(MT)-1.	FLO1010
VT1=SN(MT)+V	FLO1020
360 V=8.*RAND(0)-4.	FLO1030
D4(MT)=D4(MT)+EXP(-V*V)	FLO1040
IF(D4(MT).LT.1.) GO TO 360	FLO1050
D4(MT)=D4(MT)-1.	FLO1060
VT2=ST(MT)*B+V	FLO1070
PAU(MT,NMX)=VT1/BTA(MT)	FLO1080
PAV(MT,NMX)=(VN*A+VT2*B)/BTA(MT)	FLO1090
PAW(MT,NMX)=(-VN*B+VT2*A)/BTA(MT)	FLO1100
LCOL(MT,NMX)=0	FLO1110
ER(MT,NMX)=0.0	FLO1120

ORIGINAL PAGE IS
OF POOR QUALITY

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      IF (CHI(MT).LE.-1.) GO TO 365
225  X=9.*RAND(0)
      IF (X.EQ.0.0) GO TO 225
      XT=X**CHI(MT)*EXP(-X)
      IF (XT.GE.CMG(MT)) GO TO 226
      CNG(MT)=CNG(MT)+XT
      IF (CNG(MT).LT.CMG(MT)) GO TO 225
      CNG(MT)=CNG(MT)-CMG(MT)
226  PR(MT,NMX)=X
365  LPP(MT,NMX)=LWF(6)
370  CONTINUE
      RETURN
380  LARGE=2
      RETURN
      END

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ORIGINAL PAGE 1
OF POOR QUALITY

FLO1130
FLO1140
FLO1150
FLO1160
FLO1170
FLO1180

```

SUBROUTINE COLIDE(CN,CM,WTM,DB,DBA,NB,NCOL,LCOL,PAU,PAV,PAW,ER,T,
1LM,MT,I2,I3,NUMCEL,ETA,PHI,CHI,CN8,NP,LPP,LKW,NBP,NBM)

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COL0030

```

      INTEGER TIME
      INTEGER*2 LM(NP,1), LCOL(NP,1),LPP(NP,1),LKW(1)
      INTEGER*2 NBM,NB,NBP,NUMCEL
      DIMENSION CN(3,3,1),CM(3,3,1),WTM(1),DB(NP,1),DBA(NP,1),NB(NP,1)
      DIMENSION NCOL(3,1),T(NP,NP,1),NUMCEL(1),ETA(3,1),PHI(3,1),CHI(1)
      DIMENSION PAU(NP,1),PAV(NP,1),PAW(NP,1),ER(NP,1),CN8(3,1),WA(2)
      DIMENSION NBP(NP,1),NBM(NP,1)
      COMMON /FORTH/NBX
      COMMON /FIFTH/ND,TIME,DTM
1  FORMAT(' TIME = ',F9.4,' COLL. TIMES = ',2F9.4,' NUMBERS = ',2I5/)
2  FORMAT(' COLIDE REACHED LINE 160 IN BOX NUMBER = ',I5,' AT CPU TIM
1E = ',F9.4/' VR= ',E12.4,' REL VEL G = ',3E12.4,' EI = ',E12.4)
3  FORMAT(' COLIDE REACHED LINE 165 IN BOX NUMBER = ',I5,' AT CPU TIM
1E = ',F9.4/' VR= ',E12.4,' REL VEL G = ',3E12.4,' EI = ',E12.4)

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COL0080

COL0090

COL0100

THE PURPOSE OF THIS SUBROUTINE IS TO ADVANCE THE ELAPSED TIMES IN
CELLS BY AN AMOUNT APPROXIMATELY EQUAL TO THE PRE-SELECTED COLLIS
TIME. THERE ARE FOUR TIMES FOR EACH CELL, SAVED IN AN ARRAY CALLE
'T', CORRESPONDING TO THE FOUR TYPES OF MOLECULAR COLLISIONS WHIC
CAN OCCUR. TO ADVANCE THE VARIOUS TIMES, AN APPROPRIATE NUMBER OF
THE CORRESPONDING MOLECULAR COLLISIONS IS COMPUTED. THE ACTUAL
MOLECULES TO COLLIDE ARE SELECTED AT RANDOM, AND THEIR VELOCITY V
DIRECTIONS AFTER COLLISION ARE SELECTED AT RANDOM.

COL0180

COL0190

```

      AIME=DTM*TIME
      DO 240 MTA=1,MT
      DO 230 MTB=1,MTA
      D = WTM(MTA) + WTM(MTB)
      WA(MTA)=WTM(MTA)/D
      WA(MTB)=WTM(MTB)/D
      RM=WTM(MTA)*WTM(MTB)/D
      CHT=CHI(MTA)+CHI(MTB)+2.0
      PHT=PHI(MTA,MTB)
      ETT=ETA(MTA,MTB)
      DO 220 M=1,NBX
      N=NUMCEL(M)
      IF (N.LE.0) GO TO 220

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      AKW=LKW(N)
      IF (T(MTA,MTB,N).LT.AIME) GO TO 100
      IF (T(MTB,MTA,N).GE.AIME) GO TO 220
100  NA=NB(MTA,N)*NB(MTB,N)
      IF (MTA.EQ.MTB) NA=(NA-NB(MTA,N))/2.
      IF (NA.LT.1) GO TO 220
      KS=0
120  KC=0
      CPUT=ELTIME(0)
      KS=KS+1
      IF (KS.GT.NA) GO TO 220
130  KC=KC+1
      IF (KC.GT.NA) GO TO 220
135  I=NB(MTA,N)*RAND(0)+1+NBM(MTA,N)
      IF (I.GT.NBM(MTA,N+1)) I=NBM(MTA,N+1)
      J=LM(MTA,I)
      CR=LPP(MTA,J)/AKW
      IF (CR.GT.0.99) GO TO 140
      IF (RAND(0).GT.CR) GO TO 135
140  K=NB(MTB,N)*RAND(0)+1+NBM(MTB,N)
      IF (K.GT.NBM(MTB,N+1)) K=NBM(MTB,N+1)
      IF (MTA.EQ.MTB.AND.I.EQ.K) GO TO 140
      L=LM(MTB,K)
      CR=LPP(MTB,L)/AKW
      IF (CR.GT.0.99) GO TO 145
      IF (RAND(0).GT.CR) GO TO 140
145  CONTINUE
      GM1=WA(MTA)*PAU(MTA,J)+WA(MTB)*PAU(MTB,L)
      GM2=WA(MTA)*PAV(MTA,J)+WA(MTB)*PAV(MTB,L)
      GM3=WA(MTA)*PAW(MTA,J)+WA(MTB)*PAW(MTB,L)
      G1=PAU(MTA,J)-PAU(MTB,L)
      G2=PAV(MTA,J)-PAV(MTB,L)
      G3=PAW(MTA,J)-PAW(MTB,L)
      GS=G1**2+G2**2+G3**2
      IF (GS.LT.1.0E-8) GO TO 130
      ET=RM*GS
      EI=ER(MTA,J)+ER(MTB,L)
      VR=GS**(.5-ETT/2.)
      IF (VR.GE.CM(MTA,MTB,1)) GO TO 160
      CN(MTA,MTB,1)=CN(MTA,MTB,1)+VR
      IF (CN(MTA,MTB,1).LT.CM(MTA,MTB,1)) GO TO 130
      CN(MTA,MTB,1)=CN(MTA,MTB,1)-CM(MTA,MTB,1)
160  CONTINUE
      CPUT=ELTIME(0)
      IF (M.EQ.1196) WRITE(6,2) M,CPUT,VR,G1,G2,G3,EI
      CALL IMPACT(RM,G1,G2,G3,ET,EI,PHT,CHT,ETT,CM(MTA,MTB,2),CN(MTA,MTB
1,2))
165  CONTINUE
      CPUT=ELTIME(0)
      IF (M.EQ.1196) WRITE(6,3) M,CPUT,VR,G1,G2,G3,EI
      IF (PHT.EQ.0.) GO TO 175
      X1=0.0
      IF (CHI(MTA).EQ.-1.) GO TO 175
      X1=1.0
      IF (CHI(MTB).EQ.-1.) GO TO 175

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ORIGINAL PAGE 19
OF POOR QUALITY

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170 X1=Rand(0)
  IF ((CHI(MTA).EQ.0.).AND.(CHI(MTB).EQ.0.)) GO TO 175
  XT=X1*CHI(MTA)*(1.-X1)*CHI(MTB)
  IF(XT.GT.CN(MTA,MTB,3)) GO TO 175
  CN(MTA,MTB,3)=CN(MTA,MTB,3)+XT
  IF(CN(MTA,MTB,3).LT.CN(MTA,MTB,3)) GO TO 170
  CN(MTA,MTB,3)=CN(MTA,MTB,3)-CN(MTA,MTB,3)
175 CONTINUE
  C=DBA(MTA,N)
  D=DBA(MTB,N)
  IF(C.EQ.0.0) C=DB(MTA,N)
  IF(D.EQ.0.0) D=DB(MTB,N)
  IF(T(MTA,MTB,N).GE.AIME) GO TO 180
  PAU(MTA,J)=GM1+WA(MTB)*G1
  PAV(MTA,J)=GM2+WA(MTB)*G2
  PAW(MTA,J)=GM3+WA(MTB)*G3
  IF(PHT.GT.0.) ER(MTA,J)=X1*EI
  LCOL(MTA,J)=1
  NCOL(MTA,MTB)=NCOL(MTA,MTB)+1
  T(MTA,MTB,N)=T(MTA,MTB,N)+CN8(MTA,MTB)*LPF(MTA,J)/NBF(MTA,N)/D/VR
  IF(MTA.EQ.MTB) GO TO 190
180 IF(T(MTB,MTA,N).GE.AIME) GO TO 210
190 PAU(MTB,L)=GM1-WA(MTA)*G1
  PAV(MTB,L)=GM2-WA(MTA)*G2
  PAW(MTB,L)=GM3-WA(MTA)*G3
  IF(PHT.GT.0.) ER(MTB,L)=(1.-X1)*EI
  LCOL(MTB,L)=1
  NCOL(MTB,MTA)=NCOL(MTB,MTA)+1
  T(MTB,MTA,N)=T(MTB,MTA,N)+CN8(MTB,MTA)*LPF(MTB,L)/NBF(MTB,N)/C/VR
210 CONTINUE
  IF(M.EQ.1196) WRITE(6,1) AIME,T(MTA,MTB,N),T(MTB,MTA,N),NBF(MTA,N)
  1,NBF(MTB,N)
  IF(T(MTA,MTB,N).LT.AIME.OR.T(MTB,MTA,N).LT.AIME) GO TO 120
220 CONTINUE
230 CONTINUE
240 CONTINUE
  RETURN
  END

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ORIGINAL PAGE IS
OF POOR QUALITY

COL0810

COL0920

MOV0060
MOV0070
MOV0080
MOV0090

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SUBROUTINE MOVE(KSWCH,AKN,MJ,NS,NWEDG,THETAZ,XSTART,I2,I3,I4,I5,DEMOV0010
1LANG,NWEDGE,BTA,C2,C3,DPA,PL,HTI,HTR,JNT,KNM,NM,XCB,XLIM,MS,IWS,NTMOV0020
2CF,NTCV,PV,CTI,CTR,CNI,CNR,ALPHA,SIGMA,COEFF,HTS,HTSI,NTS,NTSF,UTLMOV0030
3,UTT,VTS,PAU,PAV,PAW,PAX,PAY,PAZ,LPF,LCOL,TB,IP,ER,CHI,CNG,CMG,I,
4UTLI,UTTI,VTSI)
  INTEGER*2 LPF(I,1),LCOL(I,1)
  INTEGER SWITCH,TIME
  LOGICAL DUMP
  REAL LAM,MU,NU
  DIMENSION DELANG(1),NWEDGE(1),BTA(1),C2(1),C3(1),PL(1),HTI(1)
  DIMENSION HTR(1),TB(1),XCB(1),ALPHA(3,1),SIGMA(3,1),COEFF(4,1)
  DIMENSION PAU(1,1),PAV(1,1),PAW(1,1),CTI(3,1),CTR(3,1)
  DIMENSION CNI(3,1),CNR(3,1),DPA(1),JNT(1),XLIM(1),KNM(1),NM(1)
  DIMENSION HTS(3,I2,I3),HTSI(3,I2,I3),NTS(3,I2,I3),NTSF(3,I2,I3)
  DIMENSION UTLI(3,I2,I3),UTTI(3,I2,I3),VTSI(3,I2,I3)
  DIMENSION UTL(3,I2,I3),UTT(3,I2,I3),VTS(3,I2,I3),NTCF(3,I4)

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DIMENSION PAX(I,1),PAY(I,1),PAZ(I,1),IWS(1),MS(1)

DIMENSION NTCV(3,I4,2,I5,3),FV(3,I4,2,I5,3)

DIMENSION ER(I,1),CNG(1),CHG(1),CHI(1)

COMMON /THIRD/PI,NREG

MOV0170

COMMON /FORTH/NBX,RM,XR,DUMP

MOV0180

COMMON /FIFTH/ND,TIME,DTM

MOV0190

COMMON /SVNTH/LAM,MU,NU,MT,N,J,XI,YI,ZI,TUSE

MOV0200

NAMELIST/CHECK/TIME,X,Y,Z,DX,DY,DZ,TLEFT,RADS,RMS,XR

MOV0210

THE PURPOSE OF THIS SUBROUTINE IS TO ADVANCE THE SPATIAL POSITION
OF ALL THE MOLECULES BY AN AMOUNT APPROPRIATE TO THEIR CURRENT VELOCITIES
LOCITIES AND THE PRE-SELECTED COLLISION TIME.

MOV0220

MOV0230

MOV0240

MOV0250

2 FORMAT(27H SOMETHING IS WRONG IN MOVE/3E20.7,4I7,E20.7)

MOV0260

MOV0270

NAREA=NREG+3

RMS=RM**2

DO 150 MT=1,IP

N=KNM(MT)

MOV0290

10 N=N+1

MOV0300

TLEFT=DTM

MOV0310

IF(KSWCH.EQ.1) TLEFT=TLEFT*RAND(0)

MOV0320

IF(N.GT.NM(MT)) GO TO 150

MOV0330

15 LAM=PAU(MT,N)

MOV0340

IF(LAM.EQ.0.) LAM=.0000001

MOV0350

MU=PAV(MT,N)

MOV0360

NU=PAW(MT,N)

MOV0370

XI=PAX(MT,N)

MOV0380

YI=PAY(MT,N)

MOV0390

ZI=PAZ(MT,N)

MOV0400

DX=TLEFT*LAM

MOV0410

DY=TLEFT*MU

MOV0420

DZ=TLEFT*NU

MOV0430

X=XI+DX

MOV0440

Y=YI+DY

MOV0450

Z=ZI+DZ

MOV0460

RADS=Y**2+Z**2

IF((RADS.GT.2.*RMS).OR.(ABS(X).GT.2.*XR)) WRITE(6,CHECK)

IF(RADS.GT.RMS) GO TO 100

RAD=SQRT(RMS)

KEY=ABS(LAM)/LAM+.5

MOV0490

DO 60 L=1,NAREA

MOV0500

IF(XI-XLIM(L)) 65,55,60

MOV0510

55 J=L+KEY-2

MOV0520

GO TO 70

MOV0530

60 CONTINUE

MOV0540

WRITE(6,2) DTM,XI,LAM,N,KEY,NAREA,L,XLIM(L)

MOV0550

IF(DUMP) CALL ABEND(16)

MOV0560

STOP

MOV0570

65 J=L-2

MOV0580

70 K=J+KEY+1

MOV0590

IF((K.EQ.0).OR.(K.EQ.NAREA+1)) GO TO 100

MOV0600

TUSE=(XLIM(K)-XI)/LAM

MOV0610

XTEMP=XLIM(K)

MOV0620

IF(TUSE.LE.TLEFT) GO TO 75

MOV0630

TUSE=TLEFT

MOV0640


```

      XTEMP=XI+TUSE*LAM
75  IF ((J.EQ.0).OR.(J.EQ.NREG+1)) GO TO 85
      CALL INTERS(AKN,MJ,NS,NWEDG,SWTCH,THETAZ,XSTART,I2,I3,I4,I5,DELANG
1,NWEDGE,BTA,C2,C3,DFA,FL,HTI,HTR,JNT,TB,XCB,CTI,CTR,CNI,CNR,ALPHA,
2SIGMA,COEPP,HTS,HTSI,NTS,NTSP,UTL,UTT,VTS,MS,IWS,NTCP,NTCV,PV,PAU,
3PAV,PAW,LPP,LCOL,IP,ER,CHI,CNG,CNG,I,UTLI,UTTI,VTSI)
      IF(SWTCH.EQ.1) GO TO 90
85  XI=XTEMP
      YI=YI+TUSE*NU
      ZI=ZI+TUSE*NU
90  PAX(MT,N)=XI
      PAY(MT,N)=YI
      PAZ(MT,N)=ABS(ZI)
      PAW(MT,N)=ABS(ZI)/ZI*PAW(MT,N)
      TLEPT=TLEPT-TUSE
      IF(TLEPT.GT.0.) GO TO 15
      GO TO 10
100 NZ=NM(MT)
      PAX(MT,N)=PAX(MT,NZ)
      PAY(MT,N)=PAY(MT,NZ)
      PAZ(MT,N)=PAZ(MT,NZ)
      PAU(MT,N)=PAU(MT,NZ)
      PAV(MT,N)=PAV(MT,NZ)
      PAW(MT,N)=PAW(MT,NZ)
      ER(MT,N)=ER(MT,NZ)
      LPP(MT,N)=LPP(MT,NZ)
      LCOL(MT,N)=LCOL(MT,NZ)
      N=N-1
      NM(MT)=NM(MT)-1
      GO TO 10
150 CONTINUE
      RETURN
      END

```

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SUBROUTINE ACCUM(I2,I3,FNB,NB,PAU,PAV,PAW,ER,TMP,TRP,XV,YV,ZV,LM,
1IP,I,LPP,NBF,NBM)
  INTEGER*2 LM(I,1),LPP(I,1),NBM(I,1)
  INTEGER*2 NB,NBF
  DIMENSION FNB(1),NB(1,1),PAU(1,1),PAV(1,1),PAW(1,1),TMP(1,1)
  DIMENSION XV(1,1),YV(1,1),ZV(1,1),ER(1,1),TRP(1,1),NBF(1,1)
  COMMON /FORTH/NBX

```

 THE PURPOSE OF THIS SUBROUTINE IS TO ACCUMULATE TEMPERATURES,
 VELOCITIES, AND DENSITIES IN VARIOUS ARRAYS FOR DETERMINING THE
 AVERAGE FLOW FIELD PROPERTIES AFTER STEADY-STATE HAS BEEN REACHED

```

N=0
DO 180 K=1,NBX
  IF(FNB(K).LE.0.) GO TO 180
  N=N+1
DO 110 MT=1,IP
  XV(MT,N)=0.0
  YV(MT,N)=0.0
  ZV(MT,N)=0.0
  TMP(MT,N)=0.

```

MOV0650
 MOV0660
 MOV0670
 MOV0680
 MOV0690
 MOV0710
 MOV0720
 MOV0730
 MOV0740
 MOV0750
 MOV0760
 MOV0770
 MOV0780
 MOV0790
 MOV0800
 MOV0810
 MOV0820
 MOV0830
 MOV0840
 MOV0850
 MOV0860
 MOV0870
 MOV0880
 MOV0890
 MOV0900
 MOV0910
 MOV0920
 MOV0930
 MOV0940
 MOV0950
 MOV0960
 ACUM050
 ACUM060
 ACUM070
 ACUM080
 ACUM090
 ACUM100
 ACUM110
 ACUM120
 ACUM130
 ACUM140
 ACUM160
 ACUM170
 ACUM180
 ACUM190

TRP(MT,N)=0.0	
TTX=0.	ACUM200
TTY=0.	ACUM210
TTZ=0.	ACUM220
TTR=0.0	
M=NB(MT,N)	ACUM230
IF(M.LT.1) GO TO 110	ACUM240
U=0.	ACUM250
V=0.	ACUM260
W=0.	ACUM270
DO 100 L=1,M	ACUM280
NA=NB(MT,N)+L	
J=LM(MT,NA)	
AKW=LPP(MT,J)	
PU=PAU(MT,J)	ACUM300
PV=PAV(MT,J)	ACUM310
PW=PAW(MT,J)	ACUM320
U=U+PU*AKW	
V=V+PV*AKW	
W=W+PW*AKW	
TTR=TTR+ER(MT,J)*AKW	
TTX=TTX+PU*PU*AKW	
TTY=TTY+PV*PV*AKW	
20 TTZ=TTZ+PW*PW*AKW	
M=NB(MT,N)	
XV(MT,N)=U/M	ACUM390
YV(MT,N)=V/M	ACUM400
ZV(MT,N)=W/M	ACUM410
TMP(MT,N)=(TTX+TTY+TTZ)/M	ACUM420
TRP(MT,N)=TTR/M	
10 CONTINUE	ACUM430
30 CONTINUE	ACUM440
RETURN	ACUM450
END	ACUM460

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SUBROUTINE AVRGE(FNB,DB,DBA,NB,NBT,XV,YV,ZV,XVA,YVA,ZVA,TMP,TMPA,
1TRP,TRPA,IP,I,NBF,NBS)
INTEGER*2 NB,NBT,NBF,NBS
DIMENSION FNB(1),DB(1),DBA(1),NB(1),NBT(1),TMP(1),
DIMENSION TMPA(1),XV(1),XVA(1),YV(1),YVA(1),ZV(1),
DIMENSION ZVA(1),TRP(1),TRPA(1),NBF(1),NBS(1)
COMMON /FORTH/NBX

AVG0050

THE PURPOSE OF THIS SUBROUTINE IS TO COMPUTE THE AVERAGE FLOW
FIELD PROPERTIES.

AVG0060

AVG0070

AVG0080

AVG0090

N=0

AVG0100

DO 110 M=1,NBX

AVG0110

IF(FNB(M).LE.0.) GO TO 110

AVG0120

N=N+1

AVG0130

DO 100 MT=1,IP

A=NBT(MT,N)

AVG0150

B=NBF(MT,N)

C=A+B

AVG0170

NBT(MT,N)=C

AVG0180


```

NBS(MT,N)=NBS(MT,N)+NB(MT,N)
IF(C.LT.1.) GO TO 100
DBA(MT,N)=(DBA(MT,N)*A+DB(MT,N)*B)/C
XVA(MT,N)=(XVA(MT,N)*A+XV(MT,N)*B)/C
YVA(MT,N)=(YVA(MT,N)*A+YV(MT,N)*B)/C
ZVA(MT,N)=(ZVA(MT,N)*A+ZV(MT,N)*B)/C
TMPA(MT,N)=(TMPA(MT,N)*A+TMP(MT,N)*B)/C
TRPA(MT,N)=(TRPA(MT,N)*A+TRP(MT,N)*B)/C
100 CONTINUE
110 CONTINUE
RETURN
END

SUBROUTINE DRAG(AKN,MJ,NS,NWEDG,THETAZ,XSTART,I2,I3,I4,I5,DELANG,
1NWEDGE,BTA,C2,C3,DFA,FL,HTI,HTR,TB,XCB,CTI,CTR,CNI,CNR,ALPHA,SIGMA
2,COEFF,HTS,HTSI,NTS,NTSF,UTL,UTT,VTS,MS,IWS,NTCF,NTCV,PV,PAU,PAV,P
3AW,LPP,LCOL,IP,ER,CHI,CNG,CMG,I,UTLI,UTTI,VTSI)
INTEGER*2 LPP,LCOL
INTEGER TIME,TST
REAL LAM,MU,NU,JAY,KAY
DIMENSION DELANG(1),NWEDGE(1),BTA(1),C2(1),C3(1),FL(1),HTI(1)
DIMENSION HTR(1),TB(1),XCB(1),ALPHA(3,1),SIGMA(3,1),COEFF(4,1)
DIMENSION LPP(1,1),PAU(1,1),PAV(1,1),PAW(1,1),CTI(3,1),CTR(3,1)
DIMENSION HTS(3,I2,I3),HTSI(3,I2,I3),NTS(3,I2,I3),NTSF(3,I2,I3)
DIMENSION UTL(3,I2,I3),UTT(3,I2,I3),VTS(3,I2,I3),LCOL(1,1)
DIMENSION UTLI(3,I2,I3),UTTI(3,I2,I3),VTSI(3,I2,I3)
DIMENSION CNI(3,1),CNR(3,1),DFA(1),IWS(1),MS(1),NTCF(3,I4)
DIMENSION NTCV(3,I4,2,I5,3),PV(3,I4,2,I5,3)
DIMENSION ER(1,1),CNG(1),CMG(1),CHI(1)
COMMON /THIRD/PI
COMMON /FIFTH/ND,TIME,DTM,TI,ITS,ITP,TST
COMMON /SVNTH/LAM,MU,NU,MT,N,J,XCL,YCL,ZCL
-----
THE PURPOSE OF THIS SUBROUTINE IS TO ACCUMULATE THE DRAG AND HEAT
TRANSFER INCREMENTS ON THE BODY CONTRIBUTED BY EACH MOLECULE WHICH
COLLIDES WITH THE BODY. IN ADDITION, EACH MOLECULE WHICH COLLIDES
WITH THE BODY IS ASSIGNED AN APPROPRIATE NEW VELOCITY (OF REFLECT
WHICH IS USED TO CONTINUE ITS SPATIAL TRANSLATION (IN SUBROUTINE
-----
CALL NORMAL(EYE,JAY,KAY,ONE,COEFF)
AKW=LPP(MT,N)
RAD=SQRT(YCL*YCL+ZCL*ZCL)
ARG=YCL/RAD
TANG=180.*(1.-ARCCOS(ARG)/PI)
IWDG=TANG/DELANG(1)+1.
IF((IWDG.GT.NWEDGE(1)).AND.(DELANG(2).NE.0.)) IWDG=(TANG-THETAZ)
1/DELANG(2)+NWEDGE(1)+1
IF(IWDG.LT.1) IWDG=1
IF(IWDG.GT.NWEDG) IWDG=NWEDG
D=(LAM*LAM+MU*MU+NU*NU)*AKW
G=ER(MT,N)*AKW
H=G
DO 100 M=1,ND
IF(XCL.LT.XCB(M)) GO TO 110
100 CONTINUE

```

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AVG0190
AVG0200
AVG0210
AVG0220
AVG0230
AVG0240
AVG0250
AVG0260
AVG0270
AVG0280

DRG0040
DRG0050
DRG0060
DRG0070

DRG0140
DRG0150
DRG0160
DRG0170

DRG0180
DRG0190
DRG0200
DRG0210
DRG0220

DRG0230
DRG0240
DRG0250
DRG0260
DRG0270
DRG0280
DRG0290

DRG0310
DRG0320

DRG0330
DRG0340
DRG0350

110 UI=LAM
 WI=(NU*JAY-MU*KAY)/ONE
 VID=LAM*EYE+MU*JAY+NU*KAY
 UID=LAM*ONE-EYE*(MU*JAY+NU*KAY)/ONE
 E=Rand(0)
 IT(E.LT.SIGMA(MT,M)) GO TO 115
 VRD=-VID
 URD=UID
 WR=WI
 GO TO 125
 115 V=4.*Rand(0)
 C2(MT)=C2(MT)+.544331*V*V*V*EXP(1.5-V*V)
 IF(C2(MT).LT.1.) GO TO 115
 C2(MT)=C2(MT)-1.
 IF(NTSF(MT,M,IWDG).NE.0) GO TO 117
 ATR=ALPHA(MT,M)*TB(M)/SIGMA(MT,M)
 GO TO 118
 117 ATR=ALPHA(MT,M)*TB(M)/SIGMA(MT,M)+(1.-ALPHA(MT,M)/SIGMA(MT,M))*HTSDRG0530
 118 ABR=SQRT(ATR)
 V=V*ABR/BTA(MT)
 120 A=Rand(0)
 C3(MT)=C3(MT)+A
 IF(C3(MT).LT.1.) GO TO 120
 C3(MT)=C3(MT)-1.
 B=SQRT(1.-A*A)
 C=2.*PI*Rand(0)
 VRD=V*A
 URD=V*B*COS(C)
 WR=V*B*SIN(C)
 IF(CHI(MT).EQ.-1.) GO TO 125
 122 X=9.*Rand(0)
 IF(X.EQ.0.0) GO TO 122
 XTEMP=1.0
 IF(CHI(MT).NE.0.0) XTEMP=X**CHI(MT)
 CNG(MT)=CNG(MT)+XTEMP*EXP(-X)
 IF(CNG(MT).LT.CMG(MT)) GO TO 122
 124 CONTINUE
 CNG(MT)=CNG(MT)-CMG(MT)
 IF(CNG(MT).GE.CMG(MT)) GO TO 124
 ER(MT,N)=X*ATR
 H=ER(MT,N)*AKW
 125 UR=EYE*VRD+ONE*URD
 PAU(MT,N)=UR
 PAV(MT,N)=JAY*VRD-(KAY*WR+EYE*JAY*URD)/ONE
 PAW(MT,N)=KAY*VRD+(JAY*WR-EYE*KAY*URD)/ONE
 IF(TIME.LE.TST) RETURN
 IF(TI.GT.0.) GO TO 130
 TI=TST*DTM
 130 XMZ=(XCL-XSTART)*AKN
 YMZ=RAD*AKN
 B=(URD*URD+VRD*VRD+WR*WR)*AKW
 UTI=UID*UID+WI*WI
 UYI=- (UID*EYE*JAY+WI*KAY)/ONE
 UYR=- (URD*EYE*JAY+WR*KAY)/ONE

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DRG0360
 DRG0370
 DRG0380
 DRG0390
 DRG0400
 DRG0410
 DRG0420
 DRG0430
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 DRG0470
 DRG0480
 DRG0490
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 DRG0680
 DRG0690
 DRG0710
 DRG0730
 DRG0740
 DRG0750
 DRG0760
 DRG0770
 DRG0780

```

FL(MT)=FL(MT)+AKW*DPA(MT)
HTI(MT)=HTI(MT)+D+G
HTR(MT)=HTR(MT)-B-H
CTI(MT,1)=CTI(MT,1)+UID*ONE*AKW
CTI(MT,2)=CTI(MT,2)+UYI*AKW
CTI(MT,3)=CTI(MT,3)+(XMZ*UYI-YMZ*UID*ONE)*AKW
CNI(MT,1)=CNI(MT,1)+VID*EYE*AKW
CNI(MT,2)=CNI(MT,2)+VID*JAY*AKW
CNI(MT,3)=CNI(MT,3)+(XMZ*JAY-YMZ*EYE)*VID*AKW
CTR(MT,1)=CTR(MT,1)-URD*ONE*AKW
CTR(MT,2)=CTR(MT,2)-UYR*AKW
CTR(MT,3)=CTR(MT,3)-(XMZ*UYR-YMZ*URD*ONE)*AKW
CNR(MT,1)=CNR(MT,1)-VRD*EYE*AKW
CNR(MT,2)=CNR(MT,2)-VRD*JAY*AKW
CNR(MT,3)=CNR(MT,3)-(XMZ*JAY-YMZ*EYE)*VRD*AKW
NTS(MT,M,IWDG)=NTS(MT,M,IWDG)+1
NTSF(MT,M,IWDG)=NTSF(MT,M,IWDG)+AKW
UTLI(MT,M,IWDG)=UTLI(MT,M,IWDG)+UID*AKW
UTL(MT,M,IWDG)=UTL(MT,M,IWDG)+(UID-URD)*AKW
UTTI(MT,M,IWDG)=UTTI(MT,M,IWDG)+WI*AKW
UTT(MT,M,IWDG)=UTT(MT,M,IWDG)+(WI-WR)*AKW
VTSI(MT,M,IWDG)=VTSI(MT,M,IWDG)-VID*AKW
VTS(MT,M,IWDG)=VTS(MT,M,IWDG)+(VRD-VID)*AKW
HTSI(MT,M,IWDG)=HTSI(MT,M,IWDG)+D+G
HTS(MT,M,IWDG)=HTS(MT,M,IWDG)+D-B+G-H
IF(NS.EQ.0) RETURN
NC=(2*LCOL(MT,N))/(1+LCOL(MT,N))+1
DO 160 L=1,NS
MTEST=MS(L)
IWT=IWS(L)
IF((M.NE.MTEST).OR.(IWDG.NE.IWT)) GO TO 160
IF(NC.EQ.1) NTCF(MT,L)=NTCF(MT,L)+AKW
DO 145 JJ=1,MJ
IF(ABS(VID).GT.FV(MT,L,NC,JJ,1)) GO TO 135
NTCV(MT,L,NC,JJ,1)=NTCV(MT,L,NC,JJ,1)+AKW
135 IF(UID.GT.FV(MT,L,NC,JJ,2)) GO TO 140
NTCV(MT,L,NC,JJ,2)=NTCV(MT,L,NC,JJ,2)+AKW
140 IF(WI*ABS(ZCL)/ZCL.GT.FV(MT,L,NC,JJ,3)) GO TO 145
NTCV(MT,L,NC,JJ,3)=NTCV(MT,L,NC,JJ,3)+AKW
145 CONTINUE
160 CONTINUE
RETURN
END

```

SUBROUTINE INTERS(AKN,MJ,NS,NWEDG,SWTCH,THETAZ,XSTART,I2,I3,I4,I5,INT0010
 1DELANG,NWEDGE,BTA,C2,C3,DPA,FL,HTI,HTR,JNT,TB,XCB,CTI,CTR,CNI,CNR,INT0020
 2ALPHA,SIGMA,COEFF,HTS,HTSI,NTS,NTSF,UTL,UTT,VTS,MS,IWS,NTCF,NTCV,PINT0030
 3V,PAU,PAV,PAW,LPP,LCOL,IP,ER,CHI,CNG,CMG,I,UTLI,UTTI,VTSI)
 INTEGER*2 LPP,LCOL
 INTEGER SWTCH
 REAL LAM,MU,NU,LINEAR
 DIMENSION DELANG(1),NWEDGE(1),BTA(1),C2(1),C3(1),FL(1),HTI(1)
 DIMENSION HTR(1),TB(1),XCB(1),ALPHA(3,1),SIGMA(3,1),COEFF(4,1)
 DIMENSION LPP(1,1),PAU(1,1),PAV(1,1),PAW(1,1),CTI(3,1),CTR(3,1)
 DIMENSION HTS(3,I2,I3),HTSI(3,I2,I3),NTS(3,I2,I3),NTSF(3,I2,I3)

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DRG0790

DRG0820

DRG0830

DRG0840

DRG0850

DRG0860

DRG0870

DRG0880

DRG0890

DRG0900

DRG0910

DRG0920

DRG0930

DRG0940

DRG0950

DRG0960

DRG0970

DRG0980

DRG1010

DRG1020

DRG1060

DRG1080

DRG1150

DRG1160

DRG1170

DRG1180

INT0050

INT0060

INT0070

INT0080

```

DIMENSION UTL(3,I2,I3),UTT(3,I2,I3),VTS(3,I2,I3),LCOL(I,1)
DIMENSION UTLI(3,I2,I3),UTTI(3,I2,I3),VTSI(3,I2,I3)
DIMENSION CNI(3,1),CNR(3,1),DPA(1),IWS(1),MS(1),NTCP(3,I4)
DIMENSION JNT(1),NTCV(3,I4,2,I5,3),PV(3,I4,2,I5,3)
DIMENSION ER(I,1),CNG(1),CMG(1),CHI(1)
COMMON /SVNTH/LAM,MU,NU,MT,N,J,XI,YI,ZI,TUSE

```

INT0150

```

-----INT0160
THE PURPOSE OF THIS SUBROUTINE IS TO DETERMINE FOR EACH MOLECULARINT0170
TRAJECTORY IF THERE IS AN INTERSECTION OF THE TRAJECTORY WITH THEINT0180
BODY SURFACE.
-----INT0190

```

INT0200

SWTCH=0

INT0210

A=COEPP(1,J)

INT0220

B=COEPP(2,J)

INT0230

C=COEPP(3,J)

INT0240

D=COEPP(4,J)

INT0250

ONE=A*LAM

INT0260

TWO=B*NU

INT0270

TRE=B*NU

INT0280

PUR=.5*C*LAM

INT0290

SQUARE=ONE*LAM+TWO*NU+TRE*NU

INT0300

LINEAR=ONE*X1+TWO*Y1+TRE*Z1+PUR

INT0310

CONST=(A*X1+C)*X1+B*(Y1*Y1+Z1*Z1)+D

INT0320

IF(SQUARE.EQ.0.) GO TO 150

INT0330

DISCR=LINEAR*LINEAR-SQUARE*CONST

INT0340

IF(DISCR.LT.0.) RETURN

INT0350

SDISC = SQRT(DISCR)

INT0360

TYME = (-LINEAR-SDISC)/SQUARE

INT0370

GO TO 250

INT0380

150 IF(LINEAR.EQ.0.) RETURN

INT0390

TYME=-.5*CONST/LINEAR

INT0400

250 IF(TYME.GT.TUSE) RETURN

INT0410

IF(TYME.LE.0.) RETURN

INT0420

JNT(MT)=JNT(MT)+1

INT0430

X1=X1+LAM*TYME

INT0440

Y1=Y1+NU*TYME

INT0450

Z1=Z1+NU*TYME

INT0460

CALL DRAG(AKN,MJ,NS,NWEDG,THETAZ,XSTAR,I2,I3,I4,I5,DELANG,NWEDGE,INT0470

1BTA,C2,C3,DPA,PL,HTI,HTR,TB,XCB,CTI,CTR,CNI,CNR,ALPHA,SIGMA,COEPP,INT0480

2HTS,HTSI,NTS,NTSP,UTL,UTT,VTS,MS,IWS,NTCP,NTCV,PV,PAU,PAV,PAW,LPP,INT0490

3LCOL,IP,ER,CHI,CNG,CMG,I,UTLI,UTTI,VTSI)

LCOL(MT,N)=2

INT0510

TUSE=TYME

INT0520

SWTCH=1

INT0530

RETURN

INT0540

END

INT0550

SUBROUTINE NORMAL(EYE,JAY,KAY,ONE,COEPP)

NORM010

REAL LAM,MU,NU,JAY,KAY

NORM020

DIMENSION COEPP(4,1)

NORM030

COMMON /SVNTH/LAM,MU,NU,MT,N,J,XCL,YCL,ZCL

NORM040

DFDX=2.*COEPP(1,J)*XCL+COEPP(3,J)

NORM050

DFDY=2.*COEPP(2,J)*YCL

NORM060

DFDZ=2.*COEPP(2,J)*ZCL

NORM070

DENOM=SQRT(DFDX*DFDX+DFDY*DFDY+DFDZ*DFDZ)

NORM080

ORIGINAL PAGE 13
OF POOR QUALITY

```

EYE=DFDX/DENOM
JAY=DFDY/DENOM
KAY=DFDZ/DENOM
ONE=SQRT(JAY*JAY+KAY*KAY)
RETURN
END

```

ORIGINAL PAGE 19
OF POOR QUALITY

NORM090
NORM100
NORM110
NORM120
NORM130
NORM140

```

FUNCTION ARCCOS(ARG)
COMMON /THIRD/PI
IF (ARG) 30, 10, 20
10 A=.5*PI
GO TO 40
20 A=ATAN(SQRT(1.-ARG*ARG)/ARG)
GO TO 40
30 A=PI+ATAN(SQRT(1.-ARG*ARG)/ARG)
40 ARCCOS=A
RETURN
END

```

ARCS010
ARCS020
ARCS030
ARCS040
ARCS050
ARCS060
ARCS070
ARCS080
ARCS090
ARCS100
ARCS110

```

FUNCTION EERF(SS)
EERF=ERFC(-SS)
RETURN
END

```

EERF010
EERF050
EERF060

```

SUBROUTINE PRINT1(DT,COSANG,SINANG,RMA,RNU,DRF,PCF,HTP,PL,HTI,HTR,
1CTI,CTR,CNI,CNR)
DIMENSION DD(3),WD(2,5),PP(4,4),QQ(4,4),RR(4,4),SS(4,4),TT(4,4)
DIMENSION UU(4,4),P1(4,4),Q1(4,4),R1(4,4),PA(4),PB(4),PC(4)
DIMENSION PL(1),HTI(1),HTR(1),CTI(3,1),CTR(3,1),CNI(3,1),CNR(3,1)
DIMENSION RMA(1),RNU(1)
DATA WD/'X-PO','RCE','Y-PO','RCE','Z-MO','MENT','DRAG',' ','LPT10070
1IPT',' '/

```

PT10080
PT10090
PT10100
PT10110
PT10120
PT10130
PT10140
PT10150
PT10160
PT10170
PT10180

THE PURPOSE OF THIS SUBROUTINE IS TO PRINT OUT THE GROSS SURFACE
COEFFICIENTS OF THE BODY.

FORMATS

```

1 FORMAT(/1X,50('*'),' GROSS SURFACE COEFFICIENTS ',50('*')/' MOLEC
1ULAR WEIGHT',12X,F8.3,3(19X,F8.3)/25X,
2 'INC. REF. TOT. INC. REF. TOT. INC.
3REF. TOT. INC. REF. TOT.')
```

PT10280
PT10290
PT10300
PT10310

PT10280
PT10290
PT10300
PT10310

ORIGINAL PAGE 11
OF POOR QUALITY

```
RMR=0.0
DO 50 MT=1,3
DD (MT)=RMA (MT)*RNU (MT)*DRP/DT
50 RMR=RMR+RMA (MT)*RNU (MT)
WRITE (6,1) (RMA (MT),MT=1,3),RMR
PF=PL (1)*PCP/DT
QF=PL (2)*PCP/DT
RF=PL (3)*PCP/DT
SF=PF+QF+RF
```

```
WRITE (6,10) PF,QF,RF,SF
```

PT10400

```
DO 200 I=1,3
```

```
PP (4,I)=0.0
```

```
QQ (4,I)=0.0
```

```
RR (4,I)=0.0
```

```
SS (4,I)=0.0
```

```
TT (4,I)=0.0
```

```
UU (4,I)=0.0
```

```
P1 (4,I)=0.0
```

```
Q1 (4,I)=0.0
```

```
R1 (4,I)=0.0
```

```
DO 150 MT=1,3
```

```
PP (MT,I)=CTI (MT,I)*DD (MT)/RMR
```

```
QQ (MT,I)=CTR (MT,I)*DD (MT)/RMR
```

```
SS (MT,I)=CNI (MT,I)*DD (MT)/RMR
```

```
TT (MT,I)=CNR (MT,I)*DD (MT)/RMR
```

```
P1 (MT,I)=PP (MT,I)+SS (MT,I)
```

PT10460

```
Q1 (MT,I)=QQ (MT,I)+TT (MT,I)
```

PT10470

```
RR (MT,I)=PP (MT,I)+QQ (MT,I)
```

PT10480

```
UU (MT,I)=SS (MT,I)+TT (MT,I)
```

PT10490

```
R1 (MT,I)=P1 (MT,I)+Q1 (MT,I)
```

```
PP (4,I)=PP (4,I)+PP (MT,I)
```

```
QQ (4,I)=QQ (4,I)+QQ (MT,I)
```

```
RR (4,I)=RR (4,I)+RR (MT,I)
```

```
SS (4,I)=SS (4,I)+SS (MT,I)
```

```
TT (4,I)=TT (4,I)+TT (MT,I)
```

```
UU (4,I)=UU (4,I)+UU (MT,I)
```

```
P1 (4,I)=P1 (4,I)+P1 (MT,I)
```

```
Q1 (4,I)=Q1 (4,I)+Q1 (MT,I)
```

```
R1 (4,I)=R1 (4,I)+R1 (MT,I)
```

```
150 CONTINUE
```

```
WRITE (6,12) (WD (J,I),J=1,2),(PP (K,I),QQ (K,I),RR (K,I),K=1,4)
```

```
WRITE (6,14) (SS (K,I),TT (K,I),UU (K,I),K=1,4)
```

```
WRITE (6,16) (P1 (K,I),Q1 (K,I),R1 (K,I),K=1,4)
```

```
200 CONTINUE
```

PT10630

```
AA=COSANG
```

PT10640

```
BB=SINANG
```

PT10650

```
DO 300 I=4,5
```

PT10660

```
DO 250 K=1,4
```

```
PP (K,4)=AA*PP (K,1)+BB*PP (K,2)
```

PT10680

```
QQ (K,4)=AA*QQ (K,1)+BB*QQ (K,2)
```

PT10690

```
RR (K,4)=AA*RR (K,1)+BB*RR (K,2)
```

PT10700

```
SS (K,4)=AA*SS (K,1)+BB*SS (K,2)
```

PT10710

```
TT (K,4)=AA*TT (K,1)+BB*TT (K,2)
```

PT10720

```
UU (K,4)=AA*UU (K,1)+BB*UU (K,2)
```

PT10730

```
P1 (K,4)=AA*P1 (K,1)+BB*P1 (K,2)
```

PT10740

```

      Q1(K,4)=AA*Q1(K,1)+BB*Q1(K,2)
250  R1(K,4)=AA*R1(K,1)+BB*R1(K,2)
      WRITE(6,12) (WD(J,I),J=1,2), (PP(K,4),QQ(K,4),RR(K,4),K=1,4)
      WRITE(6,14) (SS(K,4),TT(K,4),UU(K,4),K=1,4)
      WRITE(6,16) (P1(K,4),Q1(K,4),R1(K,4),K=1,4)
      AA=-SINANG
      BB=COSANG
300  CONTINUE
      HD=HTF/DT
      PA(4)=0.0
      PB(4)=0.0
      PC(4)=0.0
      DO 400 MT=1,3
      PA(MT)=HTI(MT)*RMA(MT)*RNU(MT)*HD/RMR
      PB(MT)=HTR(MT)*RMA(MT)*RNU(MT)*HD/RMR
      PC(MT)=PA(MT)+PB(MT)
      PA(4)=PA(4)+PA(MT)
      PB(4)=PB(4)+PB(MT)
      PC(4)=PC(4)+PC(MT)
400  CONTINUE
      WRITE(6,18) (PA(I),PB(I),PC(I),I=1,4)
      RETURN
      END

```

PT10750
PT10760

PT10800
PT10810
PT10820
PT10830

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OF POOR QUALITY

PT10950
PT10960

```

SUBROUTINE PRINT2(AKN,XSTART,DT,RNU,RMA,DRF,PCF,HTF,UTLI,UTTI,VTSI
1,HTSI,DELANG,NWEDGE,XS,XCB,YCB,HTS,NTS,NTSF,UTL,UTT,VTS,I2,I3,IP)
DIMENSION RMA(1),RNU(1),DELANG(1),NWEDGE(1),XS(1),XCB(1),YCB(1)
DIMENSION HTS(3,I2,I3),NTS(3,I2,I3),NTSF(3,I2,I3),UTL(3,I2,I3)
DIMENSION UTT(3,I2,I3),VTS(3,I2,I3),UTLI(3,I2,I3),UTTI(3,I2,I3)
DIMENSION VTSI(3,I2,I3),HTSI(3,I2,I3)
COMMON /FIFTH/ND

```

PT20060

PT20070

THE PURPOSE OF THIS SUBROUTINE IS TO PRINT OUT THE DISTRIBUTION
ON SURFACE OF THE SURFACE COEFFICIENTS

PT20080

PT20090

PT20100

PT20110

PT20120

PT20130

PT20140

PT20150

PT20160

FORMATS

```

8  FORMAT(/1X,45('*'),' DISTRIBUTION ON SURFACE ',45('*')/71X,'INC.
1  TOT.      INC.      TOT.      INC.      TOT. '/11X,' SEGMENT GEOMETRY',
214X,'MOL.      MOLE  SAMP      NUM.      SKIN      SKIN      PRES-      PRES-
3  HEAT      HEAT'/' NO.  CENTER  DELX   CENTER  DELANG',4X,'WGHT.  F
4  RACT.',10X,
5  'PLUX  PRCTN  PRCTN  SURE  SURE  TRNSF  TRNSF')
10 FORMAT(1X,I3,F8.3,F7.3,F9.3,F8.3,1X,2F8.4,I6,7F8.4)
11 FORMAT(37X,F8.4,' 1.0000',I6,7F8.4)

```

PT20230

*****> PT20240

PT20250

PT20260

RMR=0.0

DO 50 MT=1,IP

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OF POOR QUALITY

50	RMR=RMR+RMA(MT)*RNU(MT)	
	WRITE(6,8)	
	I=0	PT20280
6	DO 110 N=1,ND	PT20320
	DTY=DT*YCB(N)/180.	PT20330
	P=XS(N)	PT20340
	Q=2.*(XCB(N)-XSTART)*AKN-XS(N)	PT20350
	ANGLE=0.	PT20360
	R=0.	PT20370
7	J=0	PT20380
	DO 105 L=1,2	
	R=R+.5*ANGLE	PT20400
	ANGLE=DELANG(L)	
	R=R-.5*ANGLE	PT20420
	ICNT=NWEDGE(L)	
6	IF(ICNT.EQ.0) GO TO 105	
	PMLT=PCF/(DTY*ANGLE)	PT20440
	QMLT=DRF/(DTY*ANGLE)	PT20450
	SMLT=HTF/(DTY*ANGLE)	PT20460
	DO 100 K=1,ICNT	PT20470
	R=R+ANGLE	PT20480
8	I=I+1	PT20490
	J=J+1	PT20500
	M3=0	
	P3=0.0	
	Q3=0.0	
	Q4=0.0	
8	R3=0.0	
	R4=0.0	
	S3=0.0	
	S4=0.0	
	DO 90 MT=1,IP	
	M1=NTS(MT,N,J)	
8	M3=M3+M1	
	P1=NTSF(MT,N,J)*PMLT*RNU(MT)	
	P3=P3+P1	
	Q1=SQRT(UTLI(MT,N,J)**2+UTTI(MT,N,J)**2)*RNU(MT)*RMA(MT)*QMLT/RMR	
	Q2=SQRT(UTL(MT,N,J)**2+UTT(MT,N,J)**2)*RNU(MT)*RMA(MT)*QMLT/RMR	
	Q3=Q3+Q1	
8	Q4=Q4+Q2	
	R1=VTSI(MT,N,J)*RNU(MT)*RMA(MT)*QMLT/RMR	
	R2= VTS(MT,N,J)*RNU(MT)*RMA(MT)*QMLT/RMR	
	S1=HTSI(MT,N,J)*RNU(MT)*RMA(MT)*SMLT/RMR	
	S2= HTS(MT,N,J)*RNU(MT)*RMA(MT)*SMLT/RMR	
	R3=R3+R1	
8	R4=R4+R2	
	S3=S3+S1	
	S4=S4+S2	
90	WRITE(6,10) I,P,Q,R,ANGLE,RMA(MT),RNU(MT),M1,P1,Q1,Q2,R1,R2,S1,S2	
	WRITE(6,11) RMR,M3,P3,Q3,Q4,R3,R4,S3,S4	
100	CONTINUE	PT20660
105	CONTINUE	
110	CONTINUE	PT20670
	RETURN	PT20680
	END	PT20690


```

SUBROUTINE PRINT3(      IP,MJ,NS,NWEDG,I2,I3,I4,I5, RMA,XS,IWS,
1MS,TANGN,NTSF,NTCP,NTCV,PV) PT30020
DIMENSION RMA(1),XS(1),IWS(1),MS(1),TANGN(1),NTSF(3,I2,I3)
DIMENSION NTCF(3,I4), NTCV(3,I4,2,I5,3),PV(3,I4,2,I5,3),QUO(3,3) PT3

```

FORMATS

```

2 FORMAT (//1X,40('*'),' MOMENTS OF INCIDENT DISTRIBUTION FUNCTIONS ' PT30100
1,40('*')) PT30110
6 FORMAT (/21X,I5,' UNCOLLIDED MOLECULES',F8.4,27X,I5,' COLLIDED MOLE
1CULES',F8.4/)
8 FORMAT (I5,4X,A3,12(1X,F9.4)) PT30170
10 FORMAT (12X,12(1X,F9.4)) PT30180
12 FORMAT (1H ) PT30190

```

***** PT30220

```

WRITE(6,2) PT30230
DO 155 I=1,NS PT30240
MR=MS(I)
ITT=IWS(I)
N=(MR-1)*NWEDG+ITT
DO 150 MT=1,IP
A=0.
B=1.
IC=NTCF(MT,I)
ID=NTSF(MT,MR,ITT)-IC
E=NTSF(MT,MR,ITT)
IF(E.LE.0.) GO TO 110
A=IC/E
B=1.-A
110 WRITE(6,6) IC,A,ID,B
E=NTCF(MT,I)
DO 121 NC=1,2
DO 120 K=1,3
QUO(NC,K)=0.
IF(E.EQ.0.) GO TO 120
QUO(NC,K)=NTCV(MT,I,NC,1,K)/E
120 CONTINUE
E=NTSF(MT,MR,ITT)-NTCF(MT,I)
121 CONTINUE
WRITE(6,8) N, RMA(MT),PV(MT,I,1,1,1),QUO(1,1),PV(MT,I,1,1,2),QUO
1(1,2),PV(MT,I,1,1,3),QUO(1,3),PV(MT,I,2,1,1),QUO(2,1),PV(MT,I,2,1,
22),QUO(2,2),PV(MT,I,2,1,3),QUO(2,3)
IF(MJ.EQ.1) GO TO 150
DO 140 J=2,MJ
E=NTCF(MT,I)
DO 131 NC=1,2
DO 130 K=1,3
QUO(NC,K)=0.
IF(E.EQ.0.) GO TO 130

```

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```

PT30260
PT30270
PT30280
PT30370
PT30380
PT30400
PT30450
PT30460
PT30470
PT30480
PT30490
PT30500
PT30510
PT30520
PT30530
PT30550
PT30560
PT30570
PT30580
PT30590
PT30600
PT30610
PT30620
PT30630

```

```

      QUO(NC,K)=NTCV(MT,I,NC,J,K)/E
130 CONTINUE
      E=NTSP(MT,MR,ITT)-NTCP(MT,I)
131 CONTINUE
140 WRITE(6,10)  PV(MT,I,1,J,1),QUO(1,1),PV(MT,I,1,J,2),QUO(1,2),PV(MT,I,1,J,3),QUO(1,3),
      PV(MT,I,2,J,1),QUO(2,1),PV(MT,I,2,J,2),QUO(2,2),PV(MT,I,2,J,3),QUO(2,3)
150 WRITE(6,12)
155 CONTINUE
      RETURN
      END

```

ORIGINAL PAGE 13
OF POOR QUALITY

PT30640
PT30650
PT30660
PT30670
PT30680
PT30690
PT30700
PT30710
PT30720
PT30730

```

SUBROUTINE PRINT4(MSP,CHI,RNU,I,TRP,NUMCEL,FDN,WTM,DB,NS,TMP,XV,
1YV,ZV,KS,NB,IC,YC,ZC,LEV,LKW)
  INTEGER*2 LKW(1)
  INTEGER*2 NB,NUMCEL,NS
  DIMENSION FDN(1),RNU(1),CHI(1),WTM(1),NUMCEL(1),TMP(I,1),TRP(I,1),
  DB(I,1),NB(I,1),XV(I,1),YV(I,1),ZV(I,1),DBT(3),NS(I,1),
  IC(1),YC(1),ZC(1),LEV(1)
  COMMON /PORTH/NBX

```

PT50050
PT50060
PT50070
PT50080
PT50090
PT50100
PT50110
PT50120
PT50130
PT50140
PT50150
PT50160
PT50170

THE PURPOSE OF THIS SUBROUTINE IS TO PRINT OUT THE INSTANTANEOUS
FLOW-FIELD PROPERTIES.

FORMATS

```

1 FORMAT(/1X,45('*'),' INSTANTANEOUS FLOW FIELD INFORMATION ',45('*')
1'))
2 FORMAT(/2X,'LEVEL=',I3,3X,'WEIGHTING FACTOR(MAX)=',I3,3X,'WEDGE AN
1GLE =',F7.2,' DEGREES',3X,' RADIAL POSITION =',E11.3/2X,'BOX# X PO
2SITION SAMP DENSITY MACH NO X VEL. Y VEL. Z VEL. T(KIN) T(RO
3T) TEMP.',14X,'MOLE FRACTIONS')
3 FORMAT(/1X,46('*'),' ACCUMULATED FLOW FIELD INFORMATION ',46('*')
1)
4 FORMAT(1X,I4,E11.3,I6,8F8.3,3X,3E11.3)

```

PT50220
PT50230
PT50240
PT50250

```

*****
      IF(KS.EQ.0) WRITE(6,1)
      IF(KS.NE.0) WRITE(6,3)
      DO 40 MT=1,3
40 DBT(MT)=0.0
      FDA=0.
      CHT=0.
      DO 50 MT=1,MSP
      CHT=CHT+CHI(MT)*RNU(MT)
50 FDA=FDA+FDN(MT)*WTM(MT)
      YCT=0.0
      ZCT=0.0
      LEVEL=1

```

PT50290

```

DO 110 M=1,NBX
N=NUMCEL(M)
IF(N.LE.0) GO TO 110
IF((ZC(M).EQ.ZCT).AND.(YC(M).EQ.YCT)) GO TO 52
ZCT=ZC(M)
YCT=YC(M)
IF(M.GE.LEV(1)) LEVEL=2
IF(M.GE.LEV(2)) LEVEL=3
WRITE(6,2) LEVEL,LKW(N),ZCT,YCT
52 XCT=XC(M)
NSAMP=0
DBA=0.
XVM=0.
YVM=0.
ZVM=0.
TMPM=0.
TRPM=0.
E=0.
F=0.
DO 100 MT=1,MSP
NSAMP=NSAMP+NS(MT,N)
XVM=XVM+XV(MT,N)*RNU(MT)*WTM(MT)*NB(MT,N)
YVM=YVM+YV(MT,N)*RNU(MT)*WTM(MT)*NB(MT,N)
ZVM=ZVM+ZV(MT,N)*RNU(MT)*WTM(MT)*NB(MT,N)
DBA=DBA+DB(MT,N)*WTM(MT)
TMPM=TMPM+WTM(MT)*RNU(MT)*TMP(MT,N)*NB(MT,N)
TRPM=TRPM+RNU(MT)*NB(MT,N)*TRP(MT,N)
E=E+WTM(MT)*RNU(MT)*NB(MT,N)
100 F=F+RNU(MT)*NB(MT,N)
DBA=DBA/FDA
IF(E.EQ.0.0) GO TO 55
XVM=XVM/E
YVM=YVM/E
ZVM=ZVM/E
VS=XVM**2+YVM**2+ZVM**2
TMPM=TMPM/E-VS
TRPM=TRPM/F
55 CONTINUE
TTM=(TMPM+TRPM)/(2.5+CHT)
TMPM=TMPM/1.5
IF(CHT.NE.-1.) TRPM=TRPM/(1.+CHT)
AMS=SQRT(VS)
IF(TTM.GT.0.) AMS=SQRT((5.+2.*CHT)*VS/(TTM*(3.5+CHT)))
CCZ=COS(ZCT/57.29578)
SCZ=SQRT(1.-CCZ**2)
RVH=ZVM*SCZ-YVM*CCZ
TVH=YVM*SCZ+ZVM*CCZ
DO 60 MT=1,MSP
DBT(MT)=RNU(MT)*NB(MT,N)
IF(F.NE.0.) DBT(MT)=DBT(MT)/F
60 CONTINUE
WRITE(6,4) M,XCT,NSAMP,DBA,AMS,XVM,RVM,TVH,TMPM,TRPM,TTM,(DBT(J),
1J=1,3)
110 CONTINUE
RETURN

```

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PT50470
PT50480

FILE: GKBEKT DECK A

PRINCETON UNIVERSITY TIME-SHARING SYSTEM

END

PT50490

/LKED.SYSLMOD DD DSN=U.GKBSPACE.ROTATION(FORTH),
/ DISP=OLD,UNIT=3350,VOL=SER=RES101,SPACE=
/LKED.SYSIN DD *
INCLUDE SYSLMOD(FORTH)
ENTRY MAIN
/ EXEC COMPRESS,DSN='U.GKBSPACE.ROTATION',RLSE=RLSE

ORIGINAL
OF PGM. ONLY

APPENDIX B

MASTER'S THESIS OF Y.P. TSAI

COLLISION INDUCED VIBRATIONAL TRANSITION
PROBABILITIES IN DIATOMIC MOLECULES

ABSTRACT

In order to improve the Monte-Carlo Direct Simulation calculations for hypersonic flow in the transition regime, we have to incorporate the effects due to vibrational non-equilibrium and potential dissociation for a diatomic gas. The state-to-state transition probabilities are desired. In this paper, we model the diatomic molecule as a harmonic oscillator which collides with another molecule collinearly. Two different methods have been developed, the first one is the semi-classical treatment and the second one is the fully quantum mechanical approach. The interaction potential between two molecules is assumed to be the Lennard-Jones 12-6 interaction law which is a small perturbation to the colliding system. Some numerical results of the state-to-state transition probabilities and comparisons are presented in Chapter 4. Discussions, which present the important aspects of this kind of problem for further study, are made in Chapter 5.

Chapter 1 Introduction

The characteristic flow in a highly rarefied gas is called "free molecular flow". In this regime the mean free path is large compared to the characteristic dimensions of an aerodynamic body in the flow; and molecules that impinge on the body, and are then reemitted from it will, in general, be far away from the body before they strike another molecule. The characteristic flow in a moderately rarefied gas is called "slip flow". The flow regime intermediate between slip and molecular flow is known as the "transition flow regime". It corresponds to densities for which the mean free path has the same general order of magnitude as the characteristic dimension of the flow field. There is a dimensionless parameter called the Knudsen number Kn , introduced to serve as a criterion for determining the relative importance of these rarefaction effects.

where
$$Kn = \frac{\lambda}{d} \sim \frac{M}{Re}$$

and λ = molecular mean free path

d = characteristic dimension of vehicle

M = Mach Number

Re = Reynolds Number

Free molecular flow is usually defined as that flow for which $Kn > 10$. Slip flow is characterized by a Knudsen number of a few per cent, $0.01 < Kn < 0.1$, and the intermediate transition regime corresponds to Knudsen number in the range $0.1 < Kn < 10$.

These values are, of course, arbitrary. Since the Knudsen number is defined as the ratio of particle mean free path to body size, it therefore increases with altitude for a fixed size body travelling through the atmosphere.

Generally, for a body of the order of a meter at altitudes 150 Kilometers above the earth's surface, the mean free path of the particles is much larger than vehicle size (here the free stream mean free path $\lambda_{\infty} \approx 40\text{m}$). Hence the particle-particle collision process in the vicinity of the vehicle need not be considered and the relation of measurement on the surface to atmospheric properties is explained in a relatively straightforward manner through free molecular theory. Below 90Km, (here the free stream mean free path $\lambda_{\infty} \approx 2.5\text{ cm}$), the Knudsen number is so small that the fluid can be treated as a continuum with limited influence of transport properties and slip-flow boundary conditions. However, between these two zones, i.e. in the lower region of the thermosphere, is what we called the "transition flow regime" in which neither of the limiting theories is applicable.

Flight at very high altitudes often involves extremely high velocities and resulting high gas temperature. At velocities which correspond to effective temperature of the order of a few thousand degrees Kelvin, the so called "real gas" effects associated with vibration, dissociation, and ionization of the gas molecules can begin to be of importance. We are interested in atmospheric entry of a

vehicle such as the space shuttle. The Mach number is generally above 20 in passing through "transition flow regime" and the flow is hypersonic there. At such high speeds, any significant number of collisions between incoming particles and those reflected from the body can produce extreme changes in the environment near the surface of the vehicle. Measurements on the surface are thus strongly affected making inference of conditions in the ambient atmosphere extremely difficult. For the purpose of data interpretation, previous calculation based on direct simulation Monte-Carlo computer technique have been developed by G.A. Bird(1). By using the Monte-Carlo method, the real gas molecules are replaced by their statistical models, and the motion of one or more of the chosen particles is traced by the computer. In the "Bird" method the real gas molecules are simulated by several thousand modeled molecules, rigid spheres in the simplest version(1). Theoretical calculations of the heat transfer and aerodynamic characteristics of a body submerged in the transition flow regime may be carried out in this way. Modifications involving more realistic interaction laws have been also carried out. Molecules, however, contain internal structure, which is important primarily for its effect on the energy content of the flowing gas. Being composed of nuclei and electrons that have motion relative to the center of mass of the molecule, the molecules can possess rotational and vibrational as well as electronic

internal states. Therefore, translational energy is not necessarily conserved in all collisions. The hypersonic flow past the sharp leading edge of a flat plate incorporating the effects of rotational non-equilibrium for a diatomic gas was studied by D.I. Pullin, J.K. Harvey, and G.K. Bienkowski(2), and subsequently applied to other blunt body problems(3,4). A different model was used to incorporate the same effect in other works such as references 5 and 6. At the present stage, we want to improve the Monte-Carlo calculations by including the effects due to the vibrational degree of freedom of diatomic molecules.

It is convenient to divide vibrational energy exchange into two cases: the V-V process, in which the total vibrational quantum number of the system is unchanged, and the V-T process in which energy is exchanged between translation and vibration without conserving the vibrational quantum. For harmonic oscillators in the V-V process, the amount of vibrational energy lost by one molecule is gained by the other and no vibrational-translational energy transfer occurs. Considerable interest has been shown in the details of inelastic molecular collisions. The treatment of scattering between particles with internal structure is capable of producing differential and total cross section for state-to-state transitions. Our purpose is to calculate scattering cross sections for different transitions with known initial and final states as a function of collision energy or initial relative velocity. We emphasize here that

we need state-to-state transition cross sections rather than some overall "rates of transition" (rate coefficients) which are averaged over an equilibrium velocity distribution for the relative translational motion (e.g. Maxwellian distribution). The reasons are twofold:

1. We are dealing with a hypersonic flow system in a highly non-equilibrium state with a non-Maxwellian distribution of relative velocities of collision. The relative contributions of different energy molecules to the overall rates may therefore be drastically different than in the equilibrium state. This effect is accentuated by the steep rise of cross-sections with energy coupled to the generally decreasing magnitude of the distribution function with energy. This unifies that small changes in the fraction of molecules with high energies due to the non-equilibrium aspect of the flow can have extreme effects on inelastic processes without corresponding effects on mean properties such as density or fluid momentum.
2. The Direct Simulation Monte Carlo Code (DSMCC) consists of tracing a set of "test" molecules through a designated volume surrounding the body. The velocities and internal states of the simulated molecules are altered on the basis of collisions computed (as determined by local

collision probabilities) at fixed time steps. The positions are then advanced to new values on the basis of motion through the time step increments. This detailed computation of individual molecular collisions requires, in principle, state to state transition probabilities in order to incorporate the inelastic energy exchange into the computation of molecular velocities and internal states after collision. While the results we desire must ultimately be consistent with overall measured rates the level of detail necessary within the program is well beyond the level of availability of experimental data.

Early theoretical studies of vibrational, rotational, and translational energy transfer in collisions were based on approximate analytical solutions to the quantum mechanical and classical equations of motion. The method of Zener (7), later to become known as the distorted wave method, and the Born (8) approximation are leading examples of approximate solutions to quantum mechanical collision problems based on first order perturbation theory. A more detailed literature review on previous work in the field of vibrational collisions is given in the next Chapter.

Instead of doing an approximate treatment of a three-dimensional realistic system, in this work, we do an exact numerical treatment of a simpler model one-dimensional system

which has some important features in common with the real one. A calculation of transition probabilities for vibrational-vibrational-translational energy transfer in a collision of two diatomic molecules is to be presented. This simple collision model is approximate, utilizing a collinear collision of harmonic oscillators with an exponential repulsion between center atoms and no chemical reaction between the molecules. It may be argued that the configuration allowing the most efficient transfer of energy between translation and vibration is that in which the atoms are collinear. Collinear or head-on collisions make the most significant contribution to the transition probability. The averaged probability is equal to the probability of excitation in a head-on collision times a "steric factor" smaller than unity which takes account of unfavorable trajectories. For homonuclear molecules it is usually taken as $\frac{1}{3}$ (the average of $\cos^2\theta$ taken over a sphere). A more detailed theory for the steric factor of linear molecules has been propounded by Herzfeld (9). It is our belief that an accurate treatment of a collinear model is of more worth than an approximate result for the three-dimensional problem. The latter approach frequently contains errors which are difficult to estimate. The general magnitudes and trends of the transition probability obtained by this restricted treatment can show us qualitatively, or semi-quantitatively, some characteristic features of the problem.

Two different methods for calculating transition

probabilities are discussed in this paper. One is a semi-classical approach, and the other one is a fully quantum mechanical treatment. In the semi-classical calculation,, it is assumed that the vibrational amplitude of the harmonic oscillators are small and so the molecular oscillations do not greatly affect the external classical collision trajectory. The trajectory can be calculated from the classical equation of motion. The classical trajectory is assumed to define a time-dependent perturbation potential for the colliding system and quantum theory is used to derive the transition probability. Essentially, the Schrodinger equation is solved subject to certain initial conditions according to time-dependent perturbation theory. A detailed discussion about this theory can be found in the book Quantum Mechanic by Schiff (10).

In the quantum mechanical calculation of transition probability, the wavefunction of the whole system is expanded in terms of the complete set of eigenfunctions of vibrational states of the diatomic molecule. With the aid of the orthonormality property of these eigenfunctions, a set of coupled second order differential equations is obtained for the translational wavefunctions. The transition probability is given by the solution to this set of equations in the asymptotic region, subject to appropriate boundary condition. Several different numerical methods for solving the set of coupled equations have been developed by Diestler and Mckoy (11).

Riley and Kuppermann (12), and Gutshick et al (13). Solving a problem of quantum scattering between two diatomic molecules is then reduced to the task of finding a good numerical scheme for integrating a system of coupled differential equations accurately. In this work, we use IBM IMSL ROUTINE DGEAR to solve for the scattered wavefunction in the asymptotic region and then calculate the transition probability.

In Chapter 2, a brief review of previous semiclassical and quantum mechanical methods on vibrational collisions will be given so that we can identify the new points in the current work in Chapter 3. Chapter 3 is devoted to a discussion of the general theory. We derive some equations and expressions there which make a numerical algorithm feasible. In Chapter 4, theoretical results of state-to-state transition probability specifically for N_2-N_2 collisions at different relative velocities are presented. Comparison is made with published results. Finally, we discuss some important problems related to inelastic molecular collisions which deserve further study because they make extension to a more realistic treatment of molecular scattering possible.

Chapter 2 Literature Review

In 1931, Oldenberg (14) discussed molecular collision processes qualitatively to show the persistence of the rotational and vibrational motion. Zener (7) was the first to give a detailed mathematical treatment for collisions in which molecular vibrations are excited or de-excited. He restricted himself to collinear collisions between a diatomic molecule and an atom. His theory was based on the distorted wave method which includes direct transition from the initial state to the final state and assumes that the probability of transition is small. It is a perturbation method and cannot treat strongly coupled system. Takayanagi (15) then extended Zener's one-dimensional treatments to three-dimensional collisions. In order to save computational labor, the modified wave number approximation was introduced. Meanwhile, Schwartz, Slawsky and Herzfeld (16) gave a mathematical formulation for the vibrational transitions based on the distorted wave approximation due to Jackson and Mott (17), simplified by the modified wave number approximation due to Takayanagi, in diatom-diatom collisions. Their formulation is referred to as the SSH theory now and is a quantum mechanical result. For purpose of future reference, we describe SSH theory in more detail.

Consider the head-on collision between two diatomic molecules AB and CD (assume harmonic oscillators). For exponential intermolecular interaction between nearest

atoms B and C, one has,

$$\begin{aligned} V &= V_0 \exp(-\alpha r_{BC}) \\ &= V_0 \exp[-\alpha(R - \lambda_1 r_1 - \lambda_2 r_2)] \end{aligned} \quad (2-1)$$

where R is the distance between centers of mass of two molecules, r_1 and r_2 are vibrational coordinates, $\frac{1}{\alpha}$ is the range of the potential and V_0 is a constant.

$$\lambda_1 = \frac{m_A}{m_A + m_B}, \quad \lambda_2 = \frac{m_D}{m_C + m_D}$$

m_i is the mass of the i -th atom. Solving the Schrodinger equation, the transition probability is given in the closed analytic form:

$$P(n_1 n_2 \rightarrow n'_1 n'_2) = \frac{\pi^2}{4} |v_1(n'_1 n_1)|^2 |v_2(n'_2 n_2)|^2 \left\{ \frac{\Delta q^2}{\cosh \pi q' - \cosh \pi q} \right\}^2 \cdot \sinh \pi q' \cdot \sinh \pi q \quad (2-2)$$

where

$$v_1(n'_1 n_1) = \int_{-\infty}^{+\infty} z_1(n'_1, r) \exp[\alpha \lambda_1 (r - r_{e1})] z_1(n_1, r) dr \quad (2-3)$$

$$v_2(n'_2 n_2) = \int_{-\infty}^{+\infty} z_2(n'_2, r) \exp[\alpha \lambda_2 (r - r_{e2})] z_2(n_2, r) dr \quad (2-4)$$

$z_i(n_j, r)$ is the vibrational wave function for harmonic oscillator " i " in quantum state n_j ; r_{e1} and r_{e2} are equilibrium separations for AB and CD respectively, and

$$q = q_{n_1 n_2} = \frac{2kn_1 n_2}{\alpha} \quad (2-5)$$

$$q' = q_{n'_1 n'_2} = \frac{2kn'_1 n'_2}{\alpha} \quad (2-6)$$

$$\Delta q^2 = q'^2 - q^2 = - \left(\frac{8\mu}{h^2 \alpha^2} \right) [\epsilon_1(n'_1) + \epsilon_2(n'_2) - \epsilon_1(n_1) - \epsilon_2(n_2)] \quad (2-7)$$

The wave number of relative motion before and after collision are $k_{n_1 n_2}$ and $k_{n'_1 n'_2}$, which satisfy the energy conservation law:

$$\epsilon_1(n_1) + \epsilon_2(n_2) + \frac{h^2 k_{n_1 n_2}^2}{2\mu} = \epsilon_1(n'_1) + \epsilon_2(n'_2) + \frac{h^2 k_{n'_1 n'_2}^2}{2\mu} \quad (2-8)$$

μ is the reduced mass of the whole system; $\epsilon_i(n)$ is the energy of molecule "i" in the n-th vibration state. For the special case, V-V transition, in which $n_1 + n_2 = n'_1 + n'_2$ and $\Delta q = 0$. Applying L'Hospital's rule, we get:

$$\lim_{q' \rightarrow q} \frac{q'^2 - q^2}{\cosh \pi q' - \cosh \pi q} = \frac{2q}{\pi \sinh \pi q}$$

Then the transition probability for V-V process becomes

$$p(n_1 n_2 \rightarrow n'_1 n'_2) = |V_1(n'_1 n_1)|^2 |V_2(n'_2 n_2)|^2 q^2 \quad (2-9)$$

SSH theory has been most widely used for quantitative comparison with experimental measurements of vibrational relaxation, but the coupling between rotation has been ignored.

Zelechow, Rapp and Sharp (ZRS) (18) have developed a semi-classical method for calculating transition probabilities for V-V and V-T energy transfer in a collision of two diatomic molecules. Their basic

assumption are:

- (1) The perturbation potential is linearized in the oscillator coordinates.
- (2) The collision velocity is not too high (e.g. the upper limit for N_2-N_2 collision is 10Km/sec) and the collision induced time-varying force constant $k'(t)$ is small compared to k , the characteristic force constant of the molecule.

Under these two conditions, Kerner (19) method can be applied to solve the Schrodinger equation and closed form analytical results are obtained. However, this approach restricts itself to the transitions of processes of symmetric type only. The general formula is:



where n, m are vibration quantum numbers before collision and n', m' are that after collision. The collision is symmetric in the sense that the two B atoms are in the center.

The development of the high speed electronic computer has made it possible to solve the collision problem by direct numerical techniques. T.E. Sharp and D. Rapp (20) have calculated the vibrational transition probabilities for collisions between a diatomic molecule and an atom. In their semi-classical treatment, an N-state approximation method is used, in which the total wave function is expanded in terms of N eigenfunctions of stationary states of the system including the initial, final and all

energetically intervening states. A Runge-Kutta single-step integration method is employed in the computation program. Generally, the value of N needed in expanding the total wave function increases with collision velocity. An "exact" solution for any transition probability $P_{j \rightarrow k}$ is reached when the addition of more states to the computation results in no significant change in $P_{j \rightarrow k}$. We extend this method to collisions between two diatomic molecules (Chapter 3-A).

In quantum mechanical treatment of collisions between two diatomic molecules AB and CD, taking B and C as the inner atoms of the system, the total wave function is expanded in terms of normalized vibrational wave functions $Z_{AB}(n_1, r_1)$ and $Z_{CD}(n_2, r_2)$. That is:

$$\psi = \sum_{n_1} \sum_{n_2} f_{n_1 n_2}(R) Z_{AB}(n_1, r_1) Z_{CD}(n_2, r_2)$$

Inevitably, we have to solve a system of coupled differential equations, which are equivalent to the Schrodinger equation, of the following form (detailed discussions will be given in Chapter 3-C).

$$\frac{\hbar^2}{2\mu} \frac{d^2}{dR^2} + k_{n_1' n_2'}^2 f_{n_1 n_2}(R) = \sum_{n_1} \sum_{n_2} \langle n_1' n_2' | V | n_1 n_2 \rangle f_{n_1 n_2}(R) \quad (2-10)$$

where $\langle n_1' n_2' | V | n_1 n_2 \rangle$ is the matrix element. In principle if we can obtain the solution to equation (2-10) with the asymptotic form

$$\begin{aligned}
 f_{n_1 n_2}'(R) &\rightarrow 0 & R \rightarrow -\infty \\
 f_{n_1 n_2}'(R) &\rightarrow \delta_{n_1' n_1} \delta_{n_2' n_2} \exp(-ik_{n_1' n_2'} R) \\
 &+ A_{n_1' n_2'; n_1 n_2} \exp(+ik_{n_1' n_2'} R) & R \rightarrow +\infty
 \end{aligned}
 \tag{2-11}$$

the probability per collision for transition $(n_1, n_2) \longrightarrow (n_1', n_2')$ will be given by:

$$P(n_1 n_2 \rightarrow n_1' n_2') = \frac{k_{n_1' n_2'}}{k_{n_1 n_2}} |A_{n_1' n_2'; n_1 n_2}|^2
 \tag{2-12}$$

A number of numerical methods (7-9) have been proposed for solving the systems of equations (2-10). However, due to the rapidly oscillating wavelike solutions to the Schrodinger equation, the numerical technique is not straightforward. Riley (12) developed the initial-value technique with periodic "reorthogonalization". Gadschick et al. (13), on the other hand introduced a technique of integration using Dirichlet boundary condition and single one-step Euler integration. A new method for constructing wave function for bound states and scattering has been proposed by Roy G. Gordon (21), perhaps this procedure can save much computer time. Our quantum mechanical treatment of this molecule scattering problem is similar to the method due to Riley and Kuppermann (12). It is relatively simple and straightforward, but in our procedures, the virtual states (energetically inaccessible) are not included in the total wavefunction expansion.

Chapter 3 Theory

A. General Formalism

The collision model is shown in Fig. 3-1.

Fig. 3-1 Collision Coordinates

This figure is the collinear collision configuration between two diatoms AB and CD. Assuming that CD is the target, and AB is the incident projectile from right. The laboratory coordinates of A, B, C and D are x_A , x_B , x_C and x_D ; their masses are m_A , m_B , m_C and m_D respectively. Let V_{AB} and V_{CD} be the binding potential of molecules AB and CD. The short range interaction is assumed to be a sum of interatomic interactions,

$$V_{INT} = V_{AC}(x_A - x_C) + V_{AD}(x_A - x_D) + V_{BD}(x_B - x_D) + V_{BC}(x_B - x_C)$$

The interatomic potentials are exponentially decreasing functions, so that for the collinear configuration under consideration, only the term $V_{BC}(x_B - x_C)$ is important and thus

$$V_{INT} \approx V_{BC}(x_B - x_C)$$

The Schrodinger equation for the system is:

$$\left\{ -\frac{h^2}{2M_A} \frac{\partial^2}{\partial x_A^2} - \frac{h^2}{2M_B} \frac{\partial^2}{\partial x_B^2} - \frac{h^2}{2M_C} \frac{\partial^2}{\partial x_C^2} - \frac{h^2}{2M_D} \frac{\partial^2}{\partial x_D^2} + V_{AB}(x_A - x_B) + V_{CD}(x_C - x_D) \right. \\ \left. + V_{INT}(x_B - x_C) \right\} \psi(x_A, x_B, x_C, x_D) = E_{TOT} \psi(x_A, x_B, x_C, x_D) \quad (3-1)$$

We designate the distance between the centers of mass of two molecules as R . In molecule AB the distance between the atoms is x ; in CD it is y . x and y are internal coordinates. i.e.

$$R = \frac{m_A x_A + m_B x_B}{m_A + m_B} - \frac{m_C x_C + m_D x_D}{m_C + m_D}$$

$$x = x_A - x_B$$

$$y = x_C - x_D$$

$$\text{Let } R_{CM} = \frac{m_A x_A + m_B x_B + m_C x_C + m_D x_D}{m_A + m_B + m_C + m_D} \text{ which is the coordinate of}$$

center of mass of the whole system. In terms of the new coordinates (x, y, R, R_{CM}) the Schrodinger equation becomes:

$$\left\{ -\frac{h^2}{2\mu_{AB}} \frac{\partial^2}{\partial x^2} - \frac{h^2}{2\mu_{CD}} \frac{\partial^2}{\partial y^2} - \frac{h^2}{2\mu} \frac{\partial^2}{\partial R^2} - \frac{h^2}{2M} \frac{\partial^2}{\partial R_{CM}^2} + V_{AB}(x) + V_{CD}(y) \right. \\ \left. + V_{INT}(R - \gamma_{AB}x - \gamma_{CD}y) \right\} \psi(x, y, R, R_{CM}) = E_{TOT} \psi(x, y, R, R_{CM}) \quad (3-2)$$

where $M = m_A + m_B + m_C + m_D = \text{total mass of the system.}$

$$\mu = \frac{(m_A + m_B)(m_C + m_D)}{M} = \text{reduced mass of the system}$$

$$\mu_{AB} = \frac{m_A m_B}{m_A + m_B} \quad = \text{reduced mass of the molecule AB}$$

$$\mu_{CD} = \frac{m_C m_D}{m_C + m_D} \quad = \text{reduced mass of the molecule CD}$$

$$\gamma_{AB} = \frac{m_A}{m_A + m_B}, \quad \gamma_{CD} = \frac{m_D}{m_C + m_D}$$

Since there is no external force applied to the system, the center of mass of the system moves like a free particle and its motion can be described by a plane wave $e^{iK_{CM}R_{CM}}$, and the energy of the center of mass, $T_{CM} = \frac{h^2 K_{CM}^2}{2M}$, is also a constant of motion. This does not affect the energy transfer and need not be considered further. We can remove the R_{CM} -dependent part of the wave function (x, y, R, R_{CM}) by separation of variables.

Let

$$\Psi(x, y, R, R_{CM}) = \psi(x, y, R) e^{iK_{CM}R_{CM}}, \quad E_{TOT} = E + T_{CM}$$

Substituting these into equation (3-2), we arrive at a Schrodinger equation concerning the internal coordinates x, y and the relative motion R of the two colliding molecules as follows:

$$\left\{ -\frac{h^2}{2\mu_{AB}} \frac{\partial^2}{\partial x^2} - \frac{h^2}{2\mu_{CD}} \frac{\partial^2}{\partial y^2} - \frac{h^2}{2\mu} \frac{\partial^2}{\partial R^2} + V_{AB}(x) + V_{CD}(y) + V_{INT}(R - \gamma_{AB}x - \gamma_{CD}y) \right\}$$

$$\psi(x, y, R) = E\psi(x, y, R) \quad (3-3)$$

Define

$$X = x - x_{eq}$$

$$Y = y - y_{eq}$$

where x_{eq} and y_{eq} are equilibrium separations of AB and CD. Then X and Y are displacements from equilibrium of each oscillator. Now we introduce harmonic bonds (intramolecular potential) into AB and CD with force constants k_{AB} and k_{CD} , hence

$$V_{AB} = 1/2 k_{AB} X^2,$$

$$V_{CD} = 1/2 k_{CD} Y^2,$$

A conventional representation of the intermolecular potential energy curve is given by the Lennard-Jones 12-6 equation. Since the elementary models for energy transfer are based on exponential potential, the exponential function $V_{INT}(x_B - x_C) \sim e^{-(x_B - x_C)/L} - \epsilon$ must be fitted to the Lennard-Jones potential (Appendix 1), where L is a parameter characterizing the range of the interaction. Landau and Teller (22) assumed that only the short range repulsive part of the intermolecular potential is steep enough to influence energy transfer, so that the long-range attractive potential ϵ can be neglected. The molecular interaction is then assumed to be an exponential repulsion between atoms C and B.

Let

v = initial relative velocity

$E = 1/2 \mu v_0^2$ = initial relative kinetic energy

$$\bar{R} = R - R_T$$

where R_T is the distance at the classical turning point. The potential energy in equation (3-3) may be expressed as

(16):

$$V_{INT}(X, Y, \bar{R}) = E_0 \exp \left[\frac{-1}{L} (\bar{R} - \gamma_{AB} X - \gamma_{CD} Y) \right] \quad (3-4)$$

Equation (3-3) can be written in the form:

$$\left\{ \frac{-h^2}{2\mu_{AB}} \frac{\partial^2}{\partial X^2} - \frac{h^2}{2\mu_{CD}} \frac{\partial^2}{\partial Y^2} - \frac{h^2}{2\mu} \frac{\partial^2}{\partial \bar{R}^2} + \frac{1}{2} k_{AB} X^2 + \frac{1}{2} k_{CD} Y^2 \right. \\ \left. + E_0 \exp \left[\frac{-1}{L} (\bar{R} - \gamma_{AB} X - \gamma_{CD} Y) \right] \right\} \psi(X, Y, \bar{R}) = E \psi(X, Y, \bar{R}) \quad (3-5A)$$

or

$$\left\{ - \frac{h^2}{2\mu_{AB}} \frac{\partial^2}{\partial X^2} - \frac{h^2}{2\mu_{CD}} \frac{\partial^2}{\partial Y^2} + \frac{1}{2} k_{AB} X^2 + \frac{1}{2} k_{CD} Y^2 + E_0 \exp \left(\frac{-1}{L} (\bar{R} - \gamma_{AB} X - \gamma_{CD} Y) \right) \right\}$$

$$\psi(X, Y, \bar{R}) = i\hbar \frac{\partial}{\partial t} \psi(X, Y, \bar{R}) \quad (3-5B)$$

We shall solve the time-independent Schrodinger equation (3-5A) by purely quantum mechanical method and the time-dependent Schrodinger equation (3-5B) by semi-classical method.

B. Semi-classical Calculation

Since the deBroglie wavelength of the relative motion of two molecules is usually very small compared with atomic dimensions (e.g. for N_2 - N_2 collisions at velocity 5Km/sec, the deBroglie wavelength is of the order 10^{-15} cm, however the dimension of N_2 molecules is of the order of a few Å), it is a fairly good approximation to use the classical trajectory for relative motion. The classical equations of motion are:

$$\mu \frac{d^2 \bar{R}}{dt^2} = - \frac{\partial}{\partial \bar{R}} V_{INT}(X, Y, \bar{R}) \quad (3-6A)$$

$$\mu_{AB} \frac{d^2 X}{dt^2} = - k_{AB} X - \frac{\partial}{\partial X} V_{INT}(X, Y, \bar{R}) \quad (3-6B)$$

$$\mu_{CD} \frac{d^2 Y}{dt^2} = - k_{CD} Y - \frac{\partial}{\partial Y} V_{INT}(X, Y, \bar{R}) \quad (3-6C)$$

Generally, the incident energy is much larger than the change in the vibrational energy, or in other words, due to small transition probabilities, only a small fraction of the translational energy is transferred to vibrational energy. We may then assume that during the collision, the vibrational amplitudes of the oscillators are not driven to large values, that means

$$\begin{aligned} X &\ll L, \\ Y &\ll L. \end{aligned} \quad (3-7)$$

In a series calculation given by Wolfberg and Kelley (23), we can see that conditions (3-7) are justifiable. Wolfberg and Kelley have calculated the energy transfer for collisions involving two harmonic oscillators via an exponential collision with $L=0.22\text{\AA}$. Other parameters and data are: $m_A=m_B=m_C=m_D=12\text{a.m.u.}$, angular frequency $=2.3 \times 10^{14} \text{ sec}^{-1}$ (cf: for N_2-N_2 collisions $m_A=m_B=m_C=m_D=14 \text{ a.m.u.}$, $=4.45 \times 10^{14} \text{ sec}^{-1}$), the initial energy $E_0 = 5.078 \text{ ev}$ (corresponds to $v_0 = 9 \text{ Km/sec}$), then the vibrational energy transferred to each diatomic molecule is $\Delta E_{CD} = 1.78 \times 10^{-3} \text{ ev}$. Obviously, both ΔE_{AB} and ΔE_{CD} are much less than E_0 .

The vibrational amplitude never exceeds 0.007 \AA , and since $L \approx 0.22 \text{ \AA}$ the conditions (3-7) are fairly well satisfied. At low velocities (still high enough so that the deBroglie wavelength is much less than the atomic dimension), the energy transfer becomes smaller because the oscillator can readjust adiabatically to the perturbation caused by the incident particle. Conditions (3-7) are satisfied even better. Under conditions (3-7), the molecular oscillations do not greatly affect the external classical collision trajectory. Therefore, one can neglect the motion in X and Y in treating the motion in R . Equations (3-6A), (3-6B), and (3-6C) may then be replaced by:

$$\mu \frac{d^2 \bar{R}}{dt^2} = - \frac{\partial}{\partial \bar{R}} V_{\text{INT}}(\bar{R}, X=0, Y=0) \quad (3-8)$$

Solving equations (3-8) with V_{INT} given by (3-4), one finds that the trajectory $R(t)$ satisfies the relation

$$\exp\left(-\frac{\bar{R}(t)}{L}\right) = \text{sech}^2 \left(\frac{u_0 t}{2L} \right) \quad (3-9)$$

$R(t)$ is then inserted into the interaction potential function $V_{\text{INT}}(X, Y, \bar{R})$ to obtain $V_{\text{INT}}(X, Y, t)$. In this semi-classical treatment, $V_{\text{INT}}(X, Y, t)$ is used as a transition inducing perturbation acting upon a quantum mechanical harmonic oscillator. Finally, from equation (3-5B) we get the time-dependent Schrodinger equation to be

solved numerically by first order time-dependent perturbation theory.

$$\left\{ -\frac{\hbar^2}{2\mu_{AB}} \frac{\partial^2}{\partial X^2} + \frac{1}{2} k_{AB} X^2 - \frac{\hbar^2}{2\mu_{CD}} \frac{\partial^2}{\partial Y^2} + \frac{1}{2} k_{CD} Y^2 + E_0 \exp\left(\frac{\gamma_{AB} X + \gamma_{CD} Y}{L}\right) \right. \quad (3-10)$$

$$\left. \operatorname{sech}^2\left(\frac{u_C t}{2L}\right) \right\} \quad \psi(X, Y, t) = i\hbar \frac{\partial}{\partial t} \psi(X, Y, t)$$

There is one more comment about the assumptions of this semi-classical model. In principle, if the deBroglie wavelength associated with the relative motion is much less than the atomic dimensions, it is a good approximation to use classical mechanics for solving the relative trajectory. As just mentioned before, if the criteria (3-7) are fulfilled, in other words, each of the oscillation amplitudes of the two colliding diatomic molecules is very small compared to the characteristic range of the interaction potential, a further simplification can be made in solving for $\bar{R}(t)$, the relative trajectory. These procedures correspond to the "approximate" classical method, because the trajectory $\bar{R}(t)$ is determined with X and Y set equal to zero. As a result, these calculations do not include the conservation of energy, and E is assumed to remain as the energy in coordinate \bar{R} , regardless of how much excitation occurs in the oscillators. For more accurate semi-classical calculations, one must solve equations (3-6A), (3-6B), and (3-6C) for "exact" classical trajectory $\bar{R}(t)$ to be used as a time-dependent perturbation. Prior to the work of Wolfsberg and

Kelley (23), it has been thought that the criteria (3-7) are automatically satisfied for all low-velocity collisions. Actually, this is not the case. Wolfsberg and Kelley proved that the approximate classical method should be limited to collisions between a light particle and a heavy oscillator.

For our case of N_2-N_2 collision, Wolfsberg and Kelley's requirement is satisfied. The energy transferred to each diatomic species is very small compared to the kinetic energy of relative motion. Both molecules are negligibly distorted during collision approach and the molecular oscillations never deviate substantially from their equilibrium configurations. The semi-classical approach is a good approximation even within the high velocity range in which we are interested. The reason is that as v_0 increases, the deBroglie wavelength characterizing the relative motion in R is small compared to the distance over which the interaction potential varies significantly in R , i.e. $L.k_{n1}n_2 \gg 1$. This implies that both q and q' defined by equations (2-5) and (2-6) are much greater than unity. In this limit, $q \gg 1$ and $q' \gg 1$, the quantum mechanical results of equation (2-2) are reduced to the classical results of Landau and Teller (22).

Let ω_{AB} and ω_{CD} be the angular frequencies for the AB and CD molecules, their isolated Hamiltonians are $H^{(AB)}(X)$ and $H^{(CD)}(Y)$ respectively, then

$$\omega_{AB} = \frac{k_{AB}}{\mu_{AB}}$$

$$\omega_{CD} = \frac{k_{CD}}{\mu_{CD}}$$

$$H^{(AB)}(X) = -\frac{\hbar^2}{2\mu_{AB}} \frac{\partial^2}{\partial X^2} + \frac{1}{2} \mu_{AB} \omega_{AB}^2 X^2$$

$$H^{(CD)}(Y) = -\frac{\hbar^2}{2\mu_{CD}} \frac{\partial^2}{\partial Y^2} + \frac{1}{2} \mu_{CD} \omega_{CD}^2 Y^2$$

Suppose the individual eigenfunctions of $H^{(AB)}(X)$ and $H^{(CD)}(Y)$ are $\phi_n^{(AB)}(X)$ and $\phi_j^{(CD)}(Y)$, we have:

$$H^{(AB)}(X) \phi_n^{(AB)}(X) = (n + \frac{1}{2}) \hbar \omega_{AB} \phi_n^{(AB)}(X) \quad n = 0, 1, 2, \dots$$

$$H^{(CD)}(Y) \phi_j^{(CD)}(Y) = (j + \frac{1}{2}) \hbar \omega_{CD} \phi_j^{(CD)}(Y) \quad j = 0, 1, 2, \dots$$

Set

$$\phi_{nj}(X, Y) = \phi_n^{(AB)}(X) \phi_j^{(CD)}(Y) \quad (3-11)$$

$$W_{ij} = (n + \frac{1}{2}) \hbar \omega_{AB} + (j + \frac{1}{2}) \hbar \omega_{CD} \quad (3-12)$$

$$H_0(X, Y) = H^{(AB)}(X) + H^{(CD)}(Y) \quad (3-13)$$

$H_0(X, Y)$ is the unperturbed Hamiltonian of the system. It is obvious that $H_0(X, Y) \phi_{nj}(X, Y) = W_{nj} \phi_{nj}(X, Y)$ (3-14)
 W_{ij} are eigenvalues of the unperturbed Hamiltonian $H_0(X, Y)$. The solution to the equation (3-10), i.e. the total wavefunction for the system of two oscillators $\psi(X, Y, t)$ can be expanded in terms of the individual harmonic oscillator wavefunctions $\phi_n^{(AB)}(X)$ and $\phi_j^{(CD)}(Y)$.

$$\psi(X,Y,t) = \sum_n \sum_j a_{nj}(t) \phi_n^{(AB)}(X) \phi_j^{(CD)}(Y) e^{-i(n+\frac{1}{2})\omega_{AB}t} e^{-i(j+\frac{1}{2})\omega_{CD}t} \quad (3-15)$$

where the expansion coefficients a_{nj} depend on time.

If the oscillators AB and CD are initially in state N and J respectively, then the initial conditions are:

$$n = 0, 1, 2, \dots$$

$$a_{nj}(-\infty) = \delta_{nN} \delta_{jJ}$$

(Appendix 2)

$$j = 0, 1, 2, \dots$$

The probability of the system ending up with AB in state Q and CD in state K is

$$P_{NJ \rightarrow QK} = |a_{QK}(+\infty)|^2$$

We are now in the position to solve for the expansion coefficients $a(t)$. Substituting equation (3-15) into Schrodinger equation (3-10), and using the relations (3-11) \rightarrow (3-14), we find:

$$\begin{aligned} (ih) \sum_{nj} \frac{d}{dt} a_{nj}(t) \phi_n^{(AB)}(X) \phi_j^{(CD)}(Y) e^{-i(n+\frac{1}{2})\omega_{AB}t} e^{-i(j+\frac{1}{2})\omega_{CD}t} \\ = E_0 \exp\left(\frac{\gamma_{AB}X + \gamma_{CD}Y}{L}\right) \text{sech}^2\left(\frac{u_0 t}{2L}\right) \sum_{nj} a_{nj}(t) \phi_n^{(AB)}(X) \phi_j^{(CD)}(Y) \\ e^{-i(n+\frac{1}{2})\omega_{AB}t} e^{-i(j+\frac{1}{2})\omega_{CD}t} \end{aligned} \quad (3-16)$$

Equation (3-16) is multiplied by $\phi_{n'}^{(AB)*}(X) \phi_{j'}^{(CD)*}(Y)$ on both side, where $\phi_{n'}^{(AB)*}$ is the complex conjugate function of $\phi_{n'}^{(AB)}$ and $\phi_{j'}^{(CD)*}(Y)$ is the complex conjugate of $\phi_{j'}^{(CD)}(Y)$, and then integrated. Equation (3-16) then becomes:

$$\begin{aligned}
& (i\hbar) \sum_{nj} \frac{d}{dt} a_{nj}(t) \delta_{nn'} \delta_{jj'} e^{-i(n+\frac{1}{2})\omega_{AB}t} e^{-i(j+\frac{1}{2})\omega_{CD}t} \\
& = E_0 \operatorname{sech}^2\left(\frac{u_0 t}{2L}\right) \sum_{nj} a_{nj}(t) \left[\int_{-\infty}^{+\infty} \phi_{n'}^{(AB)*}(X) e^{\gamma_{AB}X} \phi_n^{(AB)}(X) dX \right] \quad (3-17)
\end{aligned}$$

$$\left[\int_{-\infty}^{+\infty} \phi_j^{(CD)*}(Y) e^{\gamma_{CD}Y} \phi_j^{(CD)}(Y) dY \right] e^{-i(n+\frac{1}{2})\omega_{AB}t} e^{-i(j+\frac{1}{2})\omega_{CD}t}$$

where use has been made of the orthonormal property for the set of harmonic oscillator energy eigenfunctions $\phi_n^{(AB)}(X)$ and $\phi_j^{(CD)}(Y)$.

$$\int_{-\infty}^{+\infty} \phi_{n'}^{(AB)*}(X) \phi_n^{(AB)}(X) dX = \delta_{nn'}$$

$$\int_{-\infty}^{+\infty} \phi_{j'}^{(CD)*}(Y) \phi_j^{(CD)}(Y) dY = \delta_{jj'}$$

Let
$$U_{n'n} = \int_{-\infty}^{+\infty} \phi_{n'}^{(AB)*}(X) e^{\gamma_{AB}X} \phi_n^{(AB)}(X) dX$$

$$V_{j'j} = \int_{-\infty}^{+\infty} \phi_{j'}^{(CD)*}(Y) e^{\gamma_{CD}Y} \phi_j^{(CD)}(Y) dY$$

$U_{n'n}$ and $V_{j'j}$ are the coupling terms between states, called matrix elements. Equations (3-17) can be written now as:

$$\begin{aligned}
\frac{d}{dt} a_{n',j'}(t) &= \left(\frac{E_0}{i\hbar}\right) \operatorname{sech}^2\left(\frac{u_0 t}{2L}\right) \sum_{nj} a_{nj}(t) U_{n'n} V_{j'j} e^{-i(n+\frac{1}{2})\omega_{AB}t} \\
&\quad e^{-i(j+\frac{1}{2})\omega_{CD}t} \quad (3-18)
\end{aligned}$$

Equation (3-18) is a set of coupled first order differential equations subject to certain initial conditions. Actually, this set of differential equations is equivalent to the Schrodinger equation (3-10). Next, let

$$a_{nj}(t) = A_{nj}(t) + iB_{nj}(t) \quad (3-19)$$

where $A_{nj}(t)$ is the real part of $a_{nj}(t)$, and $B_{nj}(t)$ is the imaginary part of $a_{nj}(t)$. Both $A_{nj}(t)$ and $B_{nj}(t)$ are real functions of time t . We also know that $U_{m'n'}$, V_{jj} are real, and

$$e^{i(n+\frac{1}{2})\omega_{AB}t} e^{i(j+\frac{1}{2})\omega_{CD}t} = \cos [(n+\frac{1}{2})\omega_{AB}t + (j+\frac{1}{2})\omega_{CD}t] + i \sin [(n+\frac{1}{2})\omega_{AB}t + (j+\frac{1}{2})\omega_{CD}t] \quad (3-20)$$

Inserting equations (3-19) and (3-20) into equation (3-18), with some algebra, the real part and imaginary part on each side of the equation must be equal separately. So we obtain a set of coupled differential equations for A's and B's:

$$\begin{aligned} \frac{d}{dt} A_{n',j'}(t) = & \frac{E_0}{h} \operatorname{sech}^2 \left(\frac{u_0 t}{2L} \right) \sum_{nj} \{ A_{nj}(t) \sin [(n'-n)\omega_{AB}t + (j'-j)\omega_{CD}t] \\ & + B_{nj}(t) \cos [(n'-n)\omega_{AB}t + (j'-j)\omega_{CD}t] \} U_{n'n} V_{j'j} \end{aligned} \quad (3-21)$$

$$\begin{aligned} \frac{d}{dt} B_{n',j'}(t) = & \frac{E_0}{h} \operatorname{sech}^2 \left(\frac{v_0 t}{2L} \right) \sum_{nj} \{ B_{nj}(t) \sin [(n'-n)\omega_{AB}t + (j'-j)\omega_{CD}t] \\ & - A_{nj}(t) \cos [(n'-n)\omega_{AB}t + (j'-j)\omega_{CD}t] \} U_{n'n} V_{n'j} \end{aligned}$$

$A_{rs}(+\infty)$ and $B_{rs}(+\infty)$ are to be found, hence

$$|a_{rs}(+\infty)|^2 = |A_{rs}(+\infty)|^2 + |B_{rs}(+\infty)|^2,$$

the desired transition probability.

In Chapter 2, we mentioned that in the work of ZRS (18), they treated collinear collisions between two

diatomic molecules with symmetric configuration only. Therefore, γ_{AB} is equal to γ_{CD} i.e.:

$$\gamma_{AB} = \gamma_{CD} = \gamma$$

The intermolecular potential function in equation (3-10), $E_0 \text{sech}^2 \left(\frac{v_0 t}{2L} \right) \exp \left(-\frac{\gamma_{AB} X + \gamma_{CD} Y}{L} \right)$ is expanded in a Taylor series to the second order of X and Y ;

$$E_0 \text{sech}^2 \left(\frac{v_0 t}{2L} \right) \exp \left[\frac{\gamma}{L} (x+Y) \right] = E_0 \text{sech}^2 \left(\frac{v_0 t}{2L} \right) \left\{ 1 + \frac{\gamma}{L} (X+Y) + \frac{1}{2} \left(\frac{\gamma}{L} \right)^2 (x+Y)^2 \right\} \quad (3-22)$$

neglecting high order terms in X, Y . Equation (3-10), with interaction potential given by equation (3-22), can be solved analytically by applying Kerner (19) method. Our semi-classical approach differs from the method of ZRS in that we do not expand the potential function $E_0 \text{sech}^2 \left(\frac{v_0 t}{2L} \right) \exp \left(-\frac{\gamma_{AB} X + \gamma_{CD} Y}{L} \right)$ into a Taylor series. We make a direct numerical integration of the coupled equations (3-21), which are equivalent to equation (3-10). It is appropriate to carry out classical calculations in which the approximate equation (3-8) is used, but the potential expansion (3-22) is not carried out. Numerical comparison between these two methods for N_2-N_2 collision will be given in Chapter 4.

C. Quantum Mechanical Calculation

To clarify the presentation of the theory, we shall use a scaled Schrodinger equation. Define (13):

$$y^* = \left(\frac{\mu_{CD} k_{CD}}{h^2} \right)^{\frac{1}{4}} y \quad (3-23A)$$

$$x^* = \left(\frac{\mu_{AB} k_{AB}}{h^2} \right)^{\frac{1}{4}} x \quad (3-23B)$$

Inserting equations (3-23A) and (3-23B) into equation (3-5A), it becomes:

$$\left\{ \frac{1}{2} h \omega_{CD} \left(-\frac{\partial^2}{\partial y^{*2}} + y^{*2} \right) + \frac{1}{2} h \omega_{AB} \left(-\frac{\partial^2}{\partial x^{*2}} + x^{*2} \right) - \frac{h^2}{2\mu} \frac{\partial^2}{\partial \bar{R}^2} \right. \\ \left. + E_0 \exp \left[\frac{-1}{L} \gamma_{CD} \left(\frac{h^2}{\mu_{CD} k_{CD}} \right)^{\frac{1}{4}} \left(\frac{1}{\gamma_{CD}} \left(\frac{\mu_{CD} k_{CD}}{h^2} \right)^{\frac{1}{4}} \bar{R} - y^* - \frac{\gamma_{AB}}{\gamma_{CD}} \left(\frac{\mu_{CD} k_{CD}}{\mu_{AB} k_{AB}} \right)^{\frac{1}{4}} x^* \right) \right] \right\}$$

$$\psi(x^*, y^*, \bar{R}) = E \psi(x^*, y^*, \bar{R}) \quad (3-24A)$$

Dividing both sides by $1/2h \omega_{CD}$, equation (3-24) looks simpler. If we define:

$$\omega = \frac{\omega_{AB}}{\omega_{CD}}, \quad E^* = \frac{1}{2} \frac{E}{h \omega_{CD}}, \quad r^* = \frac{1}{\gamma_{CD}} \left(\frac{\mu_{CD} k_{CD}}{h^2} \right)^{\frac{1}{4}} \bar{R} \\ L^* = \frac{1}{L} \gamma_{CD} \left(\frac{h^2}{\mu_{CD} k_{CD}} \right)^{\frac{1}{4}} = \frac{1}{L} \frac{m_D}{m_C + m_D} \left(\frac{h^2}{\mu_{CD} k_{CD}} \right)^{\frac{1}{4}} \\ \frac{1}{m} = \frac{h^2}{\mu \omega_{CD}} \left(\frac{\mu_{CD} k_{CD}}{h^2} \right)^{\frac{1}{2}} \frac{1}{\gamma_{CD}^2} = \frac{\mu_{CD}}{\mu} \left(\frac{m_C + m_D}{m_D} \right)^2 = \frac{m_C^M}{(m_A + m_B) m_D} \\ \beta = \frac{\gamma_{AB}}{\gamma_{CD}} \left(\frac{k_{CD} \mu_{CD}}{k_{AB} \mu_{AB}} \right)^{\frac{1}{4}} = \frac{m_C}{m_B} \left(\frac{\mu_{AB}}{\mu_{CD}} \right)^{\frac{1}{2}} \frac{1}{\omega^{1/2}}$$

The new equation (3-24) is dimensionless, namely;

$$\left\{ \left(-\frac{\partial^2}{\partial y^{*2}} + y^{*2} \right) + \left(-\frac{\partial^2}{\partial x^{*2}} + x^{*2} \right) - \frac{1}{\bar{m}} \frac{\partial^2}{\partial j^{*2}} + \left(\frac{E_0}{\frac{1}{2}h\omega_{CD}} \right) \exp[-L^*(r^*-y^*-\beta x^*)] \right\} \psi(x^*, y^*, r^*) = E^* \psi(x^*, y^*, r^*) \quad (3-24B)$$

Since the constant coefficient $\frac{E_0}{\frac{1}{2}h\omega_{CD}}$ can be absorbed into the argument of the exponential function $\exp(-L^*(r^*-y^*-\beta x^*))$, and the operator $\frac{\partial}{\partial r^{*2}}$ is invariant under the transformation

$$r^* \longrightarrow r^* + \text{any constant},$$

we can choose $\frac{E_0}{\frac{1}{2}h\omega_{CD}} = 1$ with no loss of generality. The notation x^* , y^* and r^* refer to dimensionless variables, we now drop the *** for convenience. Hence the scaled Schrodinger equation is:

$$\left\{ \left(-\frac{\partial^2}{\partial y^2} + y^2 \right) + \omega \left(-\frac{\partial^2}{\partial x^2} + x^2 \right) - \frac{1}{\bar{m}} \frac{\partial^2}{\partial r^2} + \exp[-L^*(r-y-\beta x)] \right\} \psi(x, y, r) = E^* \psi(x, y, r) \quad (3-25)$$

The system is specified by the five dimensionless parameters. They are ω , \bar{m} , β , L^* and E^* , E^* is energy of the system in term of ground state energy of molecule CD. The dimensionless form of the unperturbated Hamiltonian H_0 is:

$$H_0(x, y) = \left(-\frac{\partial^2}{\partial y^2} + y^2 \right) + \omega \left(-\frac{\partial^2}{\partial x^2} + x^2 \right)$$

with eigenfunctions $\phi_{ij}(x, y)$ and eigenvalues W_{ij} , then

$$H_0(x, y) \phi_{ij}(x, y) = W_{ij} \phi_{ij}(x, y) \quad \begin{array}{l} i = 0, 1, 2, \dots \\ j = 0, 1, 2, \dots \end{array}$$

where

$$\phi_{ij}(x,y) = \phi_i(y) \phi_j(x)$$

$$W_{ij} = (2i+1) + (2j+1) \omega$$

$\phi_{ij}(x,y)$ is the product of the individual harmonic oscillator wavefunction which indicates that molecule AB with internal coordinate x is the vibrational state j , and molecule CD with internal coordinate y is in the state i . Let the system be in a particular initial state (n_0, m_0) . We can expand the total stationary scattering wavefunction $\psi_{n_0 m_0}(x,y,r)$, i.e. the solution to equation (3-25) in terms of $\phi_{nm}(x,y)$ because they form a complete set.

$$\psi_{n_0 m_0}(x,y,r) = \sum_{n=0}^{N-1} \sum_{m=0}^{M-1} f_{nm, n_0 m_0}(r) \phi_n(y) \phi_m(x) \quad (3-26)$$

Substituting equation (3-26) into equation (3-25) get:

$$\begin{aligned} & \sum_{n=0}^{N-1} \sum_{m=0}^{M-1} \{ W_{nm} f_{nm, n_0 m_0}(r) \phi_n(y) \phi_m(x) - \frac{1}{m} \phi_n(y) \phi_m(x) \frac{d^2}{dr^2} f_{nm, n_0 m_0}(r) \\ & + \exp[-L^*(r-y-\beta x)] f_{nm, n_0 m_0}(r) \phi_n(y) \phi_m(x) \} \\ & = E^* \sum_{n=0}^{N-1} \sum_{m=0}^{M-1} f_{nm, n_0 m_0}(r) \phi_n(y) \phi_m(x) \end{aligned} \quad (3-27)$$

where we have used the relation

$$H_0(x,y) \phi_{ij}(x,y) = W_{ij} \phi_{ij}(x,y)$$

When equation (3-27) is multiplied by $\phi_n^*(y) \phi_m^*(x)$ on both sides and integrated over x and y , we have:

$$\begin{aligned} & \sum_{n=0}^{N-1} \sum_{m=0}^{M-1} \{ W_{nm} f_{nm, n_0 m_0}(r) \langle \phi_{n',m'}(x,y) | \phi_{nm}(x,y) \rangle - \frac{1}{m} \frac{d^2}{dy^2} f_{nm, n_0 m_0}(r) \langle \phi_{n',m'}(x,y) | \phi_{nm}(x,y) \rangle \\ & + f_{n_0 m_0, nm}(r) \langle \phi_{n',m'}(x,y) | \exp[-L^*(r-y-\beta x)] | \phi_{nm}(x,y) \rangle \} \quad (3-28) \\ & = E^* \sum_{n=0}^{N-1} \sum_{m=0}^{M-1} f_{nm, n_0 m_0}(r) \langle \phi_{n',m'}(x,y) | \phi_{nm}(x,y) \rangle \end{aligned}$$

$$\begin{aligned} \text{where } \langle \phi_{n',m'}(x,y) | \phi_{nm}(x,y) \rangle &= \int_{-\infty}^{+\infty} \phi_{n'}(y) \phi_n(y) dy \int_{-\infty}^{+\infty} \phi_{m'}(x) \phi_m(x) dx \\ &= \delta_{n'm'} \delta_{m'm} \end{aligned} \quad (3-29)$$

$$\begin{aligned} &\langle \phi_{n',m'}(x,y) | \exp [-L^*(r-y-\beta x)] | \phi_{nm}(x,y) \rangle \\ &= \exp (-L^*r) \int_{-\infty}^{+\infty} \phi_{n'}(y) \exp (L^*y) \phi_n(y) dy \int_{-\infty}^{+\infty} \phi_{m'}(x) \exp (L^*\beta x) \phi_m(x) dx \end{aligned}$$

$$\text{Defining: } V_{n',m',nm}(r) = \bar{m} \langle \phi_{n',m'}(x,y) | \exp [-L^*(r-y-\beta x)] | \phi_{nm}(x,y) \rangle \quad (3-30)$$

and using equation (3-29) and (3-30), equation (3-28) becomes

$$\begin{aligned} \bar{m} W_{n',m',n_0 m_0}(r) - \frac{d^2}{dr^2} f_{n',m',n_0 m_0}(r) + \sum_{n=0}^{N-1} \sum_{m=0}^{M-1} V_{n',m',nm}(r) f_{nm,n_0 m_0}(r) \\ = \bar{m} E^* f_{n',m',n_0 m_0}(r) \end{aligned}$$

Rearranging this equation, we get

$$\begin{aligned} \frac{d^2}{dr^2} f_{n',m',n_0 m_0}(r) &= \sum_{n=0}^{N-1} \sum_{m=0}^{M-1} V_{n',m',nm}(r) f_{nm,n_0 m_0}(r) \\ &- \bar{m} (E^* - W_{n',m'}) f_{n',m',n_0 m_0}(r) \end{aligned} \quad (3-31)$$

Define

$$\begin{aligned} k_{ij} &= \sqrt{\bar{m} (E^* - W_{ij})} \\ &= \sqrt{\bar{m} (E - (2i+1) - (2j+1))} \end{aligned} \quad (3-32)$$

N and M are the number of states of CD and AB included in the expansion. We introduce (9)

$$\begin{aligned} i &= n + m.N & n &= 0, 1, 2, \dots, N-1 & m &= 0, 1, \dots, M-1 \\ j &= n + m.N & n &= 0, 1, 2, \dots, N-1 & m &= 0, 1, \dots, M-1 \\ k &= n + m.N & n &= 0, 1, 2, \dots, N-1 & m &= 0, 1, \dots, M-1 \end{aligned}$$

to indicate the states of the system. For example, $i = (n, m)$, $j = (n, m)$ and $k = (n, m)$ etc. By incorporating the

definition of k , equation (3-32), we obtain a system of coupled second order ordinary differential equations in matrix form which is equivalent to the Schrodinger equation (3-32).

$$\frac{d^2}{dr^2} \vec{F}(r) = (\vec{V} - \vec{k}^2) \vec{F} \quad (3-33)$$

where $(\vec{F}(r))_{ij} = (\vec{F}(r))_{i(n,m)j(n'm')} = f_{nm,n'm'}(r)$

$$(\vec{V}(r))_{ij} = (\vec{V}(r))_{i(n,m)j(n'm')} = V_{nm,n'm'}(r)$$

$$\text{and } (\vec{k})_{ij} = (\vec{k})_{i(n,m)j(n'm')} = k_{nm} \delta_{mn'} \delta_{mm'}$$

In the asymptotic region where r is very large, $V(r)$ tends to zero, so integrating equation (3-33), we get

$$\lim_{r \rightarrow \infty} \vec{F}(r) = e^{-i\vec{k}r} \vec{G} + e^{i\vec{k}r} \vec{J} \quad (3-34)$$

Equation (3-34) is the asymptotic form of $\vec{F}(r)$ at large r . In principle if \vec{G} and \vec{J} are determined, then the transition probability from state $i = i(n_0, m_0)$ to state $j = j(n, m)$ is given by:

$$P_{ij} = |(\vec{J} \vec{G}^{-1})_{ji}| \frac{k_j}{k_i} \quad (3-35)$$

where \vec{G}^{-1} is the inverse matrix of \vec{G} . However, it is rather difficult to find matrices \vec{G} and \vec{J} in a straightforward manner. We go to the following alternative way:

$$\frac{d}{dr} \vec{F}(r) = \vec{E}(r) \quad (3-36)$$

Set

$$\frac{d}{dr} \vec{E}(r) = (\vec{V}(r) - \vec{k}^2) \vec{F}(r)$$

Letting r_0 be some point in the asymptotic regime, the asymptotic form of $\vec{F}(r)$ may be also written as

$$\lim_{r \rightarrow \infty} \vec{F}(r) = e^{-i\vec{k}(r-r_0)} \vec{G} + e^{i\vec{k}(r-r_0)} \vec{J} \quad (3-37)$$

and then,

$$\lim_{r \rightarrow \infty} \vec{E}(r) = \lim_{r \rightarrow \infty} \frac{d\vec{F}}{dr} = -i\vec{k}e^{-i\vec{k}(r-r_0)} \vec{G} + i\vec{k}e^{i\vec{k}(r-r_0)} \vec{J}.$$

So, as the point $r = r_0$ in the asymptotic region, it is obvious that:

$$\vec{F}(r_0) = \vec{G} + \vec{J} \quad (3-38A)$$

$$\vec{E}(r_0) = -i\vec{k} \cdot \vec{G} + i\vec{k} \cdot \vec{J} \quad (3-38B)$$

These two relations will be used later. From equation (3-37) we have:

$$\lim_{r \rightarrow \infty} \vec{F}(r) \vec{G}^{\dagger-1} e^{-i\vec{k}r_0} = e^{-i\vec{k}r} + e^{i\vec{k}r} (e^{-i\vec{k}r_0} \vec{J} \vec{G}^{\dagger-1} e^{-i\vec{k}r_0}),$$

then the transition probabilities are:

$$P_{ij} = |(e^{-i\vec{k}r_0} \vec{J} \vec{G}^{\dagger-1} e^{-i\vec{k}r_0})_{ji}|^2 \frac{k_j}{k_i} \quad (3-39)$$

Since k_j is real for open channels, that is to say the incoming particle has sufficient energy to excite the bound particle to any of its lowest N eigenstates, matrix \vec{K} is real too, and these are the only observed ones. In equation (3-39) since \vec{K} is real, we end up with

$$P_{ij} = (\vec{J} \vec{G}^{\dagger-1})_{ji} \frac{k_j}{k_i}$$

This result is exactly the same as the probabilities obtained based on the asymptotic form of $\vec{F}(r)$ in equation (3-34). We can find \vec{G} and \vec{J} in terms of $\vec{F}(r_0)$, and $\vec{E}(r_0)$ with no difficulty, since \vec{G} and \vec{J} are related to $\vec{F}(r_0)$, $\vec{E}(r_0)$ by equation (3-38A) and (3-38B). Solving for \vec{G} and

we obtain:

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$$2\vec{G} = \vec{F}(r_0) + i\vec{K}^{-1}\vec{E}(r_0)$$

$$2\vec{J} = \vec{F}(r_0) - i\vec{K}^{-1}\vec{E}(r_0)$$

Defining

$$\vec{S} = \text{Re } \vec{S} + i \cdot \text{Im } \vec{S} = \vec{J}\vec{G}^{-1}.$$

where $\text{Re } \vec{S}$ is the real part of \vec{S} and $\text{Im } \vec{S}$ is the imaginary part of \vec{S} , (both are real) then

$$|\vec{J}\vec{G}^{-1}|^2 = |\vec{S}|^2 = (\text{Re } \vec{S})^2 + (\text{Im } \vec{S})^2$$

We must find $\text{Re } \vec{S}$ and $\text{Im } \vec{S}$ in terms of $\vec{F}(r_0)$, $\vec{E}(r_0)$ and \vec{K} .

We know that (19) given three matrices \vec{A} , \vec{X} , and \vec{Y} such that

$$\vec{A} = \vec{X} + i\vec{Y}$$

where \vec{X} and \vec{Y} are real, if
$$\begin{pmatrix} \vec{X} & \vec{Y} \\ -\vec{Y} & \vec{X} \end{pmatrix}^{-1} = \begin{pmatrix} \vec{Z} & \vec{W} \\ -\vec{W} & \vec{Z} \end{pmatrix}$$

then the inverse matrix of \vec{A} is

$$\vec{A}^{-1} = \vec{Z} + i\vec{W}.$$

In our case $\vec{F}(r_0)$, $\vec{E}(r_0)$ and \vec{K} are all real, we may define

$$\vec{D} = \vec{K}_1^{-1}\vec{E}(r_0)$$

so that equation (3-40) becomes:

$$2\vec{G} = \vec{F}(r_0) + i\vec{D}$$

$$2\vec{J} = \vec{F}(r_0) - i\vec{D}$$

Here we want to find the inverse of \vec{G} , then set:

$$\begin{pmatrix} \vec{F}(r_0) & \vec{D} \\ -\vec{D} & \vec{F}(r_0) \end{pmatrix}^{-1} = \begin{pmatrix} \vec{Z} & \vec{W} \\ -\vec{W} & \vec{Z} \end{pmatrix} \quad (3-41)$$

and solve for \vec{Z} , \vec{W} in terms of $\vec{F}(r_0)$ and \vec{D} . Equation (3-41) means that

$$\begin{pmatrix} \vec{F}(r_0) & \vec{D} \\ -\vec{D} & \vec{F}(r_0) \end{pmatrix} \begin{pmatrix} \vec{Z} & \vec{W} \\ -\vec{W} & \vec{Z} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

Or equivalently

$$\vec{F}(r_0)\vec{Z} - \vec{D}\vec{W} = 1$$

$$\vec{F}(r_0)\vec{W} + \vec{D}\vec{Z} = 0$$

From the second equation of (3-42),

$$\vec{W} = -(\vec{F}(r_0))_{-1} \vec{D} \vec{Z}$$

Substituting \vec{W} into the first equation of (3-42), get:

$$(\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D}) \vec{Z} = 1$$

hence

$$\vec{Z} = (\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1}$$

and

$$\vec{W} = -(\vec{F}(r_0))^{-1} \vec{D}(\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1}$$

$$\text{in turn, } \vec{G} = 2((\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1} - i(\vec{F}(r_0))^{-1} \vec{D}(\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1})$$

The derived result is clear now:

$$\begin{aligned} \vec{J} \vec{G}^{-1} &= (\vec{F}(r_0) - i \vec{D}) ((\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1} - i(\vec{F}(r_0))^{-1} \vec{D}(\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1}) \\ &= \vec{D}(\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1} = \text{Re } \vec{S} + i \text{Im } \vec{S} \end{aligned} \quad (3-43)$$

Since the real part and imaginary part on each side of equation (3-43) are equal to each other respectively, we have:

$$\begin{aligned} \text{Re } \vec{S} &= (\vec{F}(r_0) - \vec{D}(\vec{F}(r_0))^{-1} \vec{D}) (\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1} \\ \text{Im } \vec{S} &= - (2 \vec{D}) (\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1} \end{aligned} \quad (3-44)$$

where $\vec{D} = \vec{K}^{-1} \vec{E}(r_0)$.

The transition probability from state $i = (n_0, m_0)$ to state $j = (n, m)$ is:

$$P_{i(n_0, m_0) \rightarrow j(n, m)} = ((\text{Re } \vec{S})_{ji}^2 + (\text{Im } \vec{S})_{ji}^2) \frac{k_j}{k_i}$$

where $\text{Re } \vec{S}$ and $\text{Im } \vec{S}$ are given by equation (3-44). This completes the basic principle in calculating transition probability for one-dimensional scattering problem within the quantum mechanical approach. Discussions of boundary conditions in integrating equation (3-36) will be made in the next Chapter.

The original Schrodinger equation has rapidly oscillating wavelike solutions which are difficult to represent numerically. The integration of equation (3-33)

is numerically unstable, unless special algorithms are used. Secrest and Johnson, () in their exact quantum mechanical treatment of the one-dimensional scattering problem, convert the coupled differential equations into equivalent integral equations. The integral involved is then replaced by a quadrature sum. The resulting matrix equation is then solved indirectly by numerical method to obtain the transition probabilities. Chan et al. (24), propose a different numerical approach to this problem. It involves converting the set of coupled second-order equations for the translational wavefunctions into first-order equations in matrix form and then solving it by an exponential method developed by W. Magnus (25). The idea was first conceived by Light et al. (26). The method we just discussed for quantum mechanical calculation of transition probabilities is a direct integration of the state expanded Schrodinger equation. This treatment is similar to the method due to Riley and Kupermann (12). It is relatively simple and straightforward, but in our procedures, the virtual states are not included in the total wavefunction expansion. Roy G. Gordon (21) developed another method for integrating coupled differential equations arising in bound state and scattering problems in quantum mechanics. The wavefunctions are constructed in piecewise analytic form, to any prescribed accuracy. The chief advantage of this method is that it avoids searching

for the correct initial derivatives of the wavefunction.
It is claimed to be numerically very stable.

Chapter 4 Numerical Results

A. Semi-classical Results

For the specific N^2-N^2 molecular collision, the two molecules are identical, then $\omega_O = \omega_{AB} = \omega_{CD} = 4.45 \times 10^{-14}$ sec, $\gamma_{AB} = \gamma_{CD}$, and $m_A = m_B = m_C = m_D = 14.0$ a.m.u... This implies that the matrix elements $U_{n',n}$ and $V_{n',n}$ are equal. The matrix elements $U_{n',n}$ are given in Table 4-1.

As mentioned in Chapter 3-B, we may set:

$$\begin{aligned} k &= j + (n-1) J, & j &= 1, 2, 3, \dots \\ & & n &= 1, 2, 3, \dots \end{aligned}$$

where J is an integer which is the number of states of molecules CD and AB included in the expansion of the total wavefunction ψ . The integer k is used to represent the state (j,n) which means that molecule CD is in the state j and molecule AB is in the state n . For example, if we choose $J = 4$, there are $4 \times 4 = 16$ states involved in the expansion of total wavefunction ψ in terms of the individual harmonic oscillator wavefunctions. In other words, there are $16 \times 2 = 32$ coupled first order differential equations to be solved in equation (3-21). (for J greater than 4, the extension is straightforward). In general, $J=N$, let $K=k+N.N$, where $k=1, 2, 3, \dots (N.N)$. In this way, A_{nj} and B_{nj} can be designated as:

$$\begin{aligned} Y(k) &= A_{nj} & k &= 1, 2, \dots (N.N) \\ Y(K) &= B_{nj} & K &= (N.N)+1, (N.N)+2, \dots 2(N.N) \end{aligned}$$

Y is then a vector of length $2(N.N)$. This is the suitable form for doing numerical integration of equation (3-21). In our computer program, the IBM IMSL ROUTINE DGEAR is called. On input, $Y(1), Y(2), \dots, Y(2N*N)$ supply initial values which are initial conditions for the system. One of the arguments in the subroutine DGEAR, TOL, must be chosen suitably. Otherwise the computer time is unnecessarily long. This parameter TOL, is an estimate of the local truncation error. In a series test calculations, we choose $N=11$, $V_0 = 8$ km/sec and initial state = (1,2), for three different values of TOL, 10^{-7} , 10^{-8} , and 10^{-9} . We obtain the data as shown in Table 4-2. We then choose $TOL = 10^{-8}$. The initial value of the step size H, is chosen small enough at the beginning of integration so that it can pass the error test (based on TOL). In the subsequent procedures H is adjusted by the routine itself, but changing in the step size always satisfies the error test. The number of states used in the total wavefunction expansion, $N*N$ plays a very important role in integration of equation (3-21).

In principle, we have to increase N until the final transition probabilities converge to values independent of N. Table 4-3 to Table 4-5 show the transition probability as a function of N for initial state (1,2) at low, medium, and high initial relative velocity; i.e. $v_0 = 3$ Km/sec, 6 Km/sec, and 9 Km/sec respectively. From these tables, it is obvious that for high value of velocity v_0 , we need more

states in the expansion of total wavefunction ψ . Generally speaking, for v_0 less than 6 Km/sec, $N = 7$ i.e. $7 \times 7 = 49$ states expansion is good enough for initial states (1,1), (1,2), (2,2), (3,1), and (3,2). For $v_0 = 7$ Km/sec, N should be no less than 9; and for $v_0 = 8$ Km/sec, N should be no less than 11. For $v_0 = 9$ Km/sec, N must be larger than 12. We also find that the value of N depends on the initial state. For example, for initial state (4,1) at velocity 6 Km/sec, only $N = 9$, i.e. 81 states expansion makes the transition probability converge. The integration limits are adjusted until the constraint equation (4-1) is satisfied: (for suitable N)

$$\left| \sum_{\text{final state}} P_{\text{initial state} \rightarrow \text{final state}}^{-1.0} \right| < \delta \quad (4-1)$$

where δ is an arbitrary small number, we choose $\delta = 10^{-4}$ here. The integration limits depend on the initial relative velocity v_0 , e.g. for $v_0 = 3$ Km/sec, the lower limit $T = -1.6 \times 10^{-13}$ and the upper limit $TEND = 1.6 \times 10^{-13}$; for $v_0 = 6$ Km/sec, $T = -1.0 \times 10^{-13}$, and $TEND = 1.0 \times 10^{-14}$; for $v_0 = 9$ Km/sec, $T = -5.0 \times 10^{-14}$, and $TEND = 5.0 \times 10^{-14}$. Condition (4-1) serves as a useful criterion on numerical calculations.

The numerical results of transition probabilities for different initial and final states as a function of v_0 are obtained by the semi-classical method and shown in Tables

4-6 to 4-10. Figure 4-1 through Figure 4-5 plot the transition probabilities as a function of v_0 for five initial states (1,1), (1,2), (2,2), (3,1) and (3,2). A scaling relationship can be very useful for both the theoretical and experimental analysis of molecular scattering problem. With this relationship, it is easier to write the direct Monte Carlo computation program and save computer storage space. For a systematic study of the scaling relationship, we need more data. It is too expensive to be done at this time. In Figure 4-6a and 4-6b, we plot only the scaling relationship for V-T processes (1,i) (1,i+1) at different v_0 . Since the collision is symmetric, there is nothing new in the results by changing the initial state (i,j) to (j,i). We examine Figure 4-1 through Figure 4-5 and find that at low energies the probability of transferring a given number of quanta by a V-V process is much greater than the probability of converting them by a V-T process into translational energy. For V-T process, the probability increases rapidly with increasing collision energies in the low energy regime. However, the probability of V-V transfer rises less sharply with increasing collision energy. For V-V transitions involving two quanta jumps such as (3,1) (1,3) and (3,2) (1,4), the transition probability is less than that of V-V process (1,2) (2,1), which involve only one quantum jump. Generally, at low energies, the transition probabilities are very small and multiple quantum transitions are assumed

primarily due to stepwise transitions via single collision $j \rightarrow j+1 \rightarrow \dots \rightarrow K-1 \rightarrow K$. At high collision energies the direct transition $j \rightarrow K$ has significant contribution to the transition probability of multiple quantum jump. For V-V-T transfer to an adjacent level, processes involving transfer of a single quantum (such as $(2,1) \rightarrow (3,1)$ and $(2,1) \rightarrow (2,2)$) are much more probable than processes of several quanta such as $(1,2) \rightarrow (3,1)$. Likewise, transitions $(3,2) \rightarrow (4,2)$, $(3,2) \rightarrow (2,2)$ and $(3,2) \rightarrow (3,1)$ are much more probable than the transition $(3,2) \rightarrow (2,1)$. In Figure 4-6a and 4-6b we notice that the general trend of the scaling relationship for transitions

$$(1,i) \rightarrow (1,i+1),$$

seems to be a weak v_0 -dependent function. We need more data for further analysis.

A useful check on numerical results is provided by time-reversal invariance (which leads to the principle of detailed balance). Stated classically, the principle implies that a system executes its motion in reverse if time is allowed to run backward. In quantum scattering processes this means that $P_{ij} = P_{ji}$, i.e., the probability of a transition for state i to state j is equal to that for transition for state j to i . For example, check table 4-8 and table 4-9, we have $P_{(2,2) \rightarrow (3,1)} = 0.845 \times 10^{-3}$, $P_{(3,1) \rightarrow (2,2)} = 0.843 \times 10^{-3}$, this gives

$$\frac{|P_{(3,1) \rightarrow (2,2)} - P_{(2,2) \rightarrow (3,1)}|}{P_{(3,1) \rightarrow (2,2)}} = 0.237\%, \text{ at}$$

$v_0 = 4 \text{ Km/sec.}$

At $v_0 = 6 \text{ Km/sec}$, both $P_{(3,1) \rightarrow (2,2)}$ and $P_{(2,2) \rightarrow (3,1)}$ are equal to 0.121. The principle of time reversal invariance is satisfied quite well.

B. Quantum Mechanical Results and Comparison

Referring to Chapter 3-C, it is clear that the numerical procedures for calculating transition probability by the quantum mechanical method are as follows:

- (1) Integrate equation (3-36) and solve for $\vec{F}(r_0)$ and $\vec{E}(r_0)$.
- (2) Form \vec{D} and the expression $(\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1}$.
- (3) Construct $\text{Re } \vec{S}$ and $\text{Im } \vec{S}$.
- (4) $P_{ij} = ((\text{Re } \vec{S})^2 + (\text{Im } \vec{S})^2)_{ji} \frac{k_j}{k_i}$.

For N^2 - N^2 collisions the parameters of the system are = 0.113 (corresponding to $L = 0.2$ A), $\omega = 1.0$, $\beta = 1.0$, and $\bar{m} = 0.5$. The total energy E can be assigned a suitable value which corresponds to some value of v_0 . Having all these parameters, the IBM IMSL routine DGEAR is called to integrate equation (3-36) and find $\vec{F}(r_0)$, $\vec{E}(r_0)$. There are four factors in this problem that can affect the numerical integration.

- (1) Integration error.
- (2) Number of states retained in the state expansions.
- (3) Starting point of integration r_s .
- (4) End point of integration r_0 .

The local truncation error is controlled by TOL, which is one of the arguments of the subroutine DGEAR. We choose $\text{TOL} = 1.0 \times 10^{-8}$ here. The starting point is chosen as the point $r = r_s$ at which the largest diagonal element of $\vec{V}(r_0)$

is equal to twice of the total energy there. The starting point is just beyond the classical turning point and in the classical forbidden region. Now, it is appropriate to discuss the initial conditions for this system. We set $N = M = 2$, so that there are four states involved in the total wavefunction expansion (for high energy collisions we need larger N and M to get more accurate results). The initial states of the molecule AB and CD are n_0 and m_0 definitely at the beginning, where $n_0 = 1, 2$; $m_0 = 1, 2$. Therefore the initial value of the matrix $\vec{F}(r)$ is a unit matrix:

$$\vec{F}_0 = \begin{pmatrix} 1.0 & 0 & 0 & 0 \\ 0 & 1.0 & 0 & 0 \\ 0 & 0 & 1.0 & 0 \\ 0 & 0 & 0 & 1.0 \end{pmatrix}$$

Since the point $r = r_s$ is in the classical forbidden region, the diagonal element of $\vec{V}(r_s)$ is much larger than the element of K , equation (3-33) becomes:

$$\frac{d^2 \vec{F}}{dr^2} = \vec{V}(r_s) \cdot \vec{F}$$

The diagonal element of \vec{F} , f , satisfies the equation

$$\frac{d^2 f}{dr^2} = m e^{-\alpha r_s} \cdot f(r) \quad \text{as } r \rightarrow -\infty$$

The asymptotic solution is then $\lim_{r \rightarrow -\infty} f(r) = e^{\sqrt{m e^{-\alpha r_s}} r}$. Recall that

$E = \frac{d\vec{F}}{dr}$, this gives the initial value of matrix \vec{E} at starting point r_s as:

$$\vec{E}_0 = \begin{matrix} & g & 0 & 0 & 0 \\ & 0 & g & 0 & 0 \\ & 0 & 0 & g & 0 \\ & 0 & 0 & 0 & g \end{matrix}$$

where $g = \sqrt{me^{-r_s}}$. The stopping point r_0 is chosen as the point where the diagonal element of $V(r_0)$ are less than $\frac{1.0}{5000}$ of the total energy. In principle, we have to increase the total number of states N, M , until the probabilities do not change significantly. We keep $N=2$, and $M=2$ in this work and calculate transition probabilities among states $(1,1)$, $(1,2)$ and $(2,2)$. Obviously, this four-state expansion is good for low energy collisions only. Table 4-11 and Table 4-12 display the quantum transition probabilities as a function of V_0 for initial states $(2,1)$ and $(1,1)$. In Table 4-13 through Table 4-16, we list the probabilities for transitions $(1,2) \rightarrow (1,1)$, $(1,2) \rightarrow (2,2)$, $(1,2) \rightarrow (2,1)$ and $(1,1) \rightarrow (1,2)$ by the four different methods. Since the computer expense is prohibitively large for fully quantum mechanical method, we do not have enough data for plotting purpose. Here, we show that the method just explored does work. The plots of these tables are thus given by Figure 4-7, Figure 4-8, Figure 4-9 and Figure 4-10 respectively.

We check these Figures (4-7) through (4-10) and find that for the three V-T transitions the probability obtained by the semi-classical method is almost one order of magnitude smaller than that obtained from ZRS analytic

method. For V-1 process $(1,2) \rightarrow (2,1)$, the semi-classical and ZRS results are very close to each other, however, the ZRS results are always slightly larger than the semi-classical results. The SSH theory is only good for low energy collisions, if the energy is too high the probability is greater than unity. This theory breaks down there. Since we are interested in high energy collision processes of two molecules. We concentrate on the semi-classical method. We believe the semi-classical treatment can supply a reasonable estimate in calculating transition probabilities. This is very helpful because the semi-classical method can save much computer time and the numerical algorithm is relatively simple.

Chapter 5 Discussions

In this final chapter, we discuss some important problems requiring further study for vibrational energy transfer.

A. Three-Dimensional Collisions

One of the assumptions which we have made in the collinear molecular collision model is that the target molecule is struck in the direction of its axis. To avoid this assumption in the collinear treatment, we have to average over the relative orientation of the molecule at the proper stage of calculations. However, the period of rotation is usually comparable with the duration of the collision, there is no simple way to take the average. A constant steric factor is generally used.

Since the rotational energy spacing is much smaller than the vibrational spacing, appreciable rotational scattering occurs over a range of molecule-molecule separations that is considerably longer than that for which vibrational transitions are important. The coupling between rotational and translation is usually strong too, so that the rotational state generally changes before the vibrational transition occurs. When a vibrational transition takes place the corresponding energy change will appear in either translational motion, or rotational motion of molecules, or both. It is obvious that if we calculate the vibrational transition probabilities,

effects of rotational motion have to be considered. In the collinear treatment however, we have assumed that the simultaneous rotational and vibrational transitions are not important and the impact parameter is zero. The realistic three-dimensional analyses that take rotational transitions into account should include the correction to the vibrational transition probability that results from the finite size of the rotational energy spacing in future work. Also, the incident particle is described by a plane wave which contains the partial waves of different orbital angular momentum (the one-dimensional model corresponding to an s-wave scattering problem). Usually, many partial waves have to be considered and this makes the problem very difficult.

B. Effect of Anharmonicity

It has been found by experience that the potential energy function of actual diatomic molecules can be represented quite accurately by a simple analytical function called Morse potential, which contains three adjustable parameters. If Morse potentials are used to describe the intramolecular forces, the diagonal matrix elements of the interaction potential which enter into the quantum theory of vibrational energy transfer are approximately but not identically equal. In the calculation given by F.H. Mies, (27) the consideration was restricted to the head-on collision between a structureless incident particle and a diatomic molecule. The transition

probability is found to decrease markedly when the ratio of the diagonal elements of the initial and final oscillator states is allowed to deviate even slightly from one. The deviation in turn, increases with the anharmonicity of the molecular vibrations, and an anharmonic correction factor of the order of 10^{-1} to 10^{-2} should be applied to the generally used probability expression for atom-molecule collision. There must exist a correction factor of this kind for molecule-molecule collision.

C. Interaction Potential

Choice of a potential function to be used in calculating the transition probability is a very important task since it affects the results considerably. In the theory of inelastic molecular collisions, the scattering potential to be adopted should be simple enough to make the calculations feasible as long as the essential features of the physics of collision is not lost. This requirement is relaxed if we deal with numerical calculations. The chosen interaction potential for some pair of molecules must be relatively accurate and can be used to represent the real situation. If the intermolecular interaction is strongly orientation dependent, as in the polar gas, the molecules may take a particular orientation during the encounter. For this problem of preferential orientation, a somewhat different treatment is required.

D. Exact Classical Trajectory

In our semi-classical approximation, the classical equations are first solved to obtain the relative motion of the molecules as a function of the time. The time-dependent Schrodinger equation for the internal motion under the external perturbation is then solved to obtain the probabilities of various transitions. However, the occurrence of inelastic processes are not taken into account in solving the classical equation of motion in that the effective intermolecular potential and the effective translational energy depend on the internal state. If the incident energy is much greater than the internal energy, the influence of inelastic process on the relative motion is unimportant. It is a good approximation to ignore the internal state in calculating classical trajectory. For more rigorous calculations, an exact classical trajectory must be found in which the energy conservation law is satisfied. For high energy collisions, the semi-classical treatment is, however, a fairly good approximation. It requires less computational effort and saves much computer time.

Appendix 1 Exponential Interaction Potential

A conventional representation of the intermolecular potential energy curve is given by the Lennard-Jones 12-6 power law (28);

$$V(r) = 4\epsilon \left(\left(\frac{\sigma}{r} \right)^{12} - \left(\frac{\sigma}{r} \right)^6 \right) \quad (A1-1)$$

where $V(r)$ is the potential energy at separation r , and r is the distance between atom B and C. This is shown graphically in Figure A1-1.

is the depth of the potential well at intermolecular distance r_m , where the repulsive force $\left(\frac{\sigma}{r} \right)^{12}$ takes over the long range attractive force $\left(\frac{\sigma}{r} \right)^6$, and $V(r_m)$ is the minimum of the potential function $V(r)$. σ is the

separation at zero energy, when $V(r) = 0$, sometimes loosely called the "collision diameter". The exponential function

$$V_{\text{INT}}(r) = \text{constant} \cdot e^{-r/L} - \epsilon \quad (\text{A1-2})$$

must be fitted to the Lennard-Jones potential $V(r)$, equation (A1-1). Here, the choice is made that the magnitudes and slopes of the potentials are set equal at $r = r_c$. r_c is the minimum value of r . These two potentials are illustrated in Figure A1-2.

Figure A1-2. The exponential potential $V_{\text{INT}}(r)$ fitted to the Lennard Jones potential $V(r)$. The magnitudes and slopes of the two potentials are set equal to each other at the classical turning point $r = r_c$.

We deduce an approximate formula:

$$L = \frac{\sigma}{17.5}$$

For N_2 molecule, $\sigma = 3.749 \text{ \AA}$, so $L = 0.21 \text{ \AA}$.

Appendix 2 The Choice of Initial Time Reference

Coordinate in Quantum Scattering Process

Consider the general scattering in one-dimensional space, A flux of incoming particles with mean momentum p_0 are incident from left and scattered by an arbitrary potential distribution $V(x)$ as shown in Figure A2-1, where $V(x)$ is finite and $V(x) \rightarrow 0$ as $x \rightarrow \pm\infty$

Figure A2-1. Particles scattered by an arbitrary potential.

For large and negative x , the wave packet with mean momentum p_0 can be superposed as:

$$\begin{aligned} \psi(x,t) = & \int_{-\infty}^{+\infty} dp \exp(-\alpha(p-p_0)^2) \exp\left(\frac{ipx}{\hbar}\right) \exp\left(-i \frac{p^2}{2m\hbar} t\right) \\ & + \int_{-\infty}^{+\infty} dp R(p) \exp(-\alpha(p-p_0)^2) \exp\left(-\frac{ipx}{\hbar}\right) \exp\left(-i \frac{p^2}{2m\hbar} t\right) \end{aligned} \quad (A2-1)$$

where $R(p)$ is the reflection coefficient which is a constant over a region $\Delta p \sim \frac{1}{\sqrt{\alpha}}$. For large and positive x , the transmitted wave packet is:

$$\psi(x,t) = \int_{-\infty}^{+\infty} dp T(p) \exp(-\alpha(p-p_0)^2) \exp\left(\frac{ipx}{h}\right) \exp\left(-i \frac{p^2 t}{2mh}\right) \quad (A2-2)$$

where $T(p)$ is the transmission coefficient which is constant over a region $\Delta p \sim \frac{1}{\sqrt{\alpha}}$. Since for large $|x|$ and $|t|$, the term $\exp\left(i\left(\frac{px}{h} - \frac{p^2 t}{2mh}\right)\right)$ in equation (A2-1) and (A2-2) is a very rapidly varying function of momentum p , the integrals are essentially zero unless p , x , and t are such that the stationary phase conditions are satisfied:

$$\begin{aligned} p &\approx p_0 \\ \frac{\partial}{\partial p} \left(\frac{px}{h} - \frac{p^2 t}{2mh} \right) &\approx 0 \end{aligned} \quad (A2-3)$$

From equation (A2-3) we get:

$$x \approx \frac{p_0}{m} t \quad (A2-4)$$

where $\frac{p_0}{m}$ is the classical velocity. Equation (A2-4) gives the result that $t \ll 0$, if $x \ll 0$. Let's check the second term in the right hand side of equation (A2-1). Since $R(p)$ is approximately a constant within a width $\Delta p \sim \frac{1}{\sqrt{\alpha}}$ centered at $p = p_0$, then,

$$\begin{aligned} &\int_{-\infty}^{+\infty} dp R(p) \exp(-(p-p_0)^2) \exp\left(-i \frac{px}{h}\right) \exp\left(-i \frac{p^2 t}{2mh}\right) \\ &= R(p_0) \int_{-\infty}^{+\infty} dp \exp(-(p-p_0)^2) \exp\left(-i \left(\frac{px}{h} + \frac{p^2 t}{2mh}\right)\right) \end{aligned}$$

Stationary phase conditions require that:

$$P \approx P_0$$

and

$$\frac{\partial}{\partial p} \left(\frac{px}{h} + \frac{p^2 t}{2mh} \right) \approx 0 \quad (A2-5)$$

Equation (A2-5) implies that $x \approx \frac{-p_0 t}{m}$, this means that reflection occurs only when $t > 0$ because the reflected wave exists only at $x \longrightarrow -\infty$. Combining the discussions just made, we conclude that the incident particles hit the potential at $x \approx 0$ and $t \approx 0$, and the initial conditions of the system are described at $t \longrightarrow -\infty$, the final conditions are the states at $t \longrightarrow +\infty$.

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REFERENCES

1. G.A. Bird, "Shock Wave Structure in a Rigid Sphere Gas"
Academic Press, Vol. 1, P. 216, New York 1965.
2. D.I. Pullin, J.K. Harvey, and G.K. Bienkowski
"Hypersonic Leading Edge Flow of a Diatomic Gas by the
Direct Simulation Method".

APPENDIX C

COMPUTER CODE INTERNAL

```
// JOB GKB 0367425.GKBSPACE N=WATRUN2 REG=560 T=1.0 P=100 C=0
// EXEC WATFIV
//WATFIV.FT09F001 DD DISP=OLD,DSN=U.GKBSPACE.KM115
//WATFIV.SYSIN DD DATA
$JOB      BIENKOWSKI,T=59,P=100,NOLIST
C      MAIN PROGRAM FOR MONTE CARLO 3-D ENTRANCE PROBLEM CALCULATIONS
C      OBJECTIVE OF THIS MAIN PROGRAM IS TO SET THE DIMENSIONS
C      MAIN RUNNING PROGRAM IS *** RUN ***
C
C      FOLLOWING TWO CARDS HAVE TO BE ELIMINATED FOR NON IBM MACHINES
C*****
C      INTEGER*2 LB,NBM,NBN,NB,NBT
C      INTEGER*2 LM,LCOL
C*****
C
C      THE NEXT CARD IS ASSOCIATED WITH PRINCETON RANDOM NUMBER GENERATOR
C*****
C      COMMON/RANCOM/NRAN(4)
C*****
C
C      THE FOLLOWING DIMENSION STATEMENTS SET THE MAJOR ARRAY DIMENSIONS
C      AND MUST BE CONSISTENT WITH THE FOLLOWING DATA CARD -----
C
C      NSP=NUMBER OF SPECIES - EXAMPLE BELOW NSP=1
C      NJV=NUMBER OF SUBDIVISIONS OF INPUT DISTRIBUTION FUNCTION
C      EXAMPLE BELOW NJV=22
C      NMC=NUMBER OF FINAL CELLS - EXAMPLE BELOW NMC=150
C      NMP=MAX NUMBER OF MOLECULES OF EACH SPECIES ALLOWED IN PROGRAM.
C      IF EXCEEDED, PROGRAM EITHER FAILS OR RESTARTS AT BEGINNING
C      WITH NUMBER REDUCED BY 10% - EXAMPLE BELOW NMP=5000
C      NPB=MAXIMUM NUMBER IN EACH CELL - EXAMPLE NPB=150
C
C      DIMENSION DBA(1,150),NB(1,150),NBT(1,150)
C      DIMENSION TMP(1,150),TMPA(1,150),XV(1,150),XVA(1,150)
C      DIMENSION YV(1,150),YVA(1,150),ZV(1,150),ZVA(1,150),DB(1,150)
C      DIMENSION TRP(1,150),TRPA(1,150),NBM(1,150)
C      DIMENSION NBN(150),T(1,1,150)
C      DIMENSION LB(5000),LM(1,5000),ER(1,5000)
C      DIMENSION PAU(1,5000),PAV(1,5000),PAW(1,5000)
C      DIMENSION PAX(1,5000),PAY(1,5000),PAZ(1,5000),LCOL(1,5000)
C      DIMENSION PNB(150),XC(150),YC(150),ZC(150)
C      DIMENSION VEL(22,4,1),PFV(22,4,1)
C      DATA NSP/1/,NJV/22/,NMC/150/,NMP/5000/,NPB/150/
```

```
2 FORMAT(/17X,'NORMAL TERMINATION OF THE PROGRAM')
NAMELIST/DIM/NSP,NJV,NMC,NMP,NPB,NRAN
```

INITIALIZATION OF RANDOM NUMBER GENERATOR - PRINCETON ROUTINE

```
NRAN(1)=0
NRAN(2)=0
NRAN(3)=0
NRAN(4)=0
```

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PRINTOUT OF MAJOR ARRAY DIMENSIONS USED ABOVE

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WRITE(6,DIM)

CALL OF MAIN OPERATING PROGRAM WHICH REQUIRES INPUTS:
 &CONTRL,&TIMES,&FLOREF,&MOLEC,&SHAPES,&GEOM,&INCUPL,&INOUT
 THESE INPUTS ARE ALL CURRENTLY IN THE NAMELIST FORMAT
 AND MAY HAVE TO BE CHANGED IF THAT CONVENTION IS NOT AVAILABLE
 BRIEF DESCRIPTION OF THE PARAMETERS FOLLOWS

&CONTRL - ONE OCCURRENCE (NEW OR RESTART)

PARAMETER	DEFAULT	DEFINITION OR EXPLANATION
NAME	8 BLANKS	ANY ALPHANUMERIC NAME UP TO 8 CHARACTERS
TITLE	24 BLNKS	ANY ALPHANUMERIC TITLE UP TO 24 CHARACTERS
PERCNT	.001	ACCURACY IN INTEGRATION PROCEDURES
ICOPY	1	NUMBER OF ADDITIONAL COPIES OF OUTPUT
DUMP	.TRUE.	IF TRUE WILL CAUSE SYSTEM DUMP FOR ANY OF 12 PROGRAMMER DESIGNED ERROR HALTS.
DEBUG (1)	.FALSE.	IF TRUE WILL PRINT MESSAGE WHEN CELL POP. EXCEEDS MNB
DEBUG (2)	.FALSE.	IF TRUE WILL PRINT CPU TIME AROUND EACH PART OF LOOP
DEBUG (3)	.TRUE.	IF TRUE WILL PRINT CPU TIME REMAINING AT END OF LOOP
NEW	.TRUE.	IF TRUE - NEW RUN - IF FALSE - RESTART OF RUN
SAVE	.FALSE.	IF TRUE - SNAPSHOT SAVED ON TAPE(9) FOR RESTART
REDO	.FALSE.	IF TRUE PROGRAM WILL AUTOMATICALLY RESTART WITH 90% OF TOTAL IF TOTAL CELL POPULATION EXCEEDS MNM

&TIMES - ONE OCCURRENCE (NEW OR RESTART)

PARAMETER	DEFAULT	DEFINITION OR EXPLANATION
DTM	- - -	REAL NUMBER - FRACTION OF MEAN FREE TIME PER CYCLE
ITS	- - -	INTEGER - NUMBER OF CYCLES PER SAMPLE
ITP	- - -	INTEGER - NUMBER OF CYCLES BETWEEN PRINTOUTS
TST	- - -	INTEGER - ESTIMATE OF NUMBER OF CYCLES TO STEADY STATE
TLIM	- - -	INTEGER - TOTAL NUMBER OF CYCLES TO END OF RUN - WILL TERMINATE SOONER IF CPU TIME IS TO BE EXCEEDED

&FLOREF - ONE OCCURRENCE (NEW RUN ONLY)

PARAMETER	DEFAULT	DEFINITION
LLM	- - -	INITIAL NUMBER OF MOLECULES $LLM < MNM < \text{OR} = NNP$
MNM	- - -	MAXIMUM NUMBER OF MOLECULES PER SPECIES
MNB	- - -	MAXIMUM NUMBER PER CELL - DIAGNOSTIC ONLY
MSP	- - -	NUMBER OF MOLECULAR SPECIES (MAX. IS 3)
MET	0	IF 0 - DATA IS IN SI (METRIC) UNITS IF > 0 - DATA IS IN ENGLISH UNITS
U	- - -	FLOW VELOCITY (M/SEC) OR (FT/SEC)
ANGLE	- - -	ANGLE OF ATTACK (DEGREES)
RNU	0.0	ARRAY GIVING MOLE FRACTIONS OF SPECIES IN FREE STREAM
RMA	0.0	ARRAY GIVING MOLECULAR WEIGHTS OF SPECIES ABOVE
TF	- - -	FREE STREAM TEMPERATURE (K OR R)
DENF	- - -	FREE STREAM NUMBER DENSITY (NUM/M**3 OR NUM/FT**3)

&MOLEC - ONE OCCURRENCE (NEW RUN ONLY)

PARAMETER	DEFAULT	DEFINITION
TRF	- - -	REFERENCE TEMPERATURE FOR MOLECULAR DATA
DIR	0.0	CROSS-SECTIONS AT REFERENCE TEMP. (MSPX/MSP)
ETA	0.0	PARAMETERS IN DIFFUSION AND VISCOSITY LAW (MSPX/MSP)
PHI	0.0	PARAMETERS FOR ROTATIONAL RELAXATION (MSPX/MSP)

CHI 0.0 ROTATIONAL DEGREE OF FREEDOM PARAMETER (NROT/2 - 1)
 ACR .001 ACCURACY IN MOLECULAR COLLISION CALCULATIONS

ESHAPES - ND+1 OCCURRENCES WHERE ND=NUMBER OF BODY SEGMENTS (NEW RUN)

PARAMETER DEFAULT DEFINITION

FIRST OCCURRENCE

BODY(I) I>3 - - NEED NOT BE SPECIFIED
 BODY(1) - - - START OF BODY (XSTART) IN ARBITRARY COORDINATE
 BODY(2) 0.0 TEMPERATURE AT FRONT OF TUBE IN K OR R
 BODY(3) 0.0 DIAMETER OF TUBE IN METERS OR FT.
 SUBSEQUENT OCCURRENCES (ND)
 BODY(1) - - - X COORDINATE FROM FRONT OF BODY OF THE DOWNSTREAM
 EDGE OF THE CURRENT BODY SEGMENT
 BODY(2) - - - TEMPERATURE AT THE BACK OF THIS BODY SEGMENT
 BODY(3) - - - SWITCH - IF NOT 0.0 THIS IS THE LAST SHAPES CARD
 BODY(I) I EVEN ALPHA - ENERGY ACCOMODATION COEFFICIENT FOR SPECIES
 BODY(J) J ODD SIGMA - TANGENTIAL ACCOMODATION COEFF. FOR SPECIES
 I AND J < (4+2*MSP)

EGEOM - ONE OCCURRENCE (NEW RUN ONLY)

PARAMETER DEFAULT DEFINITION

NWEDG - - - INTEGER GIVING THE NUMBER OF WEDGES WITHIN 180 DEGREES
 NW - - - NUMBER OF FIRST LEVEL CELLS IN X DIRECTION
 NH - - - NUMBER OF FIRST LEVEL CELLS IN RADIAL DIRECTION

SINCUPL - ONE OCCURRENCE (NEW RUN ONLY) - INPUT DISTRIBUTION

PARAMETER DEFAULT DEFINITION

FLUXIN 1.0 FLUX INPUT IN TERMS OF FREE STREAM FLUX - ONE
 NUMBER FOR EACH SPECIES
 PCOL 0.0 FRACTION OF ARRIVING MOLECULES THAT HAVE
 PREVIOUSLY COLLIDED
 RNP - - - RATIO OF "CAVITY" PRESSURE TO THE EFFECTIVE
 PRESSURE OF THE INCOMING STREAM AT ENTRANCE
 JV - - - THE NUMBER OF VELOCITY INTERVALS FOR DISTRIBUTION
 FUNCTION INFORMATION
 KMX 4 NUMBER OF COMPONENTS OF DISTRIBUTION
 KMX=3 IF NO ROTATIONAL ENERGY (CHI=-1)
 KMX=4 IF ROTATIONAL ENERGY IS INCLUDED (CHI>-1)

SINOUT - MSP*KMX OCCURRENCES

PARAMETER DEFAULT DEFINITION

MT DESIGNATES SPECIES
 K=1 DESIGNATES NORMAL VELOCITY
 K=2 DESIGNATES TANGENTIAL VELOCITY IN FLOW DIRECTION
 K=3 DESIGNATES TRANSVERSE TANGENTIAL VELOCITY
 K=4 DESIGNATES ROTATIONAL ENERGY
 VARG(J) - - - VELOCITY BOUNDARIES FOR DISTRIBUTION FUNCTION
 1<J<JV VELOCITY BOUNDARIES
 CURV(J) - - - PROBABILITY OF VELOCITY (OR ROTATIONAL ENERGY)
 INCIDENT AT ENTRANCE BEING BELOW VARG(J)

A SAMPLE INPUT DECK IS GIVEN BELOW:

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&CONTRL NAME='INTE','RNAL',TITLE=' PAR','ABOL','A AT',' 95K','M M','ON. ',
  DEBUG=.F.,.F.,.T.,NEW=.T.,SAVE=.F.,ICOPY=0,REDO=.T. &END
&TIMES DTM=.010,ITS=5,ITP=1000,TST=400,TLIM=1000 &END
&FLOREP LLM=2000,MNM=5000,MNB=150,MSP=1,NET=0,U=7485.9,ANGLE=28.,RNU=1.,2*0.,
  RMA=28.94,0.,0.,TF=195.51,DENF=2.52E+19 &END
&MOLEC TRP=1000,DIR=3.5E-19,ETA=.104,PHI=0.0,CHI=-1.,ACR=.001 &END
&SHAPES BODY=0.0,1000.,.00235 &END
&SHAPES BODY=.0025,555.,0.0,2*1.0 &END
&SHAPES BODY=.0050,345.,0.0,2*1.0 &END
&SHAPES BODY=.0100,300.,0.0,2*1.0 &END
&SHAPES BODY=.0200,300.,0.0,2*1.0 &END
&SHAPES BODY=.0300,300.,0.0,2*1.0 &END
&SHAPES BODY=.0400,300.,0.0,2*1.0 &END
&SHAPES BODY=.0500,300.,0.0,2*1.0 &END
&SHAPES BODY=.0600,300.,0.0,2*1.0 &END
&SHAPES BODY=.0700,300.,0.0,2*1.0 &END
&SHAPES BODY=.0800,300.,0.0,2*1.0 &END
&SHAPES BODY=.0870,300.,1.0,2*1.0 &END
&GEOM NWEDG=2,NW=20,NH=3, &END

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&INCUPL FLUXIN=2.1429,PCOL=1.0,RMP=0.0,JV=22,KMX=3 &END
&INOUT VARG=0.,1.,2.,3.,4.,5.,6.,7.,8.,9.,10.,11.,12.,13.,14.,15.,16.,17.,18.,
  19.,20.,21.,CURV=0.0,.070,.170,.282,.369,.459,.537,.599,.656,.710,.750,.785,
  .815,.845,.872,.900,.922,.951,.975,.988,.996,1.00, &END
&INOUT VARG=-20.,-19.,-17.,-15.,-13.,-11.,-9.,-7.,-5.,-3.,-1.,1.,3.,5.,7.,9.,
  11.,13.,15.,17.,19.,CURV=4*0.,.003,.013,.036,.084,.149,.250,.406,.611,.762,
  .871,.932,.962,.984,.995,.999,3*1.0, &END
&INOUT VARG=-20.,-19.,-17.,-15.,-13.,-11.,-9.,-7.,-5.,-3.,-1.,1.,3.,5.,7.,9.,
  11.,13.,15.,17.,19.,CURV=4*0.,.003,.013,.036,.084,.149,.250,.406,.611,.762,
  .871,.932,.962,.984,.995,.999,3*1.0, &END

```

```

CALL RUN(NSP,NJV,NMC,NMP,NPB,DBA,NB,NBT,TMP,TMPA,XV,XVA,YV,
  1YVA,ZV,ZVA,T,DB,FNB,XC,YC,ZC,LM,PAU,PAV,PAW,PAX,PAY,
  2PAZ,LCOL,TRP,TRPA,ER,LB,NBM,NBN,VEL,PFV)
WRITE(6,2)
STOP
END

```

```

SUBROUTINE RUN(NSP,NJV,NMC,NMP,NPB,DBA,NB,NBT,TMP,TMPA,XV,XVA,
  1 YV,YVA,ZV,ZVA,T,DB,FNB,XC,YC,ZC,LM,PAU,PAV,PAW,PAX,PAY,PAZ,
  2LCOL,TRP,TRPA,ER,LB,NBM,NBN,VEL,PFV)
  MAIN RUNNING PROGRAM ** RUN *** CALLS ALL OTHER SUBROUTINES

```

```

INTEGER*2 LM,LCOL
INTEGER*2 LB,NBM,NBN,NB,NBT
INTEGER PRT,SAMP,TST,TLIM,TIME,Q
LOGICAL DUMP,DEBUG(3),SAVE,NEW,REDO
REAL INTGRL,LAM,MU,NU,JAY,KAY
DIMENSION BTA(3),C1(3),C2(3),C3(3),C7(3),C8(3),DPA(3),FL(3)
DIMENSION FDN(3),HTI(3),HTR(3),JNT(3),KNM(3),NM(3),SR(3)
DIMENSION NAME(2),TITLE(6)
DIMENSION RNU(3),RMA(3),WTM(3),CHI(3),DIR(3,3),DAM(3,3),PHI(3,3)
DIMENSION ETA(3,3),CN8(3,3),CNG(3),CHG(3),CN(3,3,3),CH(3,3,3)

```


DIMENSION CTI (3,3),CTR (3,3),CNI (3,3),CNR (3,3),SN (3),ST (3)
 DIMENSION D1 (3),D2 (3),D3 (3),D4 (3),BODY (15),DBG1 (3,3),LIMIT (10)
 DIMENSION COEPP (4)
 DIMENSION XLIM (2),NCOL (3,3)
 DIMENSION XCB (18),XS (18),YCB (18),TB (18),ALPHA (3,18),SIGMA (3,18)
 DIMENSION WTS (3,18,12),UTL (3,18,12),UTT (3,18,12)
 DIMENSION VTS (3,18,12),HTSI (3,18,12),HTS (3,18,12)
 DIMENSION UTLI (3,18,12),UTTI (3,18,12),VTSI (3,18,12)
 DIMENSION ENT (3,2),REM (3,2),FLUXIN (3),PCOL (3)
 DIMENSION VARG (42),CURV (42),IPLUX (3,2)
 DIMENSION LB (NMP),NBN (NMC),NBH (NSP,NMC),LM (NSP,NMP)
 DIMENSION ER (NSP,NMP),TRP (NSP,NMC),TRPA (NSP,NMC)
 DIMENSION DBA (NSP,NMC),NB (NSP,NMC),NBT (NSP,NMC)
 DIMENSION TMP (NSP,NMC),TMPA (NSP,NMC),XV (NSP,NMC),XVA (NSP,NMC)
 DIMENSION YV (NSP,NMC),YVA (NSP,NMC),ZV (NSP,NMC),ZVA (NSP,NMC)
 DIMENSION T (NSP,NSP,NMC),DB (NSP,NMC)
 DIMENSION FNB (NMC),XC (NMC),YC (NMC),ZC (NMC)
 DIMENSION VEL (NJV,4,NSP),PFV (NJV,4,NSP)

DIMENSION PAU (NSP,NMP),PAV (NSP,NMP),PAW (NSP,NMP)
 DIMENSION PAX (NSP,NMP),PAY (NSP,NMP),PAZ (NSP,NMP),LCOL (NSP,NMP)

RUN0460

COMMON /RANCOM/NRAN (4),KAWLS
 COMMON /FIRST/NL,NW,NH
 COMMON /SECND/BW,BH,RMP,REN,RMF
 COMMON /THIRD/PI,NREG,S,SINANG,COSANG,AKN,AKT,AKN1,AKN2,AKT1,AKT2
 COMMON /FORTH/NBX,RM,XR,DUMP,C9,LL (3),LLM
 COMMON /FIFTH/ND,TIME,DTM,TL,ITS,ITP,TST,TLIM,RMA,RNU,DIR
 COMMON /SIXTH/RMB,XSTART,INM,MNM,MNB,NEW,SAVE,PERCNT,NSR,TR
 COMMON /SVNTH/LAM,MU,NU,MT,N,J,X,Y,Z,TUSE
 COMMON/EIGHT/DENF,U,TF,ANGLE,TRF,CHI,PHI,ETA,WTM,DAM,VELR,XREF
 NAMELIST/CONTRL/NAME,TITLE,PERCNT,ICOPY,DUMP,DEBUG,NEW,SAVE,REDO
 NAMELIST/TIMES/DTM,ITS,ITP,TST,TLIM
 NAMELIST/FLOREF/LLM,MNM,MNB,MSP,MET,U,ANGLE,RNU,RMA,TF,DENF
 NAMELIST/MOLEC/TRF,DIR,ETA,PHI,CHI,ACR
 NAMELIST/SHAPES/BODY
 NAMELIST/GEOM/NWEDG,NW,NH
 NAMELIST/INCUP/FLUXIN,PCOL,RMP,JV,KMX
 NAMELIST/INOUT/VARG,CURV

RUN0520

RUN0530

RUN0540

RUN0580

RUN0570

DATA IC/0/,ICOPY/1/
 DATA DBG1/' GAS',' AT ','110 ','FLOW',' AT ','130 ',' RUN',' AT ',
 1'303 '/
 DATA LIMIT/12,9,18,500,3600,70,900,3,20,3/
 DATA TITLE/' ',' ',' ',' ',' ',' ',' '
 DATA NAME/' ',' '
 DATA CPC/0.0/,CPM/0.0/,CPB/0.0/,CPJ/15.0/

RUN0630

RUN0640

RUN0660

RUN0670

RUN0680

RUN0690

RUN0700

RUN0710

RUN0720

RUN0730

RUN0740

RUN0750

FORMATS

1 FORMAT (1H1)

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2 FORMAT(1H1/17X,'RARIIFIED SUPERSONIC FLOW OF BINARY GAS',T74,'I') RUN0760
3 FORMAT('+',103X,'COPY ',I2) RUN0770
4 FORMAT(/17X,'FLOW THROUGH ALL THE BOUNDARIES'/5X,'MT',5X,'MASS COL
1E PR.',10X,'FCOL. FLUXIN FLUXES (ENT) ')
5 FORMAT(3X,I4,P9.2,P9.4,5X,4P10.4)
25 FORMAT(/5X,'PRESSURE RATIO (INSIDE/ENTRANCE) - EITHER TYPE =',P13.5
1,' RMP',T76,'I'/6X,'DENSITY RATIO (INSIDE/ENTRANCE) - EITHER TYPE
2=',P13.5,' RMN',T76,'I'/9X,'FLUX RATIO (INSIDE/ENTRANCE) - EITHER
3TYPE =',P13.5,' RMP',T76,'I'/)
26 FORMAT(/10X,'FLUX RATIOS FOR SPECIES RMA =',P7.2/2X,'BOUNDARY',
1' INFLUX REFLUX NET PLUX NET PLUX/RHO*U'/2X,'ENTRANCE',
24P10.4/4X,'CAVITY',4P10.4/)
30 FORMAT('1TIME =',P6.3,60X,'RANDOM NUMBER GENERATOR HAS BEEN CALLED
1',I0,' TIMES') RUN0800
31 FORMAT(' CPU TIME LEFT- ',P8.3) RUN0810
32 FORMAT(7X,'-MOLECULES-'/3X,3I6)
33 FORMAT(' TIME = ',P8.3,5X,'COLLISION LOOP=',P8.3,5X,'MOVE LOOP = '
1,P8.3,5X,'TOTAL TIME = ',P8.3/21X,'2ND MOVE LOOP = ',P8.3,5X,
2'CLEANUP LOOP=',P8.3,4X,'PARTICLE NUMBERS = ',4I6)
34 FORMAT(9X,'-MOLECULAR COLLISIONS-'/3(3I14/))
35 FORMAT(2X,'-COLLISIONS WITH SURFACE-'/3X,3I8)
36 FORMAT(' MAXIMUM NUMBER OF MOLECULES SO FAR-',I6//) RUN0880
38 FORMAT(' EXCESS MOLECULES OCCURRED IN ',3A4) RUN0890
40 FORMAT('/' SOMETHING IS WRONG WITH BOX NUMBERING IN RUN '/9I5,5E14.RUN0900
15)
44 FORMAT(' NB(' ,I2,' ',',I4,' ) POPULATION EXCEEDED ',I3,' IN MAIN AT TRUN0930
1TIME = ',P7.3) RUN0940
50 FORMAT(///' SNAP SAVED ON TAPE') RUN0950
***** RUN0960
***** RUN0970
***** RUN0980

```

```

CPA=ELTIME(0)
CALL NOUNDF
CALL TRAPS(0,1,1000000,1,1)
LIMIT(4)=NMC
LIMIT(5)=NMP
LIMIT(6)=NPF
LIMIT(7)=NJV
LIMIT(10)=NSP
KAWLS=0
PI=3.141593
PIROOT=SQRT(PI)
MET=0
LARGE=0
NL=1
DUMP=.TRUE.
DEBUG(1)=.FALSE.
DEBUG(2)=.FALSE.
DEBUG(3)=.TRUE.
SAVE=.FALSE.
NEW=.TRUE.
REDO=.FALSE.
PERCNT=.001
ACR=.001
DO 58 I=1,15

```

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```

RUN1030
RUN1040
RUN1060
RUN1080
RUN1190
RUN1230
RUN1260
RUN1240
RUN1250

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58 BODY(I)=0.0
   DO 60 I=1,3
   RNU(I)=0.0
   RMA(I)=0.0
   CHI(I)=0.0
   FLUXIN(I)=0.0
   PCOL(I)=0.0
   LL(I)=0
   DO 59 J=1,18
   ALPHA(I,J)=1.0
59 SIGMA(I,J)=1.0
   DO 60 K=1,3
   ETA(I,K)=0.0
   PHI(I,K)=0.0
60 DIR(I,K)=0.0
   WRITE(6,1)
   READ(5,CONTRL)
   WRITE(6,CONTRL)
   IF(NEW) GO TO 103
   REWIND 9
   READ(9) DENP,U,XREF,TRP,KAWLS,NL,NW,NH,BW,BH,NREG,XLB,XLC,PI,ND,
2      S,SINANG,COSANG,AKN,AKT,NBX,BM,XR,TIME,DTM,TI,ITS,ITP,TST,
3      TLM,RMA,RNU,DIR,XSTART,MNM,MNB,TR,BZC,CN7,DRP,PCF,PNA,
4      HTF,INM,LLM,NAV,NMAX,NWEDG,PRT,SAMP,AKN1,AKN2,AKT1,AKT2,
5      BTA,C1,C2,C3,C7,C8,DAM,DPA,FL,DELANG,PDN,HTI,HTR,JNT,KNM,
6      NM,WTM,C4,VRM,NCOL,CTI,CTR,CNI,CNR,SN,
7      ST,D1,D2,D3,D4,NRAN,VELR,RMP,RMN,RMF,IFLUX,FLUXIN,
8      XLIM,COEFP,XCB,XS,YCB,TB,ALPHA,SIGMA,NTS,
9      UTL,UTT,VTS,HTS,HTSI,ENT,REM,TMPA,
A      DBA,NB,NBT,TMP,XV,XVA,YV,YVA,ZV,ZVA,T,DB,FNB,XC,YC,ZC,
B      PAU,PAV,PAW,PAX,PAY,PAZ,LCOL,LM,
C      ETA,PHI,CHI,CN,CM,CNG,CMG,CN8,TRP,TRPA,MSP,ANGLE,TF,
D      UTLI,UTTI,VTSI,ER,RMB,LB,NBM,NBN,VEL,PFV,PCOL,JV
   DTMO=DTM
   READ(5,TIMES)
   WRITE(6,TIMES)
   IF(DTM.EQ.DTMO) GO TO 100
   AIME=TIME*DTMO
   TIME=AIME/DTM+0.1
   DO 99 J=1,NSP
   DO 99 L=1,2
   ENT(J,L)=ENT(J,L)*DTM/DTMO
99 CONTINUE
100 IF(TI.GT.0.0) TST=TI/DTM
   WRITE(6,2)
   WRITE(6,4)
   WRITE(6,5) (MT,RMA(MT),RNU(MT),PCOL(MT),FLUXIN(MT),(ENT(MT,K),K=1,2
1) ,MT=1,MSP)
   WRITE(6,2)
   CALL PRINTA(NWEDG,TITLE,NAME,XCB,YCB,TB,ALPHA,SIGMA,XLIM,
1COEFP,LIMIT,MSP)
   CALL PRINTB(PNA,MSP,FNB,NM,XLIM,XC,YC,ZC,NB,NSP)
   GO TO 280
103 READ(5,TIMES)
   WRITE(6,TIMES)

```

RUN1330

RUN1340

RUN1360

RUN1370

RUN1540

RUN1590

```

READ(5,FLOREP)
WRITE(6,FLOREP)
READ(5,MOLEC)
WRITE(6,MOLEC)
IF(MSP.GT.LIMIT(10))CALL DIAG(10,LIMIT(10),MSP)
CHIM=0.0
RMR=0.0
DMR=0.0
ETT=0.0
DO 105 M=1,MSP
  RMR=RMR+RMA(M)*RNU(M)
  CHIM=CHIM+CHI(M)*RNU(M)
105 CONTINUE
DO 115 K=1,MSP
DO 115 M=1,MSP
  ETT=ETT+RNU(M)*RNU(K)*ETA(M,K)
  SRMA=SQRT(.5*RMR*(1./RMA(M)+1./RMA(K)))
115 DMR=DMR+RNU(M)*RNU(K)*DIR(M,K)*(TRF/TF)**(ETA(M,K)/2.)*SRMA
  XREF=1./(DENF*DMR*1.414214)
  VELR=SQRT(16628.64*TF/RMR)
  IF(MET.NE.0)VELR=SQRT(99437.92*TF/RMR)
  TMR=XREF/VELR
  S=U/VELR
  NREG=1
  ND=1
  READ(5,SHAPES)
  WRITE(6,SHAPES)
  XCB(1)=BODY(1)/XREF
  TB(1)=BODY(2)/TF
  XLIM(1)=XCB(1)
  XSTART=XLIM(1)
  RMB=.5*BODY(3)/XREF
  TR=TB(1)
104 READ(5,SHAPES)
  WRITE(6,SHAPES)
  ND=ND+1
  IF(ND.GT.LIMIT(3))CALL DIAG(3,LIMIT(3),ND)
  XCB(ND)=BODY(1)/XREF
  TB(ND)=BODY(2)/TF
  DO 1104 M=1,MSP
  ALPHA(M,ND)=BODY(2+2*M)
  SIGMA(M,ND)=BODY(3+2*M)
1104 IF(TB(ND).GT.TR)TR=TB(ND)
  IF(BODY(3).EQ.0.0)GO TO 104
  NSTEP=NREG+1
  XLIM(NSTEP)=XCB(ND)
  COEFP(1)=0.0
  COEFP(2)=1.0
  COEFP(3)=0.0
  COEFP(4)=-RMB**2
  XR=XLIM(NSTEP)-XSTART
  AKN=1./XR
  AKT=.5/RMB
  RMFP1=TB(1)**(.5+ETT/2.)/(2.*PIROOT*S)
  RMFP2=TB(ND)**(.5+ETT/2.)/(2.*PIROOT*S)

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RUN1640
RUN1650

RUN1670

RUN1680
RUN1690

RUN1770

RUN1800

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AKN1=AKN*RMFP1
AKN2=AKN*RMFP2
AKT1=AKT*RMFP1
AKT2=AKT*RMFP2
DO 260 N=1,3
DO 260 M=1,3
NCOL(N,M)=0
260 CONTINUE
READ(5,GEOM)
READ(5,INCUPL)
WRITE(6,GEOM)
WRITE(6,INCUPL)
DO 116 J=1,42
VARG(J)=0.0
116 CURV(J)=0.0
RMN=RM*TB(1)/TB(ND)
RMP=SQRT(RMN*RMP)
RM=RMB
BW=XR/NW
BH=RM/NH
IF(NWEDG.GT.LIMIT(1)) CALL DIAG(1,LIMIT(1),NWEDG)
IF(MNM.GT.LIMIT(5)) CALL DIAG(5,LIMIT(5),MNM)
IF(MNB.GT.LIMIT(6)) CALL DIAG(6,LIMIT(6),MNB)
DELANG=180./NWEDG
SINANG=SIN(ANGLE/180.*PI)
COSANG=COS(ANGLE/180.*PI)
VOL=PI*RM*RM*XR
NBX=NW*NH*NWEDG
IF(NBX.GT.LIMIT(4)) CALL DIAG(4,LIMIT(4),NBX)
IF(JV.NE.LIMIT(7)) CALL DIAG(7,LIMIT(7),JV)
BR=SQRT(TR)
SRMX=0.0
DO 916 MT=1,MSP
WTH(MT)=RMA(MT)/RMR
BTA(MT)=SQRT(WTH(MT))
SR(MT)=S*BTA(MT)
SRT=SR(MT)*FLUXIN(MT)/RNU(MT)
IF(SRT.GT.SRMX) SRMX=SRT
916 CONTINUE
INM=LLM*SQRT(TB(1))/PIROOT/SRMX/(1.+RMN)
DDN=INM/VOL
DO 140 MT=1,MSP
FDN(MT)=RNU(MT)*DDN
DPA(MT)=RNU(MT)
SN(MT)=SR(MT)*COSANG
ST(MT)=SR(MT)*SINANG
DO 117 K=1,MSP
DAM(K,MT)=DIR(K,MT)*(TRF/TF)**(ETA(K,MT)/2.)/DMR
CN8(K,MT)=DDN/DAM(K,MT)*1.414214
BT=AMIN1(BTA(K),BTA(MT))
VR1=S+3.*(1.+SQRT(TR))/BT
VR2=3.*SQRT((1.+2.*S**2/(5.+CHIM))*(1./WTH(K)+1./WTH(MT)))
CM(K,MT,1)=AMAX1(VR1,VR2)
CN(K,MT,1)=RAND(0)*CM(K,MT,1)
DP=PHI(K,MT)*(CHI(K)+CHI(MT)+2.)-1

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ORIGINAL PAGE 19
OF POOR QUALITY

RUN2200

RUN2230

RUN2240

RUN2320

RUN2330

RUN2420

RUN2470

RUN2480

RUN2740

RUN2750

```

DS=PHI(K,MT)*(2.-.5*ETA(K,MT))-1.0
DO 917 N=2,3
XPM=ACR**AMIN1(DP,DS)
IF((DP.GT.0.).AND.(DS.GT.0.))XPM=(DP/(DP+DS))*DP*(DS/(DP+DS))*DS
XPN=ACR**AMAX1(DP,DS)
IF((DP.LT.0.).AND.(DS.LT.0.))XPN=(DP/(DP+DS))*DP*(DS/(DP+DS))*DS
CM(K,MT,N)=XPM-XPN
CN(K,MT,N)=RAND(0)*CM(K,MT,N)
DF=CHI(K)
DS=CHI(MT)
917 CONTINUE
117 CONTINUE
DO 118 K=1,KMX
READ(5,INOUT)
WRITE(6,INOUT)
DO 118 J=1,JV
VEL(J,K,MT)=VARG(J)
PFV(J,K,MT)=CURV(J)
118 CONTINUE
ENT(MT,1)=INM*S*DTM*AKN*FLUXIN(MT)/RNU(MT)
ENT(MT,2)=RMN*ENT(MT,1)*SQRT(TB(ND)/TB(1))
REM(MT,1)=0.0
REM(MT,2)=0.0
LL(MT)=INM*PIROOT*(1.+RMN)*SR(MT)/SQRT(TB(1))*FLUXIN(MT)/RNU(MT)
CHT=CHI(MT)
IF(CHT.GT.0.) CMG(MT)=CHT**CHT*EXP(-CHT)
IF(CHT.EQ.0.) CMG(MT)=1.0
IF(CHT.LT.0.) CMG(MT)=ACR**CHT*EXP(-ACR)
CNG(MT)=RAND(0)*CMG(MT)
140 CONTINUE
XS(1)=0.0
DO 155 N=2,ND
155 XS(N)=(.5*(XCB(N)+XCB(N-1))-XSTART)*AKN
YCB(1)=0.0
DO 160 N=2,ND
160 YCB(N)=2.*(XCB(N)-XCB(N-1))/RMB
CALL CELL(BW,BH,NW,NH,XSTART,DELANG,NWEDG,IC,YC,ZC,FNB)
FNA=0.0
DO 210 N=1,NBX
210 FNA=FNA+FNB(N)
NPX=NBX
220 TIME=0
LARGE=0
SAMP=0
PRT=0
NAV=0
AIME=0.
TI=0.0
DT=DTM
NMAX=0
DO 250 MT=1,3
C1(MT)=RAND(0)
C2(MT)=RAND(0)
C3(MT)=RAND(0)
C7(MT)=RAND(0)

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OF POOR QUALITY

RUN3330

RUN3380

RUN3390

RUN4250

RUN4260

RUN4270

RUN4290

RUN4300

RUN4310

RUN4380

RUN4410

RUN4420

RUN4430

RUN4440

C8(MT)=RAND(0) RUN4450
 D1(MT)=RAND(0) RUN4460
 D2(MT)=RAND(0) RUN4470
 D3(MT)=RAND(0) RUN4480
 D4(MT)=RAND(0) RUN4490
 FL(MT)=0. RUN4500
 HTI(MT)=0. RUN4510
 HTR(MT)=0. RUN4520
 JNT(MT)=0 RUN4530
 NM(MT)=0 RUN4540
 IFLUX(MT,1)=0
 IFLUX(MT,2)=0
 DO 230 N=1,3 RUN4550
 CTI(MT,N)=0. RUN4560
 CTR(MT,N)=0. RUN4570
 CNI(MT,N)=0. RUN4580
 230 CNR(MT,N)=0. RUN4590
 DO 240 N=1,ND RUN4600
 DO 240 K=1,NWEDG RUN4610
 NTS(MT,N,K)=0 RUN4620
 HTSI(MT,N,K)=0. RUN4640
 UTLI(MT,N,K)=0.
 UTTI(MT,N,K)=0.
 VTSI(MT,N,K)=0.
 UTL(MT,N,K)=0. RUN4650
 UTT(MT,N,K)=0. RUN4660
 VTS(MT,N,K)=0. RUN4670
 240 HTS(MT,N,K)=0. RUN4680
 DO 245 N=1,NPX RUN4690
 NB(MT,N)=0 RUN4700
 NBT(MT,N)=0 RUN4720
 DBA(MT,N)=0. RUN4730
 XVA(MT,N)=0. RUN4740
 YVA(MT,N)=0. RUN4750
 ZVA(MT,N)=0. RUN4760
 TMPA(MT,N)=0. RUN4790
 TRPA(MT,N)=0.0
 DO 245 NN=1,NSP
 T(MT,NN,N)=0.0
 245 CONTINUE RUN4800
 250 CONTINUE
 FND=DDN
 DRF=2./ (FND*S*S*RMB*RMB*PI) RUN4930
 PCF=1./ (FND*S*RMB*RMB*PI) RUN4940
 HTP=.5*DRF/S
 WRITE(6,2)
 WRITE(6,4)
 WRITE(6,5) (MT,RMA(MT),RNU(MT),PCOL(MT),FLUXIN(MT), (ENT(MT,K),K=1,2
 1),MT=1,MSP)
 WRITE(6,2) RUN5000
 CALL PRINTA(NWEDG,TITLE,NAME,XCB,YCB,TB,ALPHA,SIGMA,XLIM,
 1COEFP,LIMIT,MSP)
 CALL GAS(NWEDG,DELANG,ND,BTA,C1,DFA,NM,FNB,DB,NB,NBM,NBN,
 1PAU,PAV,PAW,PAX,PAY,PAZ,XLIM,COEFP,LM,LIMIT(4),LIMIT(6),XCB,TB,
 2LARGE,MNN,MNB,DEBUG(1),LCOL,MSP,ER,CHI,CNG,CMG,NSP,LB)

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RUN5060

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CPUTYM=TFIND(0)
IF (LARGE.NE.0) GO TO 345
DO 265 I=1,MSP
265 IF (NM(I).GT.NMAX) NMAX=NM(I)
CALL PRINTB(FNA,MSP,FNB,NM,XLIM,XC,YC,ZC,NB,NSP)
IF (DEBUG(2)) WRITE(6,1)
CALL ACCUM(NMC,NPB,FNB,NB,PAU,PAV,PAW,ER,TMP,TRP,XV,YV,ZV,LM,MSP,
1NSP,NBM)
CPA=ELTIME(0)
CPI=CPA
GO TO 340
280 TIME=TIME+1
IF (TIME.NE.TST+1) GO TO 285
TI=TST*DTM
DO 282 MT=1,MSP
282 IFLUX(MT,1)=0
285 LARGE=0
CPI=ELTIME(0)
AIME=TIME*DTM
DT=AIME-TI
IF (DEBUG(1)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA,(NM(I),I=1,3),NMAX
PRT=PRT+1
SAMP=SAMP+1
CALL COLIDE(CN,CM,WTM,DB,DBA,NB,NCOL,LCOL,PAU,PAV,PAW,ER,T,LM,MSP,
1LIMIT(4),LIMIT(6),ETA,PHI,CHI,CN8,NSP,NBM)
KNM(1)=0
KNM(2)=0
KNM(3)=0
CPC=ELTIME(0)
IF (DEBUG(1)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA,(NM(I),I=1,3),NMAX
CALL MOVE(0,AKN,NWEDG,XSTART,LIMIT(3),LIMIT(1),LIMIT(8),LIMIT(9),
1DELANG,BTA,C2,C3,DPA,FL,HTI,HTR,JNT,KNM,NM,XCB,XLIM,CTI,CTR,
2CNI,CNR,ALPHA,SIGMA,COEPP,HTS,HTSI,NTS,UTL,UTT,VTS,PAU,PAV,PAW,
3PAX,PAY,PAZ,LCOL,TB,MSP,ER,CHI,CNG,CMG,NSP,UTLI,UTTI,VTSI,IFLUX)
KNM(1)=NM(1)
KNM(2)=NM(2)
KNM(3)=NM(3)
CPM=ELTIME(0)
IF (DEBUG(1)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA,(NM(I),I=1,3),NMAX
TBI=TB(ND)
CALL FLOW(NWEDG,MNM,LARGE,BTA,C1,C7,C8,ENT,REM,LCOL,MSP,NM,SN,ST,
1TBI,PAU,PAV,PAW,PAX,PAY,PAZ,ER,CHI,CNG,CMG,NSP,JV,FCOL,VEL,PFV)
IF (LARGE.NE.0) GO TO 345
CPB=ELTIME(0)
IF (DEBUG(1)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA,(NM(I),I=1,3),NMAX
CALL MOVE(1,AKN,NWEDG,XSTART,LIMIT(3),LIMIT(1),LIMIT(8),LIMIT(9),
1DELANG,BTA,C2,C3,DPA,FL,HTI,HTR,JNT,KNM,NM,XCB,XLIM,CTI,CTR,
2CNI,CNR,ALPHA,SIGMA,COEPP,HTS,HTSI,NTS,UTL,UTT,VTS,PAU,PAV,PAW,
3PAX,PAY,PAZ,LCOL,TB,MSP,ER,CHI,CNG,CMG,NSP,UTLI,UTTI,VTSI,IFLUX)
CPB=CPB+ELTIME(0)
IF (DEBUG(1)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA,(NM(I),I=1,3),NMAX
DO 330 MT=1,MSP
DO 290 N=1,NBX
NB(MT,N)=0

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RUN5130

RUN5140

RUN5170

RUN5180

RUN5190

RUN5220

RUN5230

RUN5280

RUN5290

RUN5330

RUN5420

RUN5450

290	CONTINUE	RUN5470
	NG=NM(MT)	RUN5480
	N=0	RUN5490
295	N=N+1	RUN5500
	IF(N.GT.NG) GO TO 310	RUN5510
	X=PAY(MT,N)	RUN5520
	Y=PAY(MT,N)	RUN5530
	Z=PAZ(MT,N)	RUN5540
	R=SQRT(Y*Y+Z*Z)	RUN5550
	TANG=180.*ATAN2(Z,-Y)/PI	
	IWDGE=TANG/DELANG+1	
	IF(IWDGE.LT.1) IWDGE=1	
	IF(IWDGE.GT.NWEDG) IWDGE=NWEDG	
	L=X/BW	RUN5610
	IF(L.GE.NW) L=NW-1	RUN5620
	M=R/BH	RUN5630
	IF(M.GE.NH) M=NH-1	RUN5640
	K=NWEDG*(L*NH+M)+IWDGE	
	IF(K.LE.NBX) GO TO 305	RUN5660
	WRITE(6,40)L,M,IWDGE,NW,NH,NWEDG,K,MT,N,X,Y,Z,R,TANG	
	IF(DUMP) CALL ABEND(4)	RUN5680
	STOP	RUN5690
305	J=NB(MT,K)+1	
	IF(J.LE.MNB) GO TO 308	
	IF(DEBUG(1)) WRITE(6,44) MT,K,MNB,AIME	
308	NB(MT,K)=J	RUN6400
	LB(N)=K	
	GO TO 295	RUN6430
310	CONTINUE	
	NBM(MT,1)=0	RUN6450
	DO 320 M=1,NBX	RUN6480
	A=NB(MT,M)	RUN6490
	DB(MT,M)=A*DFA(MT)/FNB(M)	
	NBM(MT,M+1)=NBM(MT,M)+NB(MT,M)	
	NBN(M)=NBM(MT,M)	
320	CONTINUE	RUN6500
	IF(NM(MT).GT.NMAX) NMAX=NM(MT)	RUN6510
	DO 325 N=1,NG	
	Q=LB(N)	
	NBN(Q)=NBN(Q)+1	
	NA=NBN(Q)	
325	LM(MT,NA)=N	
330	CONTINUE	RUN6520
	IF(SAMP.LT.ITS) GO TO 335	
	CALL ACCUM(NMC,NPB,FNB,NB,PAU,PAV,PAW,ER,TMP,TRP,XV,YV,ZV,LM,MSP,	
	1NSP,NBM)	
	SAMP=0	RUN6570
	IF(TIME.LE.TST) GO TO 335	
	CALL AVRGE(FNB,DB,DBA,NB,NBT,XV,YV,ZV,XVA,YVA,ZVA,TMP,TMPA,TRP,TRP	
	1A,MSP,NSP)	
	NAV=NAV+1	RUN6600
335	CPA=ELTIME(0)	
	CPI=CPC+CPH+CPB+CPA	
	CPJ=2.*CPI+5.	
340	CPUTYM=TFIND(0)	

ORIGINAL PAGE 131
OF POOR QUALITY

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IF (DEBUG (3)) WRITE (6, 33) AIME, CPC, CPM, CPI, CPB, CPA, (NM (I), I=1, 3), NMAX
IF (TIME.EQ.0) GO TO 355
IF ((TIME.GE.TLIM).OR. (CPUTYM.LE.CPJ)) GO TO 345
IF (PRT.LT.ITP) GO TO 280
PRT=0
345 WRITE (6, 30) AIME, KAWLS
IF (DEBUG (3)) WRITE (6, 31) CPUTYM
WRITE (6, 32) (NM (I), I=1, 3)
WRITE (6, 34) ((NCOL (I, J), J=1, 3), I=1, 3)
WRITE (6, 35) (JNT (I), I=1, 3)
IF (LARGE.NE.0) GO TO 360
WRITE (6, 36) NMAX
IF (.NOT.SAVE) GO TO 355
IF (PRT.NE.0.AND.CPUTYM.GT.CPJ.AND.TIME.LT.TLIM) GO TO 355
REWIND 9
WRITE (9) DENF, U, XREF, TRF, KAWLS, NL, NW, NH, BW, BH, NREG, XLB, XLC, PI, ND,
2 S, SINANG, COSANG, AKN, AKT, NBX, RM, XR, TIME, DTM, TI, ITS, ITP, TST,
3 TLIM, RMA, RNU, DIR, XSTART, MNM, MNB, TR, BZC, CN7, DRF, PCF, PNA,
4 HTP, INM, LLM, NAV, NMAX, NWEDG, PRT, SAMP, AKN1, AKN2, AKT1, AKT2,
5 BTA, C1, C2, C3, C7, C8, DAM, DFA, FL, DELANG, PDN, HTI, HTR, JNT, KNM,
6 NM, WTM, C4, VRM, NCOL, CTI, CTR, CNI, CNR, SN,
7 ST, D1, D2, D3, D4, NRAN, VELR, RMP, RMN, RMP, IFLUX, FLUXIN,
8 XLIM, COEPP, XCB, XS, YCB, TB, ALPHA, SIGMA, NTS,
9 UTL, UTT, VTS, HTS, HTSI, ENT, REM, TMPA,
A DBA, NB, NBT, TMP, XV, XVA, YV, YVA, ZV, ZVA, T, DB, FNB, XC, YC, ZC,
B PAU, PAV, PAW, PAX, PAY, PAZ, LCOL, LM,
C ETA, PHI, CHI, CN, CM, CNG, CMG, CN8, TRP, TRPA, MSP, ANGLE, TP,
D UTLI, UTTI, VTSI, ER, RMB, LB, NEM, NBN, VEL, PFV, FCOL, JV
WRITE (6, 50)
355 CONTINUE
WRITE (6, 25) RMP, RMN, RMP
DO 356 MT=1, MSP
FIN1=ENT (MT, 1)
FIN2=ENT (MT, 2)
RF1=IFLUX (MT, 1) *DTM/DT
RF2=IFLUX (MT, 2) *DTM/DT
RNF1=1.-RF1/FIN1
RNF2=(FIN2-RF2)/FIN1
RPS1=RNF1*FLUXIN (MT) /RNU (MT)
RPS2=RNF2*FLUXIN (MT) /RNU (MT)
356 WRITE (6, 26) RMA (MT), FIN1, RF1, RNF1, RPS1, FIN2, RF2, RNF2, RPS2
IF (TIME.EQ.0) GO TO 280
IF (TIME.LE.TST) GO TO 350
CALL PRINT1 (DT, COSANG, SINANG, RMA, RNU, DRF, PCF, HTP, FL, HTI, HTR, CTI,
1 CTR, CNI, CNR)
CALL PRINT2 (AKN, XSTART, DT, RNU, RMA, DRF, PCF, HTP, UTLI, UTTI, VTSI, HTSI,
1 DELANG, NWEDG, XS, XCB, YCB, HTS, NTS, UTL, UTT, VTS, LIMIT (3), LIMIT (1), MSP)
CALL PRINT4 (MSP, CHI, RNU, NSP, TRPA, PDN, WTM, DBA, NBT, TMPA, XVA,
1 YVA, ZVA, 1, NBT, XC, YC, ZC)
GO TO 353
350 CONTINUE
CALL PRINT4 (MSP, CHI, RNU, NSP, TRP, PDN, WTM, DB, NB, TMP, XV, YV, ZV,
10, NB, XC, YC, ZC)
353 IF (DEBUG (2)) WRITE (6, 1)
IF ((TIME.LT.TLIM).AND. (CPUTYM.GT.CPJ)) GO TO 280

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RUN6650
 RUN6660
 RUN6670

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RUN6740
 RUN6750
 RUN6900

RUN7050

RUN6860

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IP(IC.EQ.ICOPY) RETURN
IC=IC+1
WRITE(6,2)
WRITE(6,4)
WRITE(6,5) (MT,RMA(MT),RNU(MT),PCOL(MT),FLUXIN(MT), (ENT(MT,K),K=1,2
1),MT=1,MSP)
WRITE(6,2)
WRITE(6,3) IC
CALL PRINTA(NWEDG,TITLE,NAME,XCB,YCB,TB,ALPHA,SIGMA,XLIM,
1COEFF,LIMIT,MSP)
CALL PRINTB(FNA,MSP,FNB,NM,XLIM,XC,YC,ZC,NB,NSP)
SAVE=.FALSE.
GO TO 345
360 WRITE(6,38) (DBG1(I,LARGE),I=1,3)
IF((REDO).AND.(TIME.LE.TST)) GO TO 364
IF(DUMP) CALL ABEND(9)
STOP
364 CONTINUE
IF(NEW) GO TO 365
REWIND 9
READ(9) DENP,U,XREF,TRP,KAWLS,NL,NW,NH,BW,BH,NREG,XLB,XLC,PI,ND,
2 S,SINANG,COSANG,AKN,AKT,NBX,RM,XR,TIME,DTM,TI,ITS,ITP,TST,
3 TLIM,RMA,RNU,DIR,XSTART,MNM,MNB,TR,BZC,CN7,DRP,PCP,PNA,
4 HTP,INM,LLM,NAV,NMAX,NWEDG,PRT,SAMP,AKN1,AKN2,AKT1,AKT2,
5 BTA,C1,C2,C3,C7,C8,DAM,DFA,PL,DELANG,PDN,HTI,HTR,JNT,KNM,
6 NM,WTM,C4,VRM,NCOL,CTI,CTR,CNI,CNR,SN,
7 ST,D1,D2,D3,D4,NRAN,VELR,RMP,RMN,RMP,IPLUX,FLUXIN,
8 XLIM,COEFF,XCB,XS,YCB,TB,ALPHA,SIGMA,NTS,
9 UTL,UTT,VTS,HTS,HTSI,ENT,REM,TMPA,
A DBA,NB,NBT,TMP,XV,XVA,YV,YVA,ZV,ZVA,T,DB,FNB,XC,YC,ZC,
B PAU,PAV,PAW,PAX,PAY,PAZ,LCOL,LM,
C ETA,PHI,CHI,CN,CM,CNG,CNG,CN8,TRP,TRPA,MSP,ANGLE,TP,
D UTLI,UTTI,VTSI,ER,RMB,LB,NBM,NBN,VEL,PFV,PCOL,JV
365 ANM=INM
INM=9*ANM/10
DDN=.9*DDN
DRP=DRP/.9
PCP=PCP/.9
HTP=HTP/.9
DO 370 MM=1,MSP
PDN(MM)=PDN(MM)*INM/ANM
LL(MM)=9*LL(MM)/10
DO 366 KK=1,MSP
366 CN8(KK,MM)=CN8(KK,MM)*.9
DO 370 NK=1,2
ENT(MM,NK)=ENT(MM,NK)*INM/ANM
370 REM(MM,NK)=0.0
IF(NEW) GO TO 220
TST=TIME+TST
TI=0.
PRT=ITP
WRITE(6,2)
WRITE(6,4)
WRITE(6,5) (MT,RMA(MT),RNU(MT),PCOL(MT),FLUXIN(MT), (ENT(MT,K),K=1,2
1),MT=1,MSP)

```

RUN7080

RUN7120

RUN7130

RUN7180

RUN7190

RUN7200

RUN7220

RUN7230

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OF POOR QUALITY

RUN7260

```

WRITE(6,2)
IF((LARGE.EQ.2).OR.(LARGE.EQ.3)) GO TO 280
REDO=.FALSE.
GO TO 360
END

```

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RUN7410

```

SUBROUTINE DIAG(N, ITEST, NUM)
REAL*8 PARAM(10) /' NWEDGE', ' NREG', ' ND', ' NPX', '
1 MNM', ' MNB', ' NBX', ' NS', ' MJ', ' MSP' /

```

DIAG010

DIAG040

DIAG050

DIAG060

DIAG070

DIAG080

DIAG090

DIAG100

FORMATS

```

32 FORMAT(9X, 'ENT, REM, ENTS, REMS, FTH, THETA, DTH')
42 FORMAT(///5X, 43H ARRAY DIMENSIONS ARE ABOUT TO BE VIOLATED./)
44 FORMAT(5X, 17H MAXIMUM VALUE IS, 15, 19H, WHEREAS YOU INPUT, 15, 3H (,
1A8, 1H))
56 FORMAT(/5X, 78H IF YOU DESIRE TO USE THIS VALUE, THE FOLLOWING ARRADIAG130
1YS MUST BE RE-DIMENSIONED./)
62 FORMAT(9X, 'HTS, HTSI, NTS, NTSP, UTL, UTT, VTS')
64 FORMAT(9X, 'XLIM, COEPP' // 11X, 'NOTE THAT THE XLIM ARRAY MUST BE DIMEDIAG160
1NSIONED TO 3 MORE THAN THE COEPP ARRAY.')
66 FORMAT(9X, 'XCB, XS, YCB, TB, ALPHA, SIGMA')
68 FORMAT(9X, 'DBA, NB, NBT, TMP, TMPA, XV, XVA, YV, YVA, ZV, ZVA, T, DB')
70 FORMAT(9X, 'PAU, PAV, PAW, PAX, PAY, PAZ, LCOL')
72 FORMAT(8X, 3H LM)
74 FORMAT(///5X, 76H IF YOU CHANGE THE ARRAY DIMENSIONS, ALSO CHANGE THDIAG220
1E 'LIMIT' DATA STATEMENT.)
75 FORMAT(9X, 'ALL ARRAYS ASSOCIATED WITH SPECIES')
76 FORMAT(9X, 'PNB, XC, YC, ZC')
78 FORMAT(9X, 'FV, NTCV, NTCF, MS, IWS, SL, DELS, TANGN')
80 FORMAT(9X, 'VEL, PFV')

```

DIAG120

DIAG130

DIAG140

DIAG150

DIAG160

DIAG170

DIAG180

DIAG190

DIAG210

DIAG220

DIAG230

DIAG240

DIAG250

DIAG260

DIAG270

DIAG280

DIAG290

DIAG300

DIAG310

DIAG320

DIAG340

DIAG350

DIAG370

DIAG390

DIAG400

DIAG420

DIAG440

DIAG460

```

WRITE(6,42)
WRITE(6,44) ITEST, NUM, PARAM(N)
WRITE(6,56)
GO TO (1,2,3,4,5,6,7,8,9,10), N
1 WRITE(6,62)
WRITE(6,32)
GO TO 11
2 WRITE(6,64)
GO TO 11
3 WRITE(6,66)
WRITE(6,62)
GO TO 11
4 WRITE(6,68)
WRITE(6,76)
GO TO 11
5 WRITE(6,70)
GO TO 11
6 WRITE(6,72)
GO TO 11

```


7 WRITE(6,80)		DIAG480
GO TO 11		
8 WRITE(6,78)		DIAG500
GO TO 11		
9 WRITE(6,80)	ORIGINAL PAGE IS	DIAG520
GO TO 11	OF POOR QUALITY	
10 WRITE(6,75)		
11 WRITE(6,74)		
STOP		DIAG540
END		DIAG550

```

SUBROUTINE PRINTA(NWEDG,TITLE,NAME,XCB,YCB,TB,ALPHA,SIGMA,XLIM,
1COEFF,LIMIT,MSP)
INTEGER TST,TLIM,TIME
LOGICAL SAVE,NEW
DIMENSION LIMIT(1),TITLE(6),NAME(2),XCB(1),YCB(1),TB(1)
DIMENSION ALPHA(3,1),SIGMA(3,1),XLIM(1),COEFF(4)
DIMENSION RNU(3),RMA(3),CHI(3),DIR(3,3),PHI(3,3),ETA(3,3)
DIMENSION WTM(3),DAM(3,3),VELS(3),XSP(3)
COMMON /FIRST/NL,NW,NH
COMMON /SECND/BW,BH,RMP,RMN,RMF
COMMON /THIRD/PI,NREG,S,SINANG,COSANG,AKN,AKT,AKN1,AKN2,AKT1,AKT2
COMMON /FOURTH/NBX,RM,XR,DUMP,C9,LL(3),LLM
COMMON /FIFTH/ND,TIME,DTM,TI,ITS,ITP,TST,TLIM,RMA,RNU,DIR
COMMON /SIXTH/RMB,XSTART,INM,MNM,MNB,NEW,SAVE,PERCNT,NSR,TR
COMMON/EIGHT/DENP,U,TP,ANGLE,TRF,CHI,PHI,ETA,WTM,DAM,VELR,XREF
DATA NOT/'NOT '/

```

PRA0030

PRA0040

PRA0100

PRA0120

PRA0130

PRA0140

PRA0150

PRA0160

PRA0170

FORMATS

1 FORMAT(16X,40(' '),T74,'I'//9X,'3-D',I2,'-FLUID PROGRAM - ')	
2 FORMAT('+',31X,A4)	PRA0190
3 FORMAT('+',35X, 'A RESTART OF A PREVIOUS RUN',T74,'I'/12X,2	PRA0200
1A4,' - ',6A4,' - ',I2,' REGIONS',T74,'I',16 (/T74,'I'))	PRA0210
4 FORMAT(7X,'FRONT OF BODY =' ,E12.4,' XSTART MAX HEIGHT =' ,E12.4,'	
1RMB',T74,'I'/7X,'BODY TEMPERATURES =' ,F12.2,' T PR.STRM.',F12.2,	
2' T ENTR.',F12.2,' T CAVITY'/7X,'X-LIMIT',T37,'BODY COEFFICIENTS',	
3T74,'I')	
6 FORMAT(5F14.6,3X,'I')	PRA0240
10 FORMAT(1X,72(' '))	PRA0250
12 FORMAT(/14X,'PARAMETERS OF SEGMENTS FOR BODY COLLISIONS',T96,'I'/	
18X,'X-COORD. TEMP. ALPHA1 ALPHA2 ALPHA3 SIGMA1 SIGMA2	
2 SIGMA3 AREAS',T96,'I')	
14 FORMAT(4X,E12.4,7F9.4,E12.4,T96,'I')	
17 FORMAT(///25X,'ARRAY STORAGE USED'/5X,I6,' *',10I6,T96,'I')	
18 FORMAT(1H1/17X,'LENGTH OF CELL IN MEAN-FREE-PATHS = ',F12.4,' BW'	
A,T76,'I'	
1/17X,'HEIGHT OF CELL IN MEAN-FREE-PATHS = ',F12.4,' BH',T76,'I'	
2/16X,'NUMBER OF L1 CELLS ALONG FLOW AXIS = ',I13,' NW',T76,'I'	
3/17X,'NUMBER OF L1 CELLS IN RADIAL DIR. = ',I13,' NH',T76,'I'	
4/21X,'NUMBER OF LEVELS OF CELL SIZE = ',I13,' NL',T76,'I')	
23 FORMAT(3X,'NUMBER OF AZIMUTHAL WEDGES WITHIN ',I3,' DEGREES =' ,IPRA0470	
113,' NWEDG I'/)	

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24 FORMAT(16X,'BASIC TIME INTERVAL FOR COLLISIONS =' ,E13.4,' DTM
1 I'/8X,'TIME INTERVAL FOR SAMPLING FLOW FIELD INFO =' ,E13.4,' DTS
2 I'/24X,'TIME INTERVAL FOR PRINTING =' ,E13.4,' DTP I'/9X,
3'TIME TO STEADY-STATE CONDITIONS (ASSUMED) =' ,E13.4,' TST I'/
419X,'TIME AT WHICH RUN IS TERMINATED =' ,E13.4,' TLIM I'/)
25 FORMAT(/5X,'PRESSURE RATIO(INSIDE/ENTRANCE) - EITHER TYPE =' ,F13.5
1,' RMP',T76,'I'/6X,'DENSITY RATIO(INSIDE/ENTRANCE) - EITHER TYPE
2=' ,F13.5,' RMN',T76,'I'/9X,'FLUX RATIO(INSIDE/ENTRANCE) - EITHER
3TYPE =' ,F13.5,' RMP',T76,'I')
26 FORMAT(5X,'FREE STREAM NUMBER OF MOLECULES - EITHER TYPE =' ,I13,'
A INM I'/9X,'INITIAL NUMBER OF MOLECULES - MAXIMUM =' ,I13,
1' LLM I'/9X,'MAXIMUM NUMBER OF MOLECULES - EITHER TYPE =' ,I13
2,' MNM I'/1X,'MAX NUMBER OF MOLECULES IN ANY CELL - EITHER TY
3PE =' ,I13,' MNB I')
27 FORMAT( /22X,'VELOCITY OF FREE STREAM FLOW =' ,E13.4,' U',T76,'I'/1
19X,'SPEED RATIO OF FREE STREAM FLOW =' ,E13.4,' S',T76,'I'/19X,'MAC
AH NUMBER OF FREE STREAM FLOW =' ,E13.4,' M',T76,'I'/19X,'SPECIFIC H
BEAT RATIO (CALCULATED) =' ,E13.4,' GAMMA',T76,'I'/ 35X,'ANG
2LE OF ATTACK =' ,F13.4,' ANGLE I'/16X,'NUMBER DENSITY OF FREE ST
3REAM FLOW =' ,E13.4,' N',T76,'I'/19X,'TEMPERATURE OF FREE STREAM FL
4OW =' ,F13.4,' TP',T76,'I'/16X,'MOLE FRACTIONS OF FREE STREAM FLOW
5=' ,3E13.4,' RNU I'/16X,'MOLECULAR WEIGHTS OF SPECIES ABOVE =' ,3F13
6.4,' RMA I'/18X,'INITIAL NUMBERS OF SPECIES ABOVE =' ,3I13,' LL I'
7)
28 FORMAT( /10X,'REFERENCE TEMPERATURE FOR MOLECULAR DATA =' ,F13.4,'
1TRF',T90,'I'/14X,'CROSS-SECTION',26X,'TEMP EXPONENT',T90,'I'/3(3X,
23E12.4,3X,3F12.6,T90,'I'/) / 5X,'CHI/2-1',11X,'ROTATIONAL PARAMETER
3 PHI',T90,'I'/3(F12.4,5X,3F12.6,T90,'I'/))
29 FORMAT(/9X,'DATA SAVED ON TAPE 9')
30 FORMAT( 31X,'REF MOLECULAR SPEED =' ,E13.4,' VELR',T76,'I'/20X,'SP
1ECIES FREE STREAM MOLECULAR SPEEDS',T76,'I'/14X,3E16.6,T76,'I'/26X
2,'REFERENCE MEAN FREE PATH =' ,E13.4,' XREP',T76,'I'/26X,'SPECIES M
3EAN FREE PATHS',T76,'I'/14X,3E16.6,T76,'I'/11X,'LONGITUDINAL KNUDS
4EN NUMBER (FREE STRM.) =' ,E13.4,' AKN',T76,'I'/13X,'TRANSVERSE KNUD
5SEN NUMBER (FREE STRM.) =' ,E13.4,' AKT',T76,'I'/11X,'LONGITUDINAL K
6NUDSEN NUMBER ( ENTRANCE ) =' ,E13.4,' AKN1',T76,'I'/13X,'TRANSVERSE
7 KNUDSEN NUMBER ( ENTRANCE ) =' ,E13.4,' AKT1',T76,'I'/11X,'LONGITUD
8INAL KNUDSEN NUMBER ( CAVITY ) =' ,E13.4,' AKN2',T76,'I'/13X,'TRAN
9SVERSE KNUDSEN NUMBER ( CAVITY ) =' ,E13.4,' AKT2',T76,'I')

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PRA0680

PRA0690

PRA0700

PRA0710

PRA0720

PRA0730

PRA0750

PRA0760

PRA0770

PRA0780

PRA0800

PRA0810

PRA0820

```

IARRAY=708+LIMIT(3)*(32+56*LIMIT(1))+20*LIMIT(2)+LIMIT(4)*(120+4*LPRA0710
1IMIT(6))+56*LIMIT(5)+LIMIT(8)*(68+96*LIMIT(9))+20*LIMIT(7)+224*LIMPRA0720
2IT(1) PRA0730
WRITE(6,1) MSP
IF(NEW) WRITE(6,2) NOT
WRITE(6,3) NAME,TITLE,NREG
WRITE(6,4) XSTART,RMB
DO 100 I=1,NREG
100 WRITE(6,6) XLIM(I+1), (COEFP(J),J=1,4)
WRITE(6,10)
WRITE(6,12)
DO 110 I=1,ND
110 WRITE(6,14) XCB(I),TB(I), (ALPHA(J,I),J=1,3), (SIGMA(J,I),J=1,3),YCB

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1 (I)
  WRITE(6,10)
  WRITE(6,17) IARRAY, (LIMIT(I), I=1,10)
  WRITE(6,18) BW,BH,NW,NH,NL
  IETAZ=180.
  WRITE(6,23) IETAZ,NWEDG
  DTS=DTM*ITS
  DTP=DTM*ITP
  AST=DTM*TST
  ALIM=DTM*TLIM
  CHT=0.0
  DO 120 J=1,MSP
120 CHT=CHT+CHI(J)*RNU(J)
  GAMMA=(7.+2.*CHT)/(5.+2.*CHT)
  AM=S*SQRT(2./GAMMA)
  WRITE(6,24) DTM,DTS,DTP,AST,ALIM
  WRITE(6,26) INM,LLM,MNM,MNB
  WRITE(6,25) RMP,RMN,RMF
  WRITE(6,27) U,S,AM,GAMMA,ANGLE,DENF,TF,(RNU(I),I=1,3),(RMA(I),I=1,3)
  1), (LL(I),I=1,3)
  WRITE(6,28) TRF,((DIR(I,K),K=1,3),(ETA(I,K),K=1,3),I=1,3),(CHI(I),
  1 (PHI(I,K),K=1,3),I=1,3)
  DO 210 I=1,3
  VELS(I)=0.0
210 XSP(I)=0.0
  DO 220 J=1,MSP
  VELS(J)=VELR/SQRT(WTM(J))
  XT=0.0
  DO 215 M=1,MSP
215 XT=XT+RNU(M)*DAM(J,M)*SQRT(1.+WTM(J)/WTM(M))
220 XSP(J)=1.414214*XREF/XT
  WRITE(6,30) VELR,(VELS(I),I=1,3),XREF,(XSP(I),I=1,3),AKN,AKT,
  1AKN1,AKT1,AKN2,AKT2
  IF(SAVE) WRITE(6,29)
  RETURN
  END

SUBROUTINE PRINTB(FNA,MSP,FNB,NM,XLIM,XC,YC,ZC,NB,N)
  INTEGER*2 NB
  DIMENSION FNB(1),NM(1),XLIM(1),XC(1)
  DIMENSION YC(1),ZC(1),NB(N,1)
  COMMON /FIRST/NL,NW,NH
  COMMON /THIRD/PI,NREG,S,SINANG,COSANG,AKN,AKT
  COMMON /PORTH/NBX
  1 FORMAT(1H1)
  2 FORMAT(2X,'-----CELL GEOMETRY-----')
  1-----'/2X,'BOX LEVEL POSITION OF CENTER VOLUME
  2INITIAL POPULATION'/2X,'NUM.',12X,'X',7X,'Y THETA',12X,'TOT
  3AL ',' EACH SPECIES ',' CELL#')
  3 FORMAT(1X,I4,I5,3X,2F8.3,2F7.1,E12.3,2X,I4,2X,3I5,3X,I4)
  4 FORMAT(2X,'-----TOTALS-----',E12.4,8X,3I5)
  WRITE(6,1)
  WRITE(6,2)
  DO 200 I=1,NBX
  X=(XC(I)-XLIM(1))*AKN

```

PRA0850

PRA0880

PRA0940

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PRA0950

PRA0960

PRA0970

PRA0980

PRA1040

PRA1050

PRA1060

PRB0070

PRB0080

PRB0090

PRB0100

PRB0150

PRB0170

PRB0180

ORIGINAL PAGE 13
OF POOR QUALITY

```

Y=YC(I)*AKT*2.0
IX=X*NW+1
IY=Y*NH+1
M1=NB(1,I)
M2=0
IF(MSP.GE.2) M2=NB(2,I)
M3=0
IF(MSP.GE.3) M3=NB(3,I)
MM=M1+M2+M3
140 WRITE(6,3) IX,IY,X,Y,ZC(I),PNB(I),MM,M1,M2,M3,I
200 CONTINUE
NM2=0
IF(MSP.GE.2) NM2=NM(2)
NM3=0
IF(MSP.GE.3) NM3=NM(3)
WRITE(6,4) PNA,NM(1),NM2,NM3
RETURN
END

```

PRB0370

PRB0390

PRB0400

```

SUBROUTINE CELL(A,B,KW,KH,XO,DELANG,NWEDG,XC,YC,ZC,PNB)
DIMENSION XC(1),YC(1),ZC(1),PNB(1)
COMMON /THIRD/PI

```

CELL030

CELL040

CELL050

CELL060

THE PURPOSE OF THIS SUBROUTINE IS TO

1. COMPUTE THE VOLUME OF EACH CELL (ALL 3 POSSIBLE LEVELS) AND STORE THE RESULT IN THE ARRAY CALLED 'PNB'.

CELL070

CELL080

2. COMPUTE THE X, R, AND THETA COORDINATES OF THE CENTER OF EACH CELL (ALL 3 POSSIBLE LEVELS) AND STORE THE RESULTS IN CELLS CALLED 'XC', 'YC', AND 'ZC'.

CELL090

CELL100

CELL110

CELL120

CELL130

```

I=0
X=XO-0.5*A
FACTOR=DELANG*PI*B*B*A/180.
DO 110 K=1,KW
X=X+A
Y=-.5*B
DO 110 L=1,KH
Y=Y+B
Z=-.5*DELANG
DO 110 M=1,NWEDG
Z=Z+DELANG
I=I+1
XC(I)=X
YC(I)=Y
ZC(I)=Z
110 PNB(I)=FACTOR*(2*L-1)
RETURN
END

```

CELL360

CELL370

```

SUBROUTINE IMPACT(RM,G1,G2,G3,ET,EI,PHI,CHI,ETA,XM,CIM)
COMMON/THIRD/PI
IF(PHI.EQ.0.) GO TO 20

```



```

      IF (CHI.EQ.0.) GO TO 20
      DP=PHI*CHI-1.
      DS=PHI*(2.-.5*ETA)-1.
      E=ET+EI
10    X=RAND(0)
      IF (X.EQ.0.0) GO TO 10
      XT=X**DP*(1.-X)**DS
      IF (XT.GT.XM) GO TO 15
      CIM=CIM+XT
      IF (CIM.LT.XM) GO TO 10
      CIM=CIM-XM
15    ET=(1.-PHI)*ET+(1.-X)*PHI*E
      EI=(1.-PHI)*EI+X*PHI*E
20    GP=SQRT(ET/RM)
      EP=2.*PI*RAND(0)
      CSX=2.*RAND(0)-1.
      SSX=SQRT(1.-CSX**2)
      G1=GP*CSX
      G2=GP*SSX*COS(EP)
      G3=GP*SSX*SIN(EP)
      RETURN
      END

```

ORIGINAL PAGE 19
OF POOR QUALITY

```

SUBROUTINE GAS(NWEDG,DELANG,ND,BTA,C1,DPA,NM,FNB,DB,NB,NBM,NBN,
1PAU,PAV,PAW,PAX,PAY,PAZ,XLIM,COEPP,LM,I2,I3,XCB,TB,LARGE,
2MNM,MNB,DEBUG1,LCOL,IP,ER,CHI,CNG,CMG,I,LB)
      INTEGER*2 LM(I,1),LCOL(I,1),LB(1),NBM(I,1),NBN(1)
      INTEGER*2 NB
      LOGICAL DUMP,DEBUG1
      DIMENSION BTA(1),C1(1),DPA(1),NM(1),FNB(1),CHI(1)
      DIMENSION DB(I,1),NB(I,1),PAU(I,1),PAV(I,1),PAW(I,1)
      DIMENSION PAX(I,1),PAY(I,1),PAZ(I,1),ER(I,1),COEPP(4),XLIM(1)
      DIMENSION CNG(1),CMG(1),XCB(1),TB(1)
      COMMON /FIRST/NL,NW,NH
      COMMON /SECND/BW,BH,RMP,RMN,RMF
      COMMON /THIRD/PI,NREG,S,SINANG,COSANG,AKN,AKT
      COMMON /FORTH/NBX,RM,XR,DUMP,C9,LL(3)

```

GAS0060

-----GAS0150

GAS0160

GAS0170

GAS0180

GAS0190

GAS0200

GAS0210

GAS0220

GAS0230

GAS0250

GAS0260

GAS0270

GAS0280

GAS0290

-----GAS0300

THE PURPOSE OF THIS SUBROUTINE IS TO

1. COMPUTE THE INITIAL VELOCITY OF EACH MOLECULE.
THE VELOCITY ARRAYS ARE 'PAU', 'PAV', AND 'PAW'.
2. COMPUTE THE INITIAL POSITION OF EACH MOLECULE.
THE POSITION ARRAYS ARE 'PAX', 'PAY', AND 'PAZ'.
3. CREATE AN ARRAY WHICH STORES THE CELL POPULATIONS-
'NB' FOR THE ACTUAL POPULATIONS
4. CREATE A CROSS-REFERENCING ARRAY (WHOSE CONSTRUCTION IS
INDICATED BY A COMMENT CARD) CALLED 'LM'.
5. COMPUTE AN ARRAY WHICH STORES THE NUMBER DENSITY IN EACH
CELL

2 FORMAT(/' SOMETHING IS WRONG WITH BOX NUMBERING IN GAS'/9I5,5E14.5
1//)

3 FORMAT(/' SOMETHING WRONG IN CELL VOLUMES IN GAS'/5X,5I5,2E14.5)

```

4  FORMAT(' NB(' ,I2,' ,',I4,' ) POPULATION EXCEEDED ',I3,' IN GAS')      GAS0340
   BP=(1.-RMN)
   CP=1.-RMN**2
   DO 180 MT=1,IP
   N=0
110  N=N+1
      IF(N.GT.LL(MT)) GO TO 180
      IF(N.GT.MNM) GO TO 190
120  P=.001+.998*RAND(0)
      IF(BP.NE.0.0) P=(1.-SQRT(1.-CP*P))/BP
      X=XLIM(1)+.X*XR
      R=RM*SQRT(RAND(0))
      D=PI*RAND(0)
      PAX(MT,N)=X
      PAY(MT,N)=R*COS(D)
      PAZ(MT,N)=R*SIN(D)
      DO 126 J=2,ND
      IF(X.LE.XCB(J)) GO TO 128
126  CONTINUE
      WRITE(6,3) J,ND,MT,LL(MT),N,X,XCB(J)
      IF(DUMP) CALL ABEND(14)
      STOP
128  TL=TB(J-1)+(TB(J)-TB(J-1))*(X-XCB(J-1))/(XCB(J)-XCB(J-1))
130  V=4.*RAND(0)
      VV=V*V
      C1(MT)=C1(MT)+VV*EXP(1.-VV)
      IF(C1(MT).LT.1.) GO TO 130
      C1(MT)=C1(MT)-1.
      A=1.-2.*RAND(0)
      B=SQRT(1.-A*A)
      C=2.*PI*RAND(0)
      V=V/BTA(MT)*SQRT(TL)
      PAU(MT,N)=V*A
      PAV(MT,N)=V*B*COS(C)
      PAW(MT,N)=V*B*SIN(C)
      EX=0.0
      IF(CHI(MT).LE.-1.) GO TO 136
135  EX=9.*RAND(0)
      IF(EX.EQ.0.0) GO TO 135
      XT=EX**CHI(MT)*EXP(-EX)
      IF(XT.GE.CMG(MT)) GO TO 136
      CNG(MT)=CNG(MT)+XT
      IF(CNG(MT).LT.CMG(MT)) GO TO 135
      CNG(MT)=CNG(MT)-CMG(MT)
136  ER(MT,N)=EX*TL
      TANG=180.*(1.-D/PI)
      IWDGE=TANG/DELANG+1
      IF(IWDGE.GT.NWEDG) IWDGE=NWEDG
      L=X/BW
      IF(L.GE.NW) L=NW-1
      M=R/BH
      IF(M.GE.NH) M=NH-1
      K=NWEDG*(L*NH+M)+IWDGE
      IF(K.LE.NBX) GO TO 165
      WRITE(6,2) L,M,IWDGE,NW,NH,NWEDG,K,MT,N,X,Y,Z,R,TANG

```

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GAS0370
GAS0380

GAS0390

GAS0400
GAS0410
GAS0420
GAS0430
GAS0440
GAS0450
GAS0460
GAS0470

GAS0510

GAS0720

GAS0780
GAS0790

IF (DUMP) CALL ABEND(11)	GAS0830
STOP	GAS0840
165 J=NB (MT, K) +1	GAS1190
LCOL (MT, N)=0	
IF (J.LE.MNB) GO TO 166	GAS1240
IF (DEJUG1) WRITE (6,4) MT,K,MNB	GAS1250
166 NB (MT, K) =J	GAS1270
LB (N)=K	
167 IF (N.LT.LL (MT)) GO TO 110	
NM (MT)=N	GAS1310
NBM (MT, 1)=0	
DO 170 N=1,NBX	GAS1330
A=NB (MT, N)	GAS1360
DB (MT, N)=A*DFA (MT) /PNB (N)	GAS1370
NBM (MT, N+1)=NBM (MT, N) +NB (MT, N)	
NBN (N)=NBM (NT, N)	
170 CONTINUE	GAS1380
NG=NM (MT)	
DO 175 N=1, NG	
NQ=LB (N)	
NBN (NQ)=NBN (NQ) +1	
NA=NBN (NQ,	
175 LM (MT, NA) =N	
180 CONTINUE	GAS1390
RETURN	GAS1400
190 LARGE=1	GAS1410
RETURN	GAS1420
END	GAS1430

ORIGINAL PAGE 13
OF POOR QUALITY

SUBROUTINE FLOW (NWEDG, MNM, LARGE, BTA, C1, C7, C8, ENT, REM, LCOL, IP, NM,
1SN, ST, TBI, PAU, PAV, PAW, PAX, PAY, PAZ, ER, CHI, CNG, CMG, I, JV, FCOL, VEL,
2PPV)

INTEGER*2 LCOL
DIMENSION BTA (1), NM (1), SN (1), ST (1)
DIMENSION C1 (1), C7 (1), C8 (1), CMG (1), FCOL (1), VELK (4)
DIMENSION PAU (I, 1), PAV (I, 1), PAW (I, 1), PAX (I, 1), PAY (I, 1), PAZ (I, 1)
DIMENSION ENT (3, 1), REM (3, 1), LCOL (I, 1), ER (I, 1), CHI (1), CNG (1)
DIMENSION VEL (JV, 4, 1), PPV (JV, 4, 1)
COMMON /THIRD/PI
COMMON /FORTH/NBX, RM, XR

FLO0100
FLO0110
FLO0120
FLO0130
FLO0140
FLO0150

THE PURPOSE OF THIS SUBROUTINE IS TO ADD A NEW BATCH OF MOLECULES
TO THE SAMPLE THROUGH THE UPSTREAM BOUNDARY.

DO 370 MT=1, IP	
XGO=0.	FLO0180
E=1.	FLO0190
PRAC=FCOL (MT)	
ARG=SN (MT)	
STT=ST (MT)	
TV=1./BTA (MT)	
TR=1.	
DO 180 NT=1, 2	FLO0200
VM=(SQRT (ARG**2+2.) +ARG) /2.	
SM=AMAX1 (0.0, VM-4.)	

```

SMM=VM+4.-SM
AM=ENT(MT,NT)+REM(MT,NT)
M=AM
REM(MT,NT)=AM-M
IF(M.EQ.0) GO TO 170
DO 160 N=1,M
IF(NM(MT).GE.MNM) GO TO 380
NM(MT)=NM(MT)+1
NMX=NM(MT)

```

ORIGINAL PAGE 12
OF POOR QUALITY

FLO0340
FLO0360

```

R=RM*SQRT(RAND(0))
D=PI*RAND(0)
PAY(MT,NMX)=R*COS(D)
PAZ(MT,NMX)=R*SIN(D)
LCOL(MT,NMX)=0
IF(FRAC.EQ.0.0) GO TO 130
PF=RAND(0)
IF(PF.GT.FRAC) GO TO 130
KMX=3

```

FLO0370
FLO0390

```

VELK(4)=0.0
IF(CHI(MT).GT.-1) KMX=4
DO 110 K=1,KMX
P=RAND(0)
DO 102 J=2,JV
IF(PFV(J,K,MT).GT.P) GO TO 105

```

FLO0420
FLO0430

102 CONTINUE

```

VELK(K)=VEL(JV,K,MT)
GO TO 110

```

105 VELK(K)=VEL(J-1,K,MT)+(P-PFV(J-1,K,MT))*(VEL(J,K,MT)-VEL(J-1,K,MT))/(PFV(J,K,MT)-PFV(J-1,K,MT))

110 CONTINUE

```

PAU(MT,NMX)=VELK(1)
PAV(MT,NMX)=VELK(2)
PAW(MT,NMX)=VELK(3)
ER(MT,NMX)=VELK(4)
GO TO 160

```

130 V=SM+RAND(0)*SMM

```

C1(MT)=C1(MT)+V*EXP(VM**2-V**2+2.*ARG*(V-VM))/VM

```

```

IF(C1(MT).LT.1.) GO TO 130

```

```

C1(MT)=C1(MT)-1.

```

```

PAU(MT,NMX)=E*V*TV

```

140 V=8.*RAND(0)-4.

```

C7(MT)=C7(MT)+EXP(-V*V)

```

```

IF(C7(MT).LT.1.) GO TO 140

```

```

C7(MT)=C7(MT)-1.

```

```

PAV(MT,NMX)=STT+V*TV

```

150 V=8.*RAND(0)-4.

```

C8(MT)=C8(MT)+EXP(-V*V)

```

```

IF(C8(MT).LT.1.) GO TO 150

```

```

C8(MT)=C8(MT)-1.

```

```

PAW(MT,NMX)=V*TV

```

```

X=0.0

```

```

IF(CHI(MT).LE.-1.) GO TO 156

```

155 X=9.*RAND(0)

```

IF(X.LE.0.0) GO TO 155

```

```

XT=X**CHI(MT)*EXP(-X)

```

FLO0480
FLO0490

FLO0510
FLO0520
FLO0530
FLO0540

FLO0560
FLO0570
FLO0580
FLO0590

```

IF (XT.GE.CNG(MT)) GO TO 156
CNG(MT)=CNG(MT)+XT
IF (CNG(MT).LT.CNG(MT)) GO TO 155
CNG(MT)=CNG(MT)-CNG(MT)
156 ER(MT,NMX)=X*TR
160 PAX(MT,NMX)=XGO
170 CONTINUE
    ARG=0.0
    XGO=XR
    E=-1.
    FRAC=0.0
    STT=0.0
    TV=SQRT(TBI)/BTA(MT)
    TR=TBI
180 CONTINUE
370 CONTINUE
    RETURN
380 LARGE=2
    RETURN
    END

```

FLO0610
FLO0620

FLO0640
FLO0650

FLO0660
FLO1140
FLO1150
FLO1160
FLO1170
FLO1180

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```

SUBROUTINE COLIDE(CN,CM,WTM,DB,DBA,NB,NCOL,LCOL,PAU,PAV,PAW,ER,T,
1LM,MT,I2,I3,ETA,PHI,CHI,CN8,NP,NBM)
    INTEGER TIME
    INTEGER*2 LM(NP,1),LCOL(NP,1)
    INTEGER*2 NBM,NB
    DIMENSION CN(3,3,1),CM(3,3,1),WTM(1),DB(NP,1),DBA(NP,1),NB(NP,1)
    DIMENSION NBM(NP,1),NCOL(3,1),T(NP,NP,1),ETA(3,1),PHI(3,1),CHI(1)
    DIMENSION PAU(NP,1),PAV(NP,1),PAW(NP,1),ER(NP,1),CN8(3,1),WA(2)
    COMMON /PORTH/NBX
    COMMON /PIPTH/ND,TIME,DTM

```

COL0030

COL0080
COL0090
COL0100

THE PURPOSE OF THIS SUBROUTINE IS TO ADVANCE THE ELAPSED TIMES IN
CELLS BY AN AMOUNT APPROXIMATELY EQUAL TO THE PRE-SELECTED COLLIS
TIME. THERE ARE FOUR TIMES FOR EACH CELL, SAVED IN AN ARRAY CALLE
'T', CORRESPONDING TO THE FOUR TYPES OF MOLECULAR COLLISIONS WHIC
CAN OCCUR. TO ADVANCE THE VARIOUS TIMES, AN APPROPRIATE NUMBER OF
THE CORRESPONDING MOLECULAR COLLISIONS IS COMPUTED. THE ACTUAL
MOLECULES TO COLLIDE ARE SELECTED AT RANDOM, AND THEIR VELOCITY V
DIRECTIONS AFTER COLLISION ARE SELECTED AT RANDOM.

COL0110
COL0120
COL0130
COL0140
COL0150
COL0160
COL0170
COL0180
COL0190

```

AIME=DTM*TIME
DO 240 MTA=1,MT
DO 230 MTB=1,MTA
D = WTM(MTA) + WTM(MTB)
WA(MTA)=WTM(MTA)/D
WA(MTB)=WTM(MTB)/D
RM=WTM(MTA)*WTM(MTB)/D
CHT=CHI(MTA)+CHI(MTB)+2.0
PHT=PHI(MTA,MTB)
ETT=ETA(MTA,MTB)
DO 220 N=1,NBX
IF (T(MTA,MTB,N).LT.AIME) GO TO 100
IF (T(MTB,MTA,N).GE.AIME) GO TO 220
100 NA=NB(MTA,N)*NB(MTB,N)

```


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```

IF (MTA.EQ.MTB) NA=(NA-NB(MTA,N))/2.
IF (NA.LT.1) GO TO 220
KS=0
#20 KC=0
CPUT=ELTIME(0)
KS=KS+1
IF (KS.GT.NA) GO TO 220
130 KC=KC+1
IF (KC.GT.NA) GO TO 220
#35 I=NB(MTA,N)*RAND(0)+1+NBM(MTA,N)
IF (I.GT.NBM(MTA,N+1)) I=NBM(MTA,N+1)
J=LM(MTA,I)
140 K=NB(MTB,N)*RAND(0)+1+NBM(MTB,N)
IF (K.GT.NBM(MTB,N+1)) K=NBM(MTB,N+1)
IF (MTA.EQ.MTB.AND.I.EQ.K) GO TO 140
L=LM(MTB,K)
GM1=WA(MTA)*PAU(MTA,J)+WA(MTB)*PAU(MTB,L)
GM2=WA(MTA)*PAV(MTA,J)+WA(MTB)*PAV(MTB,L)
GM3=WA(MTA)*PAW(MTA,J)+WA(MTB)*PAW(MTB,L)
G1=PAU(MTA,J)-PAU(MTB,L)
G2=PAV(MTA,J)-PAV(MTB,L)
G3=PAW(MTA,J)-PAW(MTB,L)
GS=G1**2+G2**2+G3**2
IF (GS.LT.1.0E-8) GO TO 130
ET=RM*GS
EI=ER(MTA,J)+ER(MTB,L)
VR=GS**(.5-ETT/2.)
IF (VR.GE.CN(MTA,MTB,1)) GO TO 160
CN(MTA,MTB,1)=CN(MTA,MTB,1)+VR
IF (CN(MTA,MTB,1).LT.CM(MTA,MTB,1)) GO TO 130
CN(MTA,MTB,1)=CN(MTA,MTB,1)-CM(MTA,MTB,1)
160 CONTINUE
CPUT=ELTIME(0)
CALL IMPACT(RM,G1,G2,G3,ET,EI,PHT,CHT,ETT,CN(MTA,MTB,2),CN(MTA,MTB
1,2))
165 CONTINUE
CPUT=ELTIME(0)
IF (PHT.EQ.0.) GO TO 175
X1=0.0
IF (CHI(MTA).EQ.-1.) GO TO 175
X1=1.0
IF (CHI(MTB).EQ.-1.) GO TO 175
170 X1=RAND(0)
IF ((CHI(MTA).EQ.0.).AND.(CHI(MTB).EQ.0.)) GO TO 175
XT=X1**CHI(MTA)*(1.-X1)**CHI(MTB)
IF (XT.GT.CN(MTA,MTB,3)) GO TO 175
CN(MTA,MTB,3)=CN(MTA,MTB,3)+XT
IF (CN(MTA,MTB,3).LT.CM(MTA,MTB,3)) GO TO 170
CN(MTA,MTB,3)=CN(MTA,MTB,3)-CM(MTA,MTB,3)
175 CONTINUE
C=DBA(MTA,N)
D=DBA(MTB,N)
IF (C.EQ.0.0) C=DB(MTA,N)
IF (D.EQ.0.0) D=DB(MTB,N)
IF (T(MTA,MTB,N).GE.AIME) GO TO 180

```

```

PAU(MTA,J)=GM1+WA(MTB)*G1
PAV(MTA,J)=GM2+WA(MTB)*G2
PAW(MTA,J)=GM3+WA(MTB)*G3
IF(PHT.GT.0.) ER(MTA,J)=X1*EI
LCOL(MTA,J)=1+LCOL(MTA,J)
NCOL(MTA,MTB)=NCOL(MTA,MTB)+1
T(MTA,MTB,N)=T(MTA,MTB,N)+CN8(MTA,MTB)/NB(MTA,N)/D/VR
IF(MTA.EQ.MTB) GO TO 190
180 IF(T(MTB,MTA,N).GE.AIME) GO TO 210
190 PAU(MTB,L)=GM1-WA(MTA)*G1
PAV(MTB,L)=GM2-WA(MTA)*G2
PAW(MTB,L)=GM3-WA(MTA)*G3
IF(PHT.GT.0.) ER(MTB,L)=(1.-X1)*EI
LCOL(MTB,L)=1+LCOL(MTB,L)
NCOL(MTB,MTA)=NCOL(MTB,MTA)+1
T(MTB,MTA,N)=T(MTB,MTA,N)+CN8(MTB,MTA)/NB(MTB,N)/C/VR
210 CONTINUE
IF(T(MTA,MTB,N).LT.AIME.OR.T(MTB,MTA,N).LT.AIME) GO TO 120
220 CONTINUE
230 CONTINUE
240 CONTINUE
RETURN
END

```

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COL0810

COL0920

```

SUBROUTINE MOVE(KSWCH,AKN,NWEDG,XSTART,I2,I3,I4,I5,DELANG,
1BTA,C2,C3,DFA,PL,HTI,HTR,JNT,KNM,NM,XCB,XLIM,CTI,CTR,CNI,
2CNR,ALPHA,SIGMA,COEFF,HTS,HTSI,NTS,UTL,UTT,VTS,PAU,PAV,PAW,PAX,
3PAY,PAZ,LCOL,TB,IP,ER,CHI,CNG,CMG,I,UTLI,UTTI,VTSI,IPLUX)

```

```

INTEGER*2 LCOL(I,1)

```

```

INTEGER TST,TIME

```

```

LOGICAL DUMP

```

```

REAL LAM,MU,NU

```

```

DIMENSION BTA(1),C2(1),C3(1),PL(1),HTI(1)

```

```

DIMENSION HTR(1),TB(1),XCB(1),ALPHA(3,1),SIGMA(3,1),COEFF(4)

```

```

DIMENSION PAU(1,1),PAV(1,1),PAW(1,1),CTI(3,1),CTR(3,1)

```

```

DIMENSION CNI(3,1),CNR(3,1),DFA(1),JNT(1),XLIM(1),KNM(1),NM(1)

```

```

DIMENSION HTS(3,I2,I3),HTSI(3,I2,I3),NTS(3,I2,I3)

```

```

DIMENSION UTLI(3,I2,I3),UTTI(3,I2,I3),VTSI(3,I2,I3)

```

```

DIMENSION UTL(3,I2,I3),UTT(3,I2,I3),VTS(3,I2,I3)

```

```

DIMENSION PAX(1,1),PAY(1,1),PAZ(1,1)

```

```

DIMENSION ER(1,1),CNG(1),CMG(1),CHI(1),IPLUX(3,2)

```

```

COMMON /THIRD/PI,NREG

```

```

COMMON /PORTH/NBX,RM,XR,DUMP

```

```

COMMON /FIFTH/ND,TIME,DTM,TI,ITS,ITP,TST

```

```

COMMON /SVNTH/LAM,MU,NU,MT,N,J,XI,YI,ZI,TUSE

```

```

NAMELIST/CHECK/TIME,X,Y,Z,DX,DY,DZ,TLEFT,RADS,RMS,XR

```

MOV0170

MOV0180

MOV0200

```

-----
THE PURPOSE OF THIS SUBROUTINE IS TO ADVANCE THE SPATIAL POSITION
OF ALL THE MOLECULES BY AN AMOUNT APPROPRIATE TO THEIR CURRENT VE
LOCITIES AND THE PRE-SELECTED COLLISION TIME.
-----

```

210

220

230

240

250

```

NAREA=NREG+1

```

```

RMS=RM**2

```

```

DO 150 MT=1,IP

```

```

N=KNM(MT)

```

MOV0270

MOV0290

```

10 N=N+1
   TLEFT=DTM
   IF (KSWCH.EQ.1) TLEFT=TLEFT*RAND(0)
   IF (N.GT.NM(MT)) GO TO 150
15 LAM=PAU(MT,N)
   MU=PAV(MT,N)
   NU=PAW(MT,N)
   XI=PAX(MT,N)
   YI=PAY(MT,N)
   ZI=PAZ(MT,N)
   DX=TLEFT*LAM
   DY=TLEFT*MU
   DZ=TLEFT*NU
   X=XI+DX
   Y=YI+DY
   Z=ZI+DZ
   RADS=Y**2+Z**2
   TUSE=TLEFT
   KS=0
   IF (RADS.GT.RMS) CALL INTERS(X,Y,Z,RMS,KS)
   IF (X.GT.XLIM(1)) GO TO 18
   IPLUX(MT,1)=IPLUX(MT,1)+1
   GO TO 100
18 IF (X.LT.XLIM(NAREA)) GO TO 19
   IPLUX(MT,2)=IPLUX(MT,2)+1
   GO TO 100
19 CONTINUE
   IF (KS.GT.0) CALL DRAG(AKN,NWEDG,XSTART,I2,I3,I4,I5,DELANG,JNT,BTA,
1C2,C3,DFA,FL,HTI,HTR,TB,XCB,CTI,CTR,CNI,CNR,ALPHA,SIGMA,RM,HTS,
2HTSI,NTS,UTL,UTT,VTS,LCOL,IP,ER,CHI,CNG,CMG,I,UTLI,UTTI,VTSI)
   IF (Z.GT.0.0) GO TO 20
   Z=-Z
   NU=-NU
20 CONTINUE
   PAX(MT,N)=X
   PAY(MT,N)=Y
   PAZ(MT,N)=Z
   PAU(MT,N)=LAM
   PAV(MT,N)=MU
   PAW(MT,N)=NU
   IF ((X.GT.0.0).AND.(X.LT.XR).AND.((Y**2+Z**2).LT.RMS)) GO TO 25
   WRITE(6,CHECK)
   GO TO 100
25 CONTINUE
   TLEFT=TLEFT-TUSE
   IF (TLEFT.GT.0.0) GO TO 15
   GO TO 10
100 NZ=NZ(MT)
   PAX(MT,N)=PAX(MT,NZ)
   PAY(MT,N)=PAY(MT,NZ)
   PAZ(MT,N)=PAZ(MT,NZ)
   PAU(MT,N)=PAU(MT,NZ)
   PAV(MT,N)=PAV(MT,NZ)
   PAW(MT,N)=PAW(MT,NZ)
   ER(MT,N)=ER(MT,NZ)

```

MOV0300
 MOV0310
 MOV0320
 MOV0330
 MOV0340
 MOV0360
 MOV0370
 MOV0380
 MOV0390
 MOV0400
 MOV0410
 MOV0420
 MOV0430
 MOV0440
 MOV0450
 MOV0460

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```

LCOL (MT, N) = LCOL (MT, NZ)
N = N - 1
NM (MT) = NM (MT) - 1
GO TO 10
150 CONTINUE
RETURN
END

```

MOV0900
MOV0910
MOV0920
MOV0930
MOV0940
MOV0950
MOV0960

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```

SUBROUTINE ACCUM (I2, I3, FNB, NB, PAU, PAV, PAW, ER, TMP, TRP, XV, YV, ZV, LM,
1IP, I, NBM)
INTEGER*2 LM (I, 1), NBM (I, 1)
INTEGER*2 NB
DIMENSION FNB (1), NB (I, 1), PAU (I, 1), PAV (I, 1), PAW (I, 1), TMP (I, 1)
DIMENSION XV (I, 1), YV (I, 1), ZV (I, 1), ER (I, 1), TRP (I, 1)
COMMON /FORTH/NBX

```

ACUM050

ACUM060

ACUM070

ACUM080

ACUM090

ACUM100

ACUM120

ACUM160

ACUM170

ACUM180

ACUM190

ACUM200

ACUM210

ACUM220

ACUM230

ACUM240

ACUM250

ACUM260

ACUM270

ACUM280

ACUM300

ACUM310

ACUM320

```

-----
THE PURPOSE OF THIS SUBROUTINE IS TO ACCUMULATE TEMPERATURES,
VELOCITIES, AND DENSITIES IN VARIOUS ARRAYS FOR DETERMINING THE
AVERAGE FLOW FIELD PROPERTIES AFTER STEADY-STATE HAS BEEN REACHED
-----
DO 180 N=1, NBX
DO 110 MT=1, IP
XV (MT, N) = 0.0
YV (MT, N) = 0.0
ZV (MT, N) = 0.0
TMP (MT, N) = 0.
TRP (MT, N) = 0.0
TTX=0.
TTY=0.
TTZ=0.
TTR=0.0
M=NB (MT, N)
IF (M.LT.1) GO TO 110
U=0.
V=0.
W=0.
DO 100 L=1, M
NA=NB (MT, N) + L
J=LM (MT, NA)
PU=PAU (MT, J)
PV=PAV (MT, J)
PW=PAW (MT, J)
U=U+PU
V=V+PV
W=W+PW
TTR=TTR+ER (MT, J)
TTX=TTX+PU*PU
TTY=TTY+PV*PV
100 TTZ=TTZ+PW*PW
M=NB (MT, N)
XV (MT, N) = U/M
YV (MT, N) = V/M
ZV (MT, N) = W/M
TMP (MT, N) = (TTX+TTY+TTZ)/M
TRP (MT, N) = TTR/M

```

ACUM390

ACUM400

ACUM410

ACUM420

110 CONTINUE
180 CONTINUE
RETURN
END

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ACUM430
ACUM440
ACUM450
ACUM460

SUBROUTINE AVRGE (FNB,DB,DBA,NB,NBT,XV,YV,ZV,XVA,YVA,ZVA,TMP,TMPA,
1TRP,TRPA,IP,I)
INTEGER*2 NB,NBT
DIMENSION FNB(1),DB(1,1),DBA(1,1),NB(1,1),NBT(1,1),TMP(1,1)
DIMENSION TMPA(1,1),XV(1,1),XVA(1,1),YV(1,1),YVA(1,1),ZV(1,1)
DIMENSION ZVA(1,1),TRP(1,1),TRPA(1,1)
COMMON /PORTH/NBX

AVG0050

THE PURPOSE OF THIS SUBROUTINE IS TO COMPUTE THE AVERAGE FLOW
FIELD PROPERTIES.

AVG0060

AVG0070

AVG0080

AVG0090

DO 110 N=1,NBX

AVG0110

DO 100 MT=1,IP

A=NBT(MT,N)

AVG0150

B=NB(MT,N)

C=A+B

AVG0170

NBT(MT,N)=C

AVG0180

IF (C.LT.1.) GO TO 100

AVG0190

DBA(MT,N)=(DBA(MT,N)*A+DB(MT,N)*B)/C

AVG0200

XVA(MT,N)=(XVA(MT,N)*A+XV(MT,N)*B)/C

AVG0210

YVA(MT,N)=(YVA(MT,N)*A+YV(MT,N)*B)/C

AVG0220

ZVA(MT,N)=(ZVA(MT,N)*A+ZV(MT,N)*B)/C

AVG0230

TMPA(MT,N)=(TMPA(MT,N)*A+TMP(MT,N)*B)/C

AVG0240

TRPA(MT,N)=(TRPA(MT,N)*A+TRP(MT,N)*B)/C

100 CONTINUE

AVG0250

110 CONTINUE

AVG0260

RETURN

AVG0270

END

AVG0280

SUBROUTINE DRAG (AKN,NWEDG,XSTART,I2,I3,I4,I5,DELANG,JNT,BTA,
1C2,C3,DFA,FL,HTI,HTR,TB,XCB,CTI,CTR,CNI,CNR,ALPHA,SIGMA,RM,HTS,
2HTSI,NTS,UTL,UTT,VTS,LCOL,IP,ER,CHI,CNG,CMG,I,UTLI,UTTI,VTSI)
INTEGER*2 LCOL

DRG0040

INTEGER TIME,TST

DRG0050

REAL LAM,MU,NU,JAY,KAY

DRG0060

DIMENSION BTA(1),C2(1),C3(1),FL(1),HTI(1),CTI(3,1),CTR(3,1)

DIMENSION HTR(1),TB(1),XCB(1),ALPHA(3,1),SIGMA(3,1)

DIMENSION HTS(3,12,13),HTSI(3,12,13),NTS(3,12,13)

DIMENSION UTL(3,12,13),UTT(3,12,13),VTS(3,12,13),LCOL(1,1)

DIMENSION UTLI(3,12,13),UTTI(3,12,13),VTSI(3,12,13)

DIMENSION CNI(3,1),CNR(3,1),DFA(1),JNT(1)

DIMENSION ER(1,1),CNG(1),CMG(1),CHI(1)

COMMON /THIRD/PI

DRG0140

COMMON /FIFTH/ND,TIME,DTM,TI,ITS,ITP,TST

DRG0150

COMMON /SVNTH/LAM,MU,NU,MT,N,J,XCL,YCL,ZCL

DRG0160

THE PURPOSE OF THIS SUBROUTINE IS TO ACCUMULATE THE DRAG AND HEAT

170

TRANSFER INCREMENTS ON THE BODY CONTRIBUTED BY EACH MOLECULE WHICH

180

COLLIDES WITH THE BODY. IN ADDITION, EACH MOLECULE WHICH COLLIDES

190

WITH THE BODY IS ASSIGNED AN APPROPRIATE NEW VELOCITY (OF REFLECTION)

200

210

WHICH IS USED TO CONTINUE ITS SPATIAL TRANSLATION (IN SUBROUTINE DRAG)

230

```

CALL NORMAL(JAY,KAY,RH)
JNT(MT)=JNT(MT)+1
TANG=180.*ATAN2(ZCL,-YCL)/PI
IWDG=TANG/DELANG+1
IF (IWDG.GT.NWEDG) IWDG=NWEDG
D=(LAM*LAM+MU*MU+NU*NU)
G=ER(MT,N)
H=G
DO 100 M=2,ND
IF (XCL.LT.XCB(M)) GO TO 110
100 CONTINUE
110 TBX=TB(M-1)+(TB(M)-TB(M-1))*(XCL-XCB(M-1))/(XCB(M)-XCB(M-1))
WI=(NU*JAY-MU*KAY)
VID=MU*JAY+NU*KAY
UID=LAM
E=RAND(0)
IF (E.LT.SIGMA(MT,M)) GO TO 115
VRD=-VID
URD=UID
WR=WI
GO TO 125
115 V=4.*RAND(0)
C2(MT)=C2(MT)+.544331*V*V*V*EXP(1.5-V*V)
IF (C2(MT).LT.1.) GO TO 115
C2(MT)=C2(MT)-1.
IF (NTS(MT,M,IWDG).NE.0) GO TO 117
ATR=ALPHA(MT,M)*TBX /SIGMA(MT,M)
GO TO 118
117 ATR=ALPHA(MT,M)*TBX /SIGMA(MT,M)+(1.-ALPHA(MT,M)/SIGMA(MT,M))*HTS
118 I(MT,M,IWDG)/NTS(MT,M,IWDG)/(3.+CHI(MT))
ABR=SQRT(ATR)
V=V*ABR/BTA(MT)
AA=RAND(0)
A=SQRT(AA)
B=SQRT(1.-AA)
C=2.*PI*RAND(0)
VRD=V*A
URD=V*B*COS(C)
WR=V*B*SIN(C)
IF (CHI(MT).EQ.-1.) GO TO 125
122 X=9.*RAND(0)
IF (X.EQ.0.0) GO TO 122
XTEMP=1.0
IF (CHI(MT).NE.0.0) XTEMP=X**CHI(MT)
CNG(MT)=CNG(MT)+XTEMP*EXP(-X)
IF (CNG(MT).LT.CMG(MT)) GO TO 122
124 CONTINUE
CNG(MT)=CNG(MT)-CMG(MT)
IF (CNG(MT).GE.CMG(MT)) GO TO 124
ER(MT,N)=X*ATR
H=ER(MT,N)
125 UR=URD
LAM=UR

```

ORIGINAL PAGE 19
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DRG0310

DRG0320

DRG0330

DRG0340

DRG0350

DRG0370

DRG0400

DRG0410

DRG0420

DRG0430

DRG0440

DRG0450

DRG0460

DRG0470

DRG0480

DRG0490

DRG0500

DRG0510

DRG0520

DRG0530

DRG0550

DRG0560

DRG0570

DRG0610

DRG0620

DRG0630

DRG0640

DRG0650

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OF POOR QUALITY

```

MU=JAY*VRD-KAY*WR
NU=KAY*VRD+JAY*WR
IF (TIME.LE.TST) RETURN
XMZ= (XCL-XSTART)*AKN
YMZ= RM*AKN
B= (URD*URD+VRD*VRD+WR*WR)
UTI=UID*UID+WI*WI
UYI=-WI*KAY
UYR=-WR*KAY
FL (MT)=FL (MT)+DPA (MT)
HTI (MT)=HTI (MT)+D+G
HTR (MT)=HTR (MT)-B-H
CTI (MT,1)=CTI (MT,1)+UID
CTI (MT,2)=CTI (MT,2)+UYI
CTI (MT,3)=CTI (MT,3)+(XMZ*UYI-YMZ*UID)
CNI (MT,2)=CNI (MT,2)+VID*JAY
CNI (MT,3)=CNI (MT,3)+XMZ*JAY*VID
CTR (MT,1)=CTR (MT,1)-URD
CTR (MT,2)=CTR (MT,2)-UYR
CTR (MT,3)=CTR (MT,3)-(XMZ*UYR-YMZ*URD)
CNR (MT,2)=CNR (MT,2)-VRD*JAY
CNR (MT,3)=CNR (MT,3)-XMZ*JAY*VRD
NTS (MT,M,IWDG)=NTS (MT,M,IWDG)+1
UTLI (MT,M,IWDG)=UTLI (MT,M,IWDG)+UID
UTL (MT,M,IWDG)=UTL (MT,M,IWDG)+(UID-URD)
UTTI (MT,M,IWDG)=UTTI (MT,M,IWDG)+WI
UTT (MT,M,IWDG)=UTT (MT,M,IWDG)+(WI-WR)
VTSI (MT,M,IWDG)=VTSI (MT,M,IWDG)-VID
VTS (MT,M,IWDG)=VTS (MT,M,IWDG)+(VRD-VID)
HTSI (MT,M,IWDG)=HTSI (MT,M,IWDG)+D+G
HTS (MT,M,IWDG)=HTS (MT,M,IWDG)+D-B+G-H
RETURN
END

```

DRG0730
DRG0740
DRG0750
DRG0760

DRG0790

DRG0820
DRG0830
DRG0840
DRG0860

DRG0880
DRG0890
DRG0900
DRG0920

DRG0940
DRG0960
DRG0970
DRG0980

DRG1170
DRG1180

```

SUBROUTINE INTERS(X,Y,Z,RMS,KS)
REAL LAM,MU,NU
COMMON /SVNTH/LAM,MU,NU,MT,N,J,XI,YI,ZI,SE
1 FORMAT (/' SOMETHING IS WRONG IN INTERS'/6E15.6/4E15.6)
2 FORMAT (/' TROUBLE IN INTERS - TYME =' ,E15.6/6E15.6/4E15.6)
A=MU**2+NU**2
B= (YI*MU+ZI*NU)/A
C= (RMS-YI**2-ZI**2)/A
IF (C.LT.0.0) C=0.0
DISCR=B**2+C
IF (DISCR.GE.0.0) GO TO 10
WRITE (6,1) XI,YI,ZI,MU,NU,RMS,A,B,C,DISCR
STOP
10 TYME=SQRT (DISCR)-B
IF (C.EQ.0.0) TYME=0.0
IF ((TYME.LE.TUSE).AND. (TYME.GE.0.0)) GO TO 11
WRITE (6,2) TYME,XI,YI,ZI,MU,NU,RMS,A,B,C,DISCR
IF (TYME.GT.TUSE) TYME=TUSE
IF (TYME.LT.0.0) TYME=0.0
11 KS=1
TUSE=TYME

```

```

XI=XI+LAM*TYME
YI=YI+MU*TYME
ZI=ZI+NU*TYME
IF (ZI.GT.0.0) GO TO 20
ZI=-ZI
NU=-NU

```

ORIGINAL PAGE 13
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```

20 CONTINUE
YI=.9999*YI
ZI=.9999*ZI
X=XI
Y=YI
Z=ZI
RETURN
END

```

```

SUBROUTINE NORMAL(JAY,KAY,RM)
REAL LAM,MU,NU,JAY,KAY
COMMON /SVNTH/LAM,MU,NU,MT,N,J,XCL,YCL,ZCL
JAY=-YCL/RM
KAY=-ZCL/RM
RETURN
END

```

NORM020
NORM040

NORM130
NORM140

```

SUBROUTINE PRINT1(DT,COSANG,SINANG,RMA,RNU,DRP,PCP,HTP,PL,HTI,HTR,
1CTI,CTR,CNI,CNR)
DIMENSION DD(3),WD(2,5),PP(4,4),QQ(4,4),RR(4,4),SS(4,4),TT(4,4)
DIMENSION UU(4,4),P1(4,4),Q1(4,4),R1(4,4),PA(4),PB(4),PC(4)
DIMENSION PL(1),HTI(1),HTR(1),CTI(3,1),CTR(3,1),CNI(3,1),CNR(3,1)
DIMENSION RMA(1),RNU(1)
DATA WD/'X-PO','RCE ','Y-PO','RCE ','Z-MO','MENT','DRAG',' ','LPT10070
1IPT',' '/

```

PT10080
PT10090
PT10100
PT10110
PT10120
PT10130
PT10140
PT10150
PT10160
PT10170
PT10180

THE PURPOSE OF THIS SUBROUTINE IS TO PRINT OUT THE GROSS SURFACE
COEFFICIENTS OF THE BODY.

FORMATS

```

1 FORMAT(//1X,50('*'),' GROSS SURFACE COEFFICIENTS ',50('*')/' MOLEC
1ULAR WEIGHT',12X,P8.3,3(19X,P8.3)/25X,
2 'INC. REF. TOT. INC. REF. TOT. INC.
3REF. TOT. INC. REF. TOT. ')
10 FORMAT(' NUMBER FLUX ',4(F9.3,18X))
12 FORMAT(1X,2A4,2X,'SHEAR ',4(3F8.3,3X))
14 FORMAT(11X,'PRESSURE ',4(3F8.3,3X))
16 FORMAT(11X,'TOTAL ',4(3F8.3,3X)/)
18 FORMAT(' HEAT TRANSFER',7X,4(3F8.3,3X)/)

```

PT10280
PT10290
PT10300
PT10310

RM=0.0

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```

DO 50 MT=1,3
DD(MT)=RMA(MT)*RNU(MT)*DRF/DT
50 RMR=RMR+RMA(MT)*RNU(MT)
WRITE(6,1) (RMA(MT),MT=1,3),RMR
PF=PL(1)*FCF/DT
QF=PL(2)*FCF/DT
RF=PL(3)*FCF/DT
SF=PF+QF+RF
WRITE(6,10) PF,QF,RF,SF
DO 200 I=1,3
PP(4,I)=0.0
QQ(4,I)=0.0
RR(4,I)=0.0
SS(4,I)=0.0
TT(4,I)=0.0
UU(4,I)=0.0
P1(4,I)=0.0
Q1(4,I)=0.0
R1(4,I)=0.0
DO 150 MT=1,3
PP(MT,I)=CTI(MT,I)*DD(MT)/RMR
QQ(MT,I)=CTR(MT,I)*DD(MT)/RMR
SS(MT,I)=CNI(MT,I)*DD(MT)/RMR
TT(MT,I)=CNR(MT,I)*DD(MT)/RMR
P1(MT,I)=PP(MT,I)+SS(MT,I)
Q1(MT,I)=QQ(MT,I)+TT(MT,I)
RR(MT,I)=PP(MT,I)+QQ(MT,I)
UU(MT,I)=SS(MT,I)+TT(MT,I)
R1(MT,I)=P1(MT,I)+Q1(MT,I)
PP(4,I)=PP(4,I)+PP(MT,I)
QQ(4,I)=QQ(4,I)+QQ(MT,I)
RR(4,I)=RR(4,I)+RR(MT,I)
SS(4,I)=SS(4,I)+SS(MT,I)
TT(4,I)=TT(4,I)+TT(MT,I)
UU(4,I)=UU(4,I)+UU(MT,I)
P1(4,I)=P1(4,I)+P1(MT,I)
Q1(4,I)=Q1(4,I)+Q1(MT,I)
R1(4,I)=R1(4,I)+R1(MT,I)
150 CONTINUE
WRITE(6,12) (WD(J,I),J=1,2), (PP(K,I),QQ(K,I),RR(K,I),K=1,4)
WRITE(6,14) (SS(K,I),TT(K,I),UU(K,I),K=1,4)
WRITE(6,16) (P1(K,I),Q1(K,I),R1(K,I),K=1,4)
200 CONTINUE
AA=COSANG
BB=SINANG
DO 300 I=4,5
DO 250 K=1,4
PP(K,4)=AA*PP(K,1)+BB*PP(K,2)
QQ(K,4)=AA*QQ(K,1)+BB*QQ(K,2)
RR(K,4)=AA*RR(K,1)+BB*RR(K,2)
SS(K,4)=AA*SS(K,1)+BB*SS(K,2)
TT(K,4)=AA*TT(K,1)+BB*TT(K,2)
UU(K,4)=AA*UU(K,1)+BB*UU(K,2)
P1(K,4)=AA*P1(K,1)+BB*P1(K,2)
Q1(K,4)=AA*Q1(K,1)+BB*Q1(K,2)

```

PT10400

PT10460

PT10470

PT10480

PT10490

PT10630

PT10640

PT10650

PT10660

PT10680

PT10690

PT10700

PT10710

PT10720

PT10730

PT10740

PT10750

```

250 R1(K,4)=AA*R1(K,1)+BB*R1(K,2)
WRITE(6,12) (WD(J,I),J=1,2),(PP(K,4),QQ(K,4),RR(K,4),K=1,4)
WRITE(6,14) (SS(K,4),TT(K,4),UU(K,4),K=1,4)
WRITE(6,16) (P1(K,4),Q1(K,4),R1(K,4),K=1,4)
AA=-SINANG
BB=COSANG
300 CONTINUE
HD=HTP/DT
PA(4)=0.0
PB(4)=0.0
PC(4)=0.0
DO 400 MT=1,3
PA(MT)=HTI(MT)*RMA(MT)*RNU(MT)*HD/RMR
PB(MT)=HTR(MT)*RMA(MT)*RNU(MT)*HD/RMR
PC(MT)=PA(MT)+PB(MT)
PA(4)=PA(4)+PA(MT)
PB(4)=PB(4)+PB(MT)
PC(4)=PC(4)+PC(MT)
400 CONTINUE
WRITE(6,18) (PA(I),PB(I),PC(I),I=1,4)
RETURN
END

```

PT10760

PT10800

PT10810

PT10820

PT10830

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PT10950

PT10960

```

SUBROUTINE PRINT2(AKN,XSTART,DT,RNU,RMA,DRF,PCF,HTP,UTLI,UTTI,VTSI
1,HTSI,DELANG,NWEDG,XS,XCB,YCB,HTS,NTS,UTL,UTT,VTS,I2,I3,IP)
DIMENSION RMA(1),RNU(1),XS(1),XCB(1),YCB(1)
DIMENSION HTS(3,I2,I3),NTS(3,I2,I3),UTL(3,I2,I3)
DIMENSION UTT(3,I2,I3),VTS(3,I2,I3),UTLI(3,I2,I3),UTTI(3,I2,I3)
DIMENSION VTSI(3,I2,I3),HTSI(3,I2,I3)
COMMON /FIFTH/ND

```

PT20060

PT20070

THE PURPOSE OF THIS SUBROUTINE IS TO PRINT OUT THE DISTRIBUTION
ON SURFACE OF THE SURFACE COEFFICIENTS

PT20080

PT20090

PT20100

PT20110

PT20120

PT20130

PT20140

PT20150

PT20160

FORMATS

```

8 FORMAT(/1X,45('*'),' DISTRIBUTION ON SURFACE ',45('*')/71X,'INC.
1 TOT. INC. TOT. INC. TOT. '/11X,'SEGMENT GEOMETRY',
214X,'MOL. MOLE SAMP NUM. SKIN SKIN PRES- PRES-
3 HEAT HEAT'/' NO. CENTER DELX CENTER DELANG',4X,'WGHT. F
4RACT.',10X,
5 'FLUX PRCTN PRCTN SURE SURE TRNSF TRNSF')
10 FORMAT(1X,I3,F8.3,F7.3,F9.3,F8.3,1X,2F8.4,I6,7F8.4)
11 FORMAT(37X,F8.4,' 1.0000',I6,7F8.4)

```

PT20230

PT20240

PT20250

PT20260

```

RMR=0.0
DO 50 MT=1,IP
50 RMR=RMR+RMA(MT)*RNU(MT)

```

```
WRITE(6,8)
I=0
DO 110 N=2,ND
DTY=DT*YCB(N)/180.
P=XS(N)
Q=2.*((XCB(N)-XSTART)*AKN-XS(N))
J=0
```

PT20280
PT20320
PT20330
PT20340
PT20350
PT20380

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```
R=-.5*DELANG
PMLT=FCF/(DTY*DELANG)
QMLT=DRF/(DTY*DELANG)
SMLT=HTF/(DTY*DELANG)
DO 100 K=1,NWEDG
```

```
R=R+DELANG
```

```
I=I+1
```

```
J=J+1
```

PT20490
PT20500

```
M3=0
```

```
P3=0.0
```

```
Q3=0.0
```

```
Q4=0.0
```

```
R3=0.0
```

```
R4=0.0
```

```
S3=0.0
```

```
S4=0.0
```

```
DO 90 MT=1,IP
```

```
M1=NTS(MT,N,J)
```

```
M3=M3+M1
```

```
P1=NTS(MT,N,J)*PMLT*RNU(MT)
```

```
P3=P3+P1
```

```
Q1=SQRT(UTLI(MT,N,J)**2+UTTI(MT,N,J)**2)*RNU(MT)*RMA(MT)*QMLT/RMR
```

```
Q2=SQRT(UTL(MT,N,J)**2+UTT(MT,N,J)**2)*RNU(MT)*RMA(MT)*QMLT/RMR
```

```
Q3=Q3+Q1
```

```
Q4=Q4+Q2
```

```
R1=VTSI(MT,N,J)*RNU(MT)*RMA(MT)*QMLT/RMR
```

```
R2=VTS(MT,N,J)*RNU(MT)*RMA(MT)*QMLT/RMR
```

```
S1=HTSI(MT,N,J)*RNU(MT)*RMA(MT)*SMLT/RMR
```

```
S2=HTS(MT,N,J)*RNU(MT)*RMA(MT)*SMLT/RMR
```

```
R3=R3+R1
```

```
R4=R4+R2
```

```
S3=S3+S1
```

```
S4=S4+S2
```

```
90 WRITE(6,10)I,P,Q,R,DELANG,RMA(MT),RNU(MT),M1,P1,Q1,Q2,R1,R2,S1,S2
```

```
WRITE(6,11)RMR,M3,P3,Q3,Q4,R3,R4,S3,S4
```

```
100 CONTINUE
```

```
110 CONTINUE
```

```
RETURN
```

```
END
```

PT20660
PT20670
PT20680
PT20690

```
SUBROUTINE PRINT4(MSP,CHI,RNU,I,TRP,PDN,WTM,DB,NS,TMP,XV,  
1YV,ZV,KS,NB,XC,YC,ZC)
```

```
INTEGER*2 NB,NS
```

```
DIMENSION PDN(1),RNU(1),CHI(1),WTM(1),TMP(I,1),TRP(I,1)
```

```
DIMENSION DB(I,1),NB(I,1),XV(I,1),YV(I,1),ZV(I,1),DBT(3),NS(I,1)
```

```
DIMENSION XC(1),YC(1),ZC(1)
```

```
COMMON /SECND/BW,BH,RMP,RMN,RMF
```

```
COMMON /PORTH/NBX,RM,XR
```


-----PT50060
THE PURPOSE OF THIS SUBROUTINE IS TO PRINT OUT THE INSTANTANEOUS PT50070
FLOW-FIELD PROPERTIES. PT50080
-----PT50090

FORMATS

1 FORMAT(/1X,45('*'),' INSTANTANEOUS FLOW FIELD INFORMATION ',45('*') PT50100
1')) PT50110
2 FORMAT(/2X,'HORIZONTAL NUMBER=',I3,3X,'VERTICAL NUMBER=',I3,3X,'HO PT50120
1RIZONTAL POSITION=',F8.5,3X,' RADIAL POSITION =',F8.5/2X,'BOX# PT50130
2ANGLE SAMP DENSITY MACH NO X VEL. Y VEL. Z VEL. T(KIN) T(RO PT50140
3T) TEMP.',14X,'MOLE FRACTIONS') PT50150
3 FORMAT(/1X,46('*'),' ACCUMULATED FLOW FIELD INFORMATION ',46('*') PT50160
1) PT50170
4 FORMAT(1X,I4,E11.3,I6,8F8.3,3X,3E11.3)

*****PT50220
PT50230
PT50240
PT50250

IF(KS.EQ.0) WRITE(6,1)
IF(KS.NE.0) WRITE(6,3)
DO 40 MT=1,3
40 DBT(MT)=0.0
FDA=0.
CHT=0.
DO 50 MT=1,MSP
CHT=CHT+CHI(MT)*RNU(MT)
50 FDA=FDA+FDN(MT)*WTM(MT)
YCT=0.0
XCT=0.0
DO 110 N=1,NBX
IF((XC(N).EQ.XCT).AND.(YC(N).EQ.YCT)) GO TO 52
XCT=XC(N)
YCT=YC(N)
XXT=XCT/XR
YYT=YCT/RM
IXC=XCT/BW + 1
IYC=YCT/BH + 1
WRITE(6,2) IXC,IYC,XXT,YYT
52 ZCT=ZC(N)
NSAMP=0
DBA=0.
XVM=0.
YVM=0.
ZVM=0.
TMPM=0.
TRPM=0.
E=0.
F=0.
DO 100 MT=1,MSP
NSAMP=NSAMP+NS(MT,N)

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PT50290

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```

XVM=XVM+XV(MT,N)*RNU(MT)*WTM(MT)*NB(MT,N)
YVM=YVM+YV(MT,N)*RNU(MT)*WTM(MT)*NB(MT,N)
ZVM=ZVM+ZV(MT,N)*RNU(MT)*WTM(MT)*NB(MT,N)
DBA=DBA+DB(MT,N)*WTM(MT)
TMPM=TMPM+WTM(MT)*RNU(MT)*TMP(MT,N)*NB(MT,N)
TRPM=TRPM+RNU(MT)*NB(MT,N)*TRP(MT,N)
E=E+WTM(MT)*RNU(MT)*NB(MT,N)
100 P=P+RNU(MT)*NB(MT,N)
DBA=DBA/FDA
IF(E.EQ.0.0) GO TO 55
XVM=XVM/E
YVM=YVM/E
ZVM=ZVM/E
VS=XVM**2+YVM**2+ZVM**2
TMPM=TMPM/E-VS
TRPM=TRPM/P
55 CONTINUE
TTM=(TMPM+TRPM)/(2.5+CHT)
TMPM=TMPM/1.5
IF(CHT.NE.-1.) TRPM=TRPM/(1.+CHT)
AMS=SQRT(VS)
IF(TTM.GT.0.) AMS=SQRT((5.+2.*CHT)*VS/(TTM*(3.5+CHT)))
CCZ=COS(ZCT/57.29578)
SCZ=SQRT(1.-CCZ**2)
RVM=ZVM*SCZ-YVM*CCZ
TVM=YVM*SCZ+ZVM*CCZ
DO 60 MT=1,MSP
DBT(MT)=RNU(MT)*NB(MT,N)
IF(F.NE.0.) DBT(MT)=DBT(MT)/F
60 CONTINUE
WRITE(6,4) N,ZCT,NSAMP,DBA,AMS,XVM,RVM,TVM,TMPM,TRPM,TTM,(DBT(J),
1J=1,3)
110 CONTINUE
RETURN
END
ENTRY
CONTRL NAME='INTE','RNAL',TITLE=' PAR','ABOL','A AT',' 95K','M M','ON. ',
DEBUG=.F.,.F.,.T.,NEW=.F.,SAVE=.T.,ICOPY=0,REDO=.T. &END
TIMES DTM=.005,ITS=6,ITP=6,TST=2,TLIM=12 &END
FLOREF LLH=2000,MNH=5000,MNB=150,MSP=1,MET=0,U=7485.9,ANGLE=28.,RNU=1.,2*0.,
RMA=28.94,0.,0.,TF=195.51,DENF=2.52E+19 &END
MOLEC TRF=1000,DIR=3.5E-19,ETA=.104,PHI=0.0,CHI=-1.,ACR=.001 &END
SHAPES BODY=0.0,1000.,.00235 &END
SHAPES BODY=.0025,555.,0.0,2*1.0 &END
SHAPES BODY=.0050,345.,0.0,2*1.0 &END
SHAPES BODY=.0100,300.,0.0,2*1.0 &END
SHAPES BODY=.0200,300.,0.0,2*1.0 &END
SHAPES BODY=.0300,300.,0.0,2*1.0 &END
SHAPES BODY=.0400,300.,0.0,2*1.0 &END
SHAPES BODY=.0500,300.,0.0,2*1.0 &END
SHAPES BODY=.0600,300.,0.0,2*1.0 &END
SHAPES BODY=.0700,300.,0.0,2*1.0 &END
SHAPES BODY=.0800,300.,0.0,2*1.0 &END
SHAPES BODY=.0870,300.,1.0,2*1.0 &END
GEOM NWEDG=2,NW=20,NH=3, &END

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PT50470

PT50480

PT50490

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;INCUPL FLUXIN=2.1429,PCOL=1.0,RMP=0.3,JV=22,KMX=3 &END
;INOUT VARG=0.,1.,2.,3.,4.,5.,6.,7.,8.,9.,10.,11.,12.,13.,14.,15.,16.,17.,18.,
19.,20.,21.,CURV=0.0,.070,.170,.282,.369,.459,.537,.599,.656,.710,.750,.785,
.815,.845,.872,.900,.922,.951,.975,.988,.996,1.00, &END
;INOUT VARG=-20.,-19.,-17.,-15.,-13.,-11.,-9.,-7.,-5.,-3.,-1.,1.,3.,5.,7.,9.,
11.,13.,15.,17.,19.,CURV=4*0.,.003,.013,.036,.084,.149,.250,.406,.611,.762,
.871,.932,.962,.984,.995,.999,3*1.0, &END
;INOUT VARG=-20.,-19.,-17.,-15.,-13.,-11.,-9.,-7.,-5.,-3.,-1.,1.,3.,5.,7.,9.,
11.,13.,15.,17.,19.,CURV=4*0.0,.003,.013,.036,.084,.149,.250,.406,.611,.762,
.871,.932,.962,.984,.995,.999,3*1.0, &END
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